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# GEOMETRIC AND ERGODIC PROPERTIES OF THE STABLE FOLIATION

## Ursula HAMENSTÄDT

In 1983 appeared an article of **Lucy Garnett** in the Journal of Functional Analysis ([G]) in which she studies ergodic properties of a foliation  $\mathcal{F}$  on a compact manifold N. Principal assumption is that for every  $x \in N$  the leaf  $\mathcal{F}(x)$  of  $\mathcal{F}$  through x is a smoothly immersed submanifold of N depending continuously on  $x \in N$  in the  $C^{\infty}$ -topology. (In her paper she only considers smooth foliations, but her arguments immediately carry over to foliations  $\mathcal{F}$  satisfying the assumptions just mentioned, see [Y]).

Any smooth Riemannian metric g on N restricts to a leafwise smooth Riemannian metric  $g_{\mathcal{F}}$  on the tangent bundle  $T\mathcal{F}$  of  $\mathcal{F}$ . With respect to this metric the leaves of  $\mathcal{F}$  are smooth Riemannian manifolds of bounded geometry. In particular each leaf carries a natural Laplacean, and these Laplaceans group together to define a global second order differential operator  $\Delta_{\mathcal{F}}$  on N with continuous coefficients which is leafwise elliptic.

For every  $x \in N$  the Laplacean of  $g_{\mathcal{F}}$  on the leaf  $\mathcal{F}(x)$  of  $\mathcal{F}$  through x induces a Brownian motion on  $\mathcal{F}(x)$ , described by the heat kernel  $p(x, y, t)(y \in \mathcal{F}(x), t > 0)$ and the Lebesgue measure  $\lambda_{\mathcal{F}}$  on  $\mathcal{F}(x)$  induced by  $g_{\mathcal{F}}$ . For each t > 0 we now obtain a Borel-probability measure  $\omega_t$  on N whose support equals the closure  $\overline{\mathcal{F}(x)}$  of  $\mathcal{F}(x)$  in N by defining

$$\omega_t(A) = \frac{1}{t} \int_0^t \left( \int_A p(x, y, s) \ d\lambda_{\mathcal{F}}(y) \right) \ ds.$$

This measure is the time-t-average of the diffusions of the Dirac mass at x. Since  $\overline{\mathcal{F}(x)} \subset N$  is compact we can find a sequence  $\{t_j\}_j$  such that  $t_j \to \infty(j \to \infty)$  and such that the measures  $\omega_{t_j}$  converge weakly to a Borel-probability measure  $\omega$  on  $\overline{\mathcal{F}(x)} \subset N$ . This measure  $\omega$  is *stationary* for the process obtained by considering simultaneously all Brownian motions on all leaves of  $\mathcal{F}$ .

More precisely, if  $P^x$  denotes the Wiener measure on paths defined by Brownian

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#### **U. HAMENSTÄDT**

motion on  $\mathcal{F}(x)$  with starting point x then the measure P on the space of paths  $\Omega$  in N which is defined by  $P(A) = \int P^x(A) d\omega(x)$  is invariant under the one-parameter group  $\{T^t \mid t \geq 0\}$  of shift transformations  $T^t\xi(s) = \xi(s+t)$  (see [G] and the survey [Y]).

The stationary measure  $\omega$  is also called a *harmonic measure* for the operator  $\Delta_{\mathcal{F}}$  since it is characterized by the property that  $\int \Delta_{\mathcal{F}}(f) d\omega = 0$  for every smooth function f on N. It disintegrates locally into a transversal sum of leaf measures, where almost every leaf measure is a positive harmonic function times the Riemannian leaf measure ([G]).

In contrast to the case of the trivial foliation (dim  $\mathcal{F} = \dim N$ ) a harmonic measure for  $\Delta_{\mathcal{F}}$  needs not be unique. If  $\mathcal{F}$  has two distinct compact leaves  $\mathcal{F}(x_1), \mathcal{F}(x_2)$  then the normalized Lebesgue measures on these leaves are harmonic measures for  $\Delta_{\mathcal{F}}$  which are mutually singular.

However there are nontrivial foliations for which a harmonic measure is unique. Once again, the first example of such a foliation was described by Garnett ([G]).

Namely let M be a compact Riemannian manifold of negative sectional curvature. The geodesic flow  $\Phi^t$  is a smooth dynamical system on the unit tangent bundle  $T^1M$  of M generated by the geodesic spray X.

There are four  $\Phi^t$ -invariant Hölder continuous foliations on  $T^1M$  with smooth leaves which depend continuously in the  $C^{\infty}$ -topology on the points in  $T^1M$  ([S]). These foliations can be described as follows: Let d be a distance on  $T^1M$  defined by a smooth Riemannian metric. Then for  $v \in T^1M$  the leaf  $W^{ss}(v)$  through v of the strong stable foliation  $W^{ss}$  is the set  $\{w \in T^1M \mid d(\Phi^t v, \Phi^t w) \to 0 \ (t \to \infty)\}$ . Its tangent bundle  $TW^{ss}$  is a Hölder continuous subbundle of  $TT^1M$ . The tangent bundle  $TW^s$  of the stable foliation  $W^s$  is  $TW^s = \mathbb{R}X \oplus TW^{ss}$ , and the strong unstable foliation  $W^{su}$  (resp. the unstable foliation  $W^u$ ) is the image of  $W^{ss}$ (resp.  $W^s$ ) under the flip  $v \to -v$  of  $T^1M$ .

The canonical projection  $P: T^1M \to M$  maps each leaf of  $W^s$  locally diffeomorphically onto M. Thus the Riemannian metric on M lifts to a Riemannian metric  $g^s$  on  $TW^s$  which gives rise to a *stable Laplacean*  $\Delta^s$  along the leaves of  $W^s$  whose coefficients as a global operator on  $T^1M$  are Hölder continuous.

Garnett showed ([G]) that if M is a surface of constant curvature, then  $\Delta^s$  admits a unique harmonic measure. However her arguments are also valid for an arbitrary compact negatively curved manifold M, a fact which was explicitly pointed out by Yue ([Y]). Ledrappier gave independently a proof using the same arguments ([L2]).

The above considerations indicate that the structure of the convex compact space of harmonic measures for a leafwise Laplacean  $\Delta_{\mathcal{F}}$  reflects ergodic properties of the foliation  $\mathcal{F}$ , but it might not depend in a sensitive way on the Riemannian metric on N used to define the operator  $\Delta_{\mathcal{F}}$ . Some additional evidence for this was given by Kaimanovich ([K]). To formulate his result, recall that a *completely invariant* transverse measure for a foliation  $\mathcal{F}$  is a measure defined on transversals T for  $\mathcal{F}$  and such that the following holds: If T, T' are transversals, if  $\varphi: T \to T'$  is a homeomorphism such that  $\varphi(x) \in \mathcal{F}(x)$  for all  $x \in T$  (i.e.  $\varphi$  is defined by sliding T along the leaves of  $\mathcal{F}$ ) then

56

 $\varphi$  maps the measure on T to the measure on T'. According to Plante ([Pl]), completely invariant measures exist if  $\mathcal{F}$  has sub-exponential growth. Any such invariant transverse measure can be combined with the Lebesgue measure on the leaves of  $\mathcal{F}$  to define a finite Borel-measure on N which we call a *completely invariant harmonic measure*. Now Kaimanovich showed the following ([K]):

THEOREM 1. — If  $\mathcal{F}$  has sub-exponential growth, then every harmonic measure  $\nu$  for  $\Delta_{\mathcal{F}}$  is a completely invariant harmonic measure, and  $\nu$ -almost every leaf of  $\mathcal{F}$  is Liouville.

The arguments of Kaimanovich go as follows: First he induces a notion of entropy for the leafwise diffusion, the so called *Kaimanovich entropy*  $h_K$  which depends on the choice of a stationary measure  $\nu$ . He shows that  $h_K = 0$  if and only if almost every leaf of  $\mathcal{F}$  is Liouville, i.e. does not admit nonconstant bounded harmonic functions.

It  $\mathcal{F}$  is of subexponential growth, then necessarily  $h_K = 0$  for every harmonic measure for  $\Delta_{\mathcal{F}}$ . Now let  $\nu$  be a harmonic measure for  $\Delta_{\mathcal{F}}$  and consider reversal of time of the diffusion induced by  $\Delta_{\mathcal{F}}$  with respect to  $\nu$ . Since  $h_K = 0$ , this reversal of time coincides with the diffusion itself, and hence  $\nu$  is a *self-adjoint harmonic measure* for  $\Delta_{\mathcal{F}}$ , i.e.

$$\int f(\Delta_{\mathcal{F}} u) \; d
u = \int u(\Delta_{\mathcal{F}} f) \; d
u$$

for all smooth functions f, u on  $T^1M$  (compare [H1]). But a self adjoint harmonic measure corresponds to constant harmonic functions on the leaves of  $\mathcal{F}$  in the description of harmonic measures by Garnett ([G]) and hence is completely invariant (see [H1] for a detailed discussion).

Consider now the strong stable foliation  $W^{ss}$  on  $T^1M$  as above, equipped with the restriction  $g^{ss}$  of the Riemannian metric  $g^s$  on  $TW^s$  and the induced strong stable Laplacean  $\Delta^{ss}$ . The foliation  $W^{ss}$  is of subexponential growth and by a classical result of Bowen and Marcus ([B-M]) it admits a unique transverse invariant measure (defined by conditionals of unstable manifolds of the Bowen-Margulis measure on  $T^1M$ ). Thus the above considerations imply.

COROLLARY. — The strong stable Laplacean  $\Delta^{ss}$  admits a unique harmonic measure.

In the terminology of Knieper ([Kn]) the harmonic measure for  $\Delta^{ss}$  is just the horospherical measure  $\nu$  given with respect to a local product structure by  $d\nu = d\lambda^s \times d\mu^{su}$  where  $\mu^{su}$  is a family of conditional measures on strong unstable manifolds for the Bowen-Margulis measure and  $\lambda^s$  is the family of Lebesgue measures on stable manifolds induced by the Riemannian metric  $g^s$ . Moreover, if we change the metric on the leaves of  $W^{ss}$  to obtain a new leafwise Laplacean  $\overline{\Delta}$ , then  $\overline{\Delta}$  admits again a unique harmonic measure which is self-adjoint and contained in the measure class of  $\nu$ .

To summarize the above considerations, we see that a leafwise Laplacean on a foliation  $\mathcal{F}$  on N of subexponential growth with the additional property that every leaf of  $\mathcal{F}$  is dense in N induces a leafwise diffusion process whose ergodic properties are easy to describe and from which we can not hope to extract geometric properties of  $\mathcal{F}$  or N.

### U. HAMENSTÄDT

For foliations of exponential growth, however, the situation is much more complicated and interesting. A principal and relatively well understood example is the stable foliation  $W^s$  of a compact negatively curved manifold M. In fact, if the leaf  $W^s(v)$ through  $v \in T^1 M$  does not contain a periodic orbit of the geodesic flow (and there are only countably many such leaves), then  $(W^s(v), g^s)$  is isometric to the universal covering  $\widetilde{M}$  of M. This means that we can study Brownian motion on  $\widetilde{M}$  by studying the diffusion induced by  $\Delta^s$  on the compact space  $T^1 M$ .

Recall that  $\Delta^s$  admits a unique harmonic measure  $\omega$ , and properties of  $\omega$  reflect geometric properties of M and  $\widetilde{M}$ . One of the first observations in this direction is due to Ledrappier ([L1]); it can be combined with deep results of Benoist, Foulon, Labourie ([B-F-L], [F-L]) and Besson, Courtois, Gallot ([B-C-G]) to show:

THEOREM 2. — The unique harmonic measure  $\omega$  for  $\Delta^s$  is invariant under the geodesic flow if and only if M is locally symmetric.

If  $\omega$  is  $\Phi^t$ -invariant, then  $\omega$  necessarily coincides with the Lebesgue Liouville measure  $\lambda$  on  $T^1M$ . An open question is whether M is locally symmetric if only  $\omega$  is contained in the Lebesgue measure class.

Following Ledrappier ([L1]), the measure  $\omega$  can explicitly be constructed as follows:

Denote by  $\Delta$  the Laplace operator on the universal covering  $\widetilde{M}$  of M. The operator  $\Delta$  is *weakly coercive* in the sense of Ancona ([A]), i.e. there is a number  $\varepsilon > 0$  such that  $\Delta + \varepsilon$  admits a positive superharmonic function (i.e. a positive function f such that  $(\Delta + \varepsilon)(f) \ge 0$ ). By the results of Ancona ([A])  $\Delta$  admits a Green's function and the Martin boundary of  $\Delta$  can naturally be identified with the *ideal boundary*  $\partial \widetilde{M}$  of  $\widetilde{M}$ . This means that for every  $\xi \in \partial \widetilde{M}$  and every  $x \in \widetilde{M}$  there is a unique minimal positive  $\Delta$ -harmonic function  $y \to K(x, y, \xi)$  on  $\widetilde{M}$  with *pole* at  $\xi$  which is normalized by  $K(x, x, \xi) = 1$ . The function  $K: \widetilde{M} \times \widetilde{M} \times \partial \widetilde{M} \to (0, \infty)$  is Hölder continuous.

Now recall that the stable foliation on  $T^1M$  lifts to a foliation on  $T^1\widetilde{M}$  which we denote again by  $W^s$ . This foliation defines a natural closed equivalence relation  $\sim$  on  $T^1\widetilde{M}$  by writing  $v \sim w$  if and only if  $w \in W^s(v)$ . The ideal boundary  $\partial \widetilde{M}$  of  $\widetilde{M}$  is then naturally homeomorphic to the quotient  $T^1\widetilde{M}/\sim$ . In other words, there is a natural projection  $\pi: T^1\widetilde{M} \to \partial \widetilde{M}$  such that for every  $\zeta \in \partial \widetilde{M}$  the pre-image  $\pi^{-1}(\zeta)$  is a leaf of  $W^s$ .

For every  $v \in T^1 \widetilde{M}$  the restriction to  $W^s(v)$  of the canonical projection  $P: T^1 \widetilde{M} \to \widetilde{M}$  is a diffeomorphism. Thus for every fixed  $x \in \widetilde{M}$  and every  $\zeta \in \partial \widetilde{M}$  the gradient of the logarithm of the function  $y \to K(x, y, \zeta)$  lifts to a vector field on  $\pi^{-1}(\zeta)$  not depending on the base-point x.

These vector fields group together to a Hölder continuous leafwise smooth section  $\widetilde{Y}$  of  $TW^s$  over  $T^1\widetilde{M}$  which is equivariant under the action of the fundamental group  $\pi_1(M)$  of M on  $T^1\widetilde{M}$  and hence projects to a Hölder continuous leafwise smooth section Y of  $TW^s$  over  $T^1M$ .

Let  $\mathcal{M}$  be the convex compact space of  $\Phi^t$ -invariant Borel-probability measures

on  $T^1M$  equipped with the weak\*-topology. For  $\eta \in \mathcal{M}$  denote by  $h_\eta$  the entropy of  $\eta$  (see [W]). The pressure of a Hölder continuous function f on  $T^1M$  is defined by  $pr(f) = \sup\{h_\eta - \int f \, d\eta \mid \eta \in \mathcal{M}\}$ . There is a unique measure  $\nu_f \in \mathcal{M}$ , the so called Gibbs equilibrium state of f, such that  $h_{\nu_f} = \int f \, d\nu_f = pr(f)$  (see [W]). The measure  $\nu_f$  admits a family  $\nu_f^{su}$  of conditional measures on strong unstable manifolds which transform under the geodesic flow via  $\frac{d}{dt}\Phi^t \circ \nu^{su}|_{t=0} = f + pr(f)$ .

Let again X be the geodesic spray and let Y be the section of  $TW^s$  over  $T^1M$ as above. Then the pressure of the function  $g^s(X,Y)$  is zero ([L1]) and the unique harmonic measure  $\omega$  for  $\Delta^s$  is of the form  $d\omega = d\lambda^s \times d\nu^{su}$  where  $\nu^{su}$  is a family of conditional measures of the Gibbs equilibrium state induced by  $g^s(X,Y)$ .

For  $v \in T^1 M$  denote now by  $P^v$  the Wiener measure on paths on  $W^s(v)$  induced by  $\Delta^s |_{W^s(v)}$  with starting point v. Let  $\tilde{v} \in T^1 \widetilde{M}$  be a lift of v to  $T^1 \widetilde{M}$  and let  $P^x$  be the Wiener measure on paths on  $\widetilde{M}$  induced by Brownian motion on  $\widetilde{M}$  with starting point  $x = P\tilde{v}$ . If A is a family of paths on  $W^s(v)$  starting at v, then A lifts to a unique family  $\widetilde{A}$  of paths on  $W^s(\tilde{v})$  starting at  $\tilde{v}$ , and we have  $P^x \{Pc \mid c \in \widetilde{A}\} = P^v(A)$ .

By a result of Prat ([P]), for  $P^x$ -almost every path c in  $\widetilde{M}$  the limit  $\lim_{t\to\infty} c(t)$  exists in  $\widetilde{M} \cup \partial \widetilde{M}$  and is contained in  $\partial \widetilde{M}$ . Thus  $P^x$  projects to a hitting measure  $\omega^x$  on  $\partial \widetilde{M}$  defined by  $\omega^x(A) = P^x\{c \mid c(\infty) \in A\}$ . The measures  $\omega^x, \omega^y$  for  $x, y \in \widetilde{M}$  are equivalent and do not have atoms. Moreover the above convergence is with positive speed, which means that  $\liminf_{t\to\infty} \frac{1}{t} \operatorname{dist}(c(0), c(t)) > 0$  for  $P^x$ -almost every path c.

While the result of Prat is valid for every simply connected Riemannian manifold  $\widetilde{M}$  of bounded negative curvature, more can be said for the universal covering of a compact space using methods from ergodic theory applied to the diffusion on  $(T^1M, \omega)$  induced by  $\Delta^s$ . Namely for  $w \in T_x^1 \widetilde{M}$  let  $\Theta_w$  be the *Busemann function* at  $\pi(w)$  normalized by  $\Theta_w(P\dot{w}) = 0$ . The lift of  $\Theta_w$  to  $(W^s(w), g^s)$  is a function whose gradient is just the negative -X of the geodesic spray X.

For  $v \in T^1M$  denote by trU(v) the trace of the second fundamental form of the horosphere  $PW^{ss}(v)$  at Pv, normalized to be positive. Then for every  $x \in \widetilde{M}, w \in T_x^1\widetilde{M}$  and  $P^x$ -almost every path c in  $\widetilde{M}$  the limit  $\lim_{t\to\infty} \frac{1}{t}\Theta_w(c(t))$  exists and equals  $l = \int (trU) d\omega$  (see [K], [L1]). This means that the asymptotic escape rate for a typical path c does not depend on c, moreover Brownian motion does not have a preferred escape direction. For the diffusion on  $T^1M$  induced by  $\Delta^s$  this shows that a typical paths follows (roughly) an orbit of the geodesic flow with positive speed, but in the negative direction (recall that the gradient of  $\Theta_w$  on  $W^s(w) \subset T^1\widetilde{M}$  is -X). We call such a diffusion a diffusion of positive escape and say also in short that  $\Delta^s$  is of positive escape with respect to its (unique) harmonic measure  $\omega$ .

Let now f be a minimal positive harmonic function on  $\widetilde{M}$  with pole at  $\zeta \in \partial \widetilde{M}$ . The diffusion induced by the operator  $\Delta + 2\nabla \log f$  on  $\widetilde{M}$  is a conditional Brownian motion, and a typical path c satisfies  $\lim_{t\to\infty} c(t) = \zeta$  in  $\widetilde{M} \cup \partial \widetilde{M}$ .

The collection of all those diffusions given by all possible positive minimal harmonic functions can be described by the diffusion induced by  $\Delta^s + 2\tilde{Y}$  on  $T^1\tilde{M}$ . The

#### U. HAMENSTÄDT

operator  $\Delta^s + 2\tilde{Y}$  projects to the operator  $\Delta^s + 2Y$  on  $T^1M$  (notations as above), and the diffusion induced by  $\Delta^s + 2Y$  can again be studied using ergodic theory on the compact space  $T^1M$ . Now a typical path of  $\Delta^s + 2Y$  follows a flow line of the geodesic flow with positive speed in the positive direction. We say that this diffusion is of *negative escape* and call  $\Delta^s + 2Y$  of *negative escape*. Observe here that this qualitative behaviour may depend on the choice of a harmonic measure for  $\Delta^s + 2Y$ . One particular harmonic measure for  $\Delta^s + 2Y$  is just  $\omega$ , the harmonic measure for  $\Delta^s$ . In fact we have ([H1]):

LEMMA. — The reversal of time of the diffusion induced by  $\Delta^s$  on  $(T^1M, \omega)$  is the diffusion induced by  $\Delta^s + 2Y$  on  $(T^1M, \omega)$ .

The above considerations are valid in a larger context. Let now g be any smooth Riemannian metric on  $T^1M$  and denote by  $\Delta$  the leafwise Laplacean along the stable foliation induced by g. Recall that g induces an isomorphism of  $TW^s$  with its dual  $T^*W^s$ . Let Z be Hölder continuous section of  $TW^s$  which is differentiable along the leaves of  $W^s$  and such that its restriction to every leaf of  $W^s$  is dual with respect to g to a closed one-form along the leaf. Write  $L = \Delta + Z$  and call L weakly coercive if there is  $v \in T^1M$ such that the restriction of L to  $W^s(v) \sim \widetilde{M}$  is weakly coercive in the sense of Ancona. Observe that  $\Delta^s + 2Y$  is an operator of this type which is weakly coercive.

Call an operator L of this form of positive escape resp. negative escape if a typical path with respect to every harmonic measure for L follows (roughly) an orbit of the geodesic flow with positive speed in the negative direction (resp. the positive direction). As we have seen,  $\Delta^s$  is of positive escape, and the fact that  $\Delta^s + 2Y$  is of negative escape is contained in the following theorem ([H1]):

THEOREM 3.

- 1) If pr(g(X,Z)) > 0 then  $L = \Delta + Z$  admits a unique harmonic measure  $\nu$ . Moreover L is weakly coercive, of positive escape, and the Kaimanovich entropy  $h_K$  of the diffusion induced by L on  $(T^1M, \nu)$  is positive.
- 2) If pr(g(X, Z)) = 0 then L admits a unique self-adjoint harmonic measure  $\nu$ . Moreover L is not weakly coercive, of zero escape, and the Kaimanovich entropy of the diffusion induced by L on  $(T^1M, \nu)$  vanishes.

3) If pr(g(X, Z)) < 0 then L is weakly coercive, of negative escape with respect to every harmonic measure  $\nu$ , and the Kaimanivich entropy vanishes.

In the case pr(g(X,Z)) < 0 a harmonic measure for L needs not be unique; we'll describe an example for this in Theorem 5 below.

Operators of the above type are suitable to study eigenfunctions of the Laplacean  $\Delta$  on  $\widetilde{M}$ . Namely let  $\delta_0 > 0$  be the bottom of the positive spectrum for  $\Delta$ . Ledrappier related  $\delta_0$  to the *topological entropy* h of the geodesic flow on  $T^1M$ ; he showed:

THEOREM 4 [L3]. —  $\delta_0 \leq \frac{\hbar^2}{4}$ , with equality if and only if M is asymptotically harmonic and hence locally symmetric.

For  $\varepsilon > 0$  the operator  $\Delta_{\varepsilon} = \Delta + \delta_0 - \varepsilon$  on  $\widetilde{M}$  is weakly coercive and hence as before its Martin kernel is a Hölder continuous function  $K_{\varepsilon}: \widetilde{M} \times \widetilde{M} \times \partial \widetilde{M} \to (0, \infty)$  which gives rise to a Hölder continuous section  $\xi_{\varepsilon}$  of  $TW^s$  over  $T^1M$  as before.

Let p(x, y, t) be the heat kernel of  $\Delta$  and let f be a positive solution of the equation  $\Delta_{\varepsilon}=0$ . Then the fundamental solution of the parabolic equation  $\frac{\partial}{\partial t} - \Delta - 2\nabla \log f = 0$  equals the function  $(x, y, t) \in \widetilde{M} \times \widetilde{M} \times (0, \infty) \rightarrow e^{(\delta_0 - \varepsilon)t} p(x, y, t) f(y) / f(x)$ .

In other words, we may study the operator  $\Delta + 2\nabla \log f$  without zero order term to find properties of  $\Delta_{\varepsilon}$ .

Recall the definition of the section  $\tilde{\xi}_{\varepsilon}$  of  $TW^s$  over  $T^1\widetilde{M}(\varepsilon \in (0, \delta_0])$  from above. For every  $v \in T^1\widetilde{M}$  the restriction of  $\Delta^s + 2\xi_{\varepsilon}$  to  $W^s(v)$  is an operator of the kind just described. But  $\tilde{\xi}_{\varepsilon} \mid_{W^s(v)}$  is the gradient of the logarithm of a minimal positive  $\Delta_{\varepsilon}$ -harmonic function on  $W^s(v) \sim \widetilde{M}$  and hence a typical path of the diffusion induced by  $\Delta^s + 2\tilde{\xi}_{\varepsilon} \mid_{W^s(v)}$  converges as  $t \to \infty$  to the distinguished point  $\pi(v) \in \partial \widetilde{M}$  (with positive speed). Thus the operator  $\Delta^s + 2\xi_{\varepsilon}$  on  $T^1M$  falls into category 3) in Theorem 3 above. In fact it admits many harmonic measures (see [H2]):

THEOREM 5. — Let  $\bar{\eta}$  a Gibbs equilibrium state of a flip invariant Hölder continuous function on  $T^1M$ . Let  $\bar{\eta}^{su}$  be a family of conditional measures on strong unstable manifolds for  $\bar{\eta}$ . Then for every  $\varepsilon \in (0, \delta_0)$  the operator  $\Delta^s + 2\xi_{\varepsilon}$  admits a harmonic measure in the measure class of  $\eta$  where  $d\eta = d\lambda^s \times d\bar{\eta}^{su}$ .

Fix now a point  $w \in T^1 \widetilde{M}$  and consider the restriction of  $\tilde{\xi}_{\varepsilon}$  to  $W^s(w)$ . The operators  $\Delta_{\varepsilon}$  satisfy a uniform infinitesimal Harnack inequality, independent of  $\varepsilon \in (0, \delta]$ , and hence there is a sequence  $\{\varepsilon_i\} \subset (0, \delta]$  such that  $\varepsilon_i \to 0$   $(i \to \infty)$  and that  $\tilde{\xi}_{\varepsilon_i} \mid_{W^s(w)}$  converge uniformly on compact subsets of  $W^s(w) \sim \widetilde{M}$  to a vector field  $\tilde{\xi}_0$  on  $W^s(w) \sim \widetilde{M}$ . Then  $\tilde{\xi}_0$  is the gradient of the logarithm of a positive  $\Delta_0 = \Delta + \delta_0$ -harmonic function on  $\widetilde{M}$  is in fact a combination of function of this kind; it is contained in [H2]:

THEOREM 6. — The sections  $\tilde{\xi}_{\varepsilon}$  of  $TW^s$  over  $T^1\widetilde{M}$  converge uniformly to a section  $\tilde{\xi}_0$ . The restriction of  $\tilde{\xi}_0$  to a leaf  $W^s(w)$  is the gradient of the logarithm of a minimal positive  $\Delta_0$ -harmonic function on  $W^s(w) \sim \widetilde{M}$  with pole at  $\pi(w)$ . Every minimal positive  $\Delta_0$ -harmonic function is of this kind.

The vector fields  $\xi_0$  projects to a section  $\xi_0$  of  $TW^s$  over  $T^1M$ . The operator  $\Delta^s + 2\xi_0$  admits a unique self-adjoint harmonic measure.

The above describes the minimal Martin boundary for  $\Delta_0$ ; it can be identified with the ideal boundary  $\partial \widetilde{M}$  of  $\widetilde{M}$ . We do not know however how the full Martin boundary of  $\Delta_0$  looks like. We also do not know whether the Martin topology for the minimal Martin boundary  $\partial \widetilde{M}$  of  $\Delta_0$  induces on  $\partial \widetilde{M}$  a Hölder structure compatible with the usual Hölder structure of  $\partial \widetilde{M}$  (which is the case for the Martin boundary of the operators  $\Delta_{\varepsilon}$ for  $\varepsilon > 0$ ).

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