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## A. Melin <br> The Lippman-Schwinger equation treated as a characteristic Cauchy problem

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## ÉQUATIONS AUX DÉRIVÉES PARTIELLES

## THE LIPPMAN-SCHWINGER EQUATION TREATED AS A CHARACTERISTIC CAUCHY PROBLEM.

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# THE LIPPMANN-SCHWINGER EQUATION TREATED AS A 

## CHARACTERISTIC CAUCHY PROBLEM.

by Anders Melin

## Introduction.

We shall consider a real-valued function $v \in C^{\infty}\left(\mathbf{R}^{n}\right)$ when $n>1$ is odd. In order to have sufficiently regular scattering data associated to the Schrödinger operator $H_{v}=-\Delta_{x}+v(x)$ we shall assume that $v$ satisfies the following short-range condition:

$$
\begin{equation*}
\int_{\mathbf{R}^{n}}(1+|x|)^{2-n+|\alpha|}\left|v^{(\alpha)}(x)\right| d x<\infty \quad \text { for any } \alpha \tag{1}
\end{equation*}
$$

The class of such potentials will be denoted $\mathcal{V}$. By using polar coordinates in the frequency variables one may write the the Lippmann-Schwinger equation on the form

$$
\begin{equation*}
\left(-\Delta_{x}+v(x)\right) \phi(x, \theta, k)=k^{2} \phi(x, \theta, k), \quad x \in \mathbf{R}^{n}, \theta \in S^{n-1}, k \in \mathbf{R} \tag{2}
\end{equation*}
$$

One has also to impose some condition on $\phi(x, \theta, k)$ as $|x| \rightarrow \infty$ in order to obtain a unique solution of (2). We shall always consider $\phi$ as a perturbation of the function $\phi_{0}(x, \theta, k)=e^{i k(x, \theta\rangle}$ which solves (2) when $v=0$. Moreover, $\phi$ will be a continuous function of $k \in \mathbf{R} \backslash 0$ with a meromorhic extension to the upper half-plane. If $0<\Im k$ is small then

$$
\phi=\phi_{0}-\left(H_{0}-k^{2}\right)^{-1}(v \phi),
$$

where $\left(H_{0}-k^{2}\right)^{-1}$ is the $L^{2}$ - bounded inverse of $H_{0}-k^{2}$. In the case of a compactly supported potential $v$ this leads to the formula

$$
\begin{equation*}
\phi(x, \theta, k)-\phi_{0}(x, \theta, k)=2^{-1}\left(\frac{4 \pi}{i k|x|}\right)^{(n-1) / 2} e^{i k|x|} T(k, x /|x|, \theta)+O\left(|x|^{-(n+1) / 2}\right) \tag{3}
\end{equation*}
$$

where $T$ is the scattering amplitude. We also remark that $\phi$ can be defined in terms of the distribution kernels of the wave operators $W_{ \pm}=\lim _{t \rightarrow \pm \infty} e^{i t H_{v}} e^{-i t H_{0}}$, and one often calls the solutions of (2) generalized eigenfunctions.

In this note we show how $\phi$, or rather its Fourier transform w.r.t. the variable $k$, can be obtained as the solution of a characteristic Cauchy problem for the differential operator $\Delta_{x}-\partial_{t}^{2}-v(x)$. This viewpoint will give us extra information about $\phi$ and enables us to prove that

$$
\begin{equation*}
e^{-i k\langle x, \theta\rangle} \phi(x, \theta, k)=1+\int_{0}^{\infty} w_{\theta}(x, t) e^{i k t} d t \tag{4}
\end{equation*}
$$

where $w_{\theta}(x, t)$ is a smooth function. In particular we shall recover an identity which is usually referred to as the miracle (cf [ $\mathrm{N} 1, \mathrm{~N} 2, \mathrm{C}]$ ). We also remark that part of the discussions here can be carried over to the case of more general short range potentials.

Construction of $\phi$ by means of intertwining operators.
We shall first consider the equation

$$
\begin{equation*}
\left(\Delta_{x}-\Delta_{y}-v(x)\right) A_{\theta}(x, y)=0 \tag{5}
\end{equation*}
$$

In [M5] it was proved that this equation has a solution which is supported in the set $\langle y-x, \theta\rangle \geq 0$ and given by a series

$$
\sum_{0}^{\infty} U_{N, \theta}(x, y)
$$

where

$$
\left(\Delta_{x}-\Delta_{y}\right) U_{N+1, \theta}(x, y)=v(x) U_{N}(x, y), \quad U_{0}(x, y)=\delta(x-y)
$$

In order to describe the regularity of the solution one introduces the set $\mathcal{P}_{\boldsymbol{\lambda}}$ of all seminorms

$$
p(U)=\sup _{x} \int_{\mathbf{R}^{n}} e^{-\lambda(y-x, \theta\rangle}\left|\left(\partial_{x}+\partial_{y}\right)^{\alpha}\left(\left\langle x, \partial_{x}\right\rangle+\left\langle y, \partial_{y}\right\rangle\right)^{\beta} U(x, y)\right| d y
$$

Then for each $v$ which satisfies (1) there is a $\lambda=\lambda_{v} \geq 0$ so that

$$
\begin{equation*}
\sum_{1}^{\infty} p\left(U_{N, \theta}\right)<\infty, \quad p \in \mathcal{P}_{\lambda} \tag{6}
\end{equation*}
$$

Moreover, for each $m \geq 0$ there is a positive integer $N(m)$ so that

$$
\begin{equation*}
\sum_{N(m)}^{\infty} p\left(\partial_{x}^{\alpha} \partial_{y}^{\beta} U_{N, \theta}\right)<\infty, \quad|\alpha+\beta| \leq m, p \in \mathcal{P}_{\lambda} \tag{7}
\end{equation*}
$$

In order to make the $U_{N, \theta}$ unique one also has to introduce some conditions at infinity which will exclude from the considerations functions which are constant in the direction of $(\theta, \theta)$. We shall not discuss these details here.

Next we introduce

$$
V_{N, \theta}(x, t)=\int_{\langle y-x, \theta\rangle=t} U_{N, \theta}(x, y) d y
$$

and we let $\mathcal{Q}_{\boldsymbol{\lambda}}$ be the family of semi-norms

$$
q(V)=\sup _{x} \int_{0}^{\infty} e^{-\lambda t}\left|\partial_{x}^{\alpha}\left(\left\langle x, \partial_{x}\right\rangle+t \partial_{t}\right)^{\beta} V(x, t)\right| d t
$$

It follows from (6) and (7) then that

$$
\begin{equation*}
\sum_{1}^{\infty} q\left(V_{N, \theta}\right)<\infty, \quad q \in \mathcal{Q}_{\lambda} \tag{6}
\end{equation*}
$$

and

$$
\begin{equation*}
\sum_{N(m)}^{\infty} q\left(\partial_{x}^{\alpha} \partial_{t}^{\beta} V_{N, \theta}\right)<\infty, \quad|\alpha|+\beta \leq m, q \in \mathcal{Q}_{\lambda} \tag{7}
\end{equation*}
$$

It was proved in [M5] that

$$
\phi(x, \theta, k)=\int A_{\theta}(x, y) e^{i k\langle y, \theta\rangle} d y
$$

Hence if $V_{\theta}(x, t)=\sum_{0}^{\infty} V_{N, \theta}(x, t)$ we must (in view of the definition of $V_{N, \theta}$ ) have

$$
\begin{equation*}
e^{-i k\langle x, \theta\rangle} \phi(x, \theta, k)=\int e^{i t k} V_{\theta}(x, t) d t \tag{8}
\end{equation*}
$$

It follows from (6)' that the integrand is continuous w.r.t. $x$ and integrable w.r.t. $t$ when $\Im k$ is large enough. Moreover, $t \geq 0$ in the support of $V_{\theta}$, and $V_{\theta}(x, t)=\delta(t)$ if $v=0$.

## The main result.

It follows immediately from (6) that $V_{\theta}(x, t)$ is a smooth function of $x$ and $t$ when $t>0$, and the next result implies that one may write $V_{\theta}(x, t)=\delta(t)+Y_{+}(t) w_{\theta}(x, t)$, where $Y_{+}$ is the Heaviside function and $w_{\theta}(x, t)$ is smooth when $x \in \mathbf{R}^{n}$ and $t \geq 0$.
Theorem 1. There is a positive number $\lambda$ such that

$$
\begin{equation*}
\sup _{x, \theta} \int_{+0}^{\infty} e^{-\lambda t}\left|\partial_{x}^{\alpha} \partial_{t}^{\beta} V_{\theta}(x, t)\right|\langle x\rangle^{-\beta} d t<\infty \tag{9}
\end{equation*}
$$

for any $\alpha$ and $\beta$. (Here $\langle x\rangle=\left(1+|x|^{2}\right)^{1 / 2}$.)
We have already seen that $V_{\theta}(x, t)$ is smooth when $t>0$, and it follows from (6) ${ }^{\prime}$ also that we need only consider the integral over the interval ( 0,1 ) in (9). Moreover, the estimates (7)' imply that it it suffices to prove a similar result for each of the $V_{N, \theta}$. Hence Theorem 1 results from the following
Theorem 1'. If $N \geq 1$, then

$$
\begin{equation*}
\sup _{x, \theta} \int_{+0}^{1}\langle x\rangle^{-\beta}\left|\partial_{x}^{\alpha} \partial_{t}^{\beta} V_{N, \theta}(x, t)\right| d t<\infty . \tag{10}
\end{equation*}
$$

We have now come to the point where one has to study the wave equation. In fact, the equation $\left(\Delta_{x}-\Delta_{y}\right) U_{N+1, \theta}(x, y)=v(x) U_{N, \theta}(x, y)$ implies that

$$
\begin{equation*}
\mathcal{L}_{\theta} V_{N+1, \theta}(x, t)=v(x) V_{N, \theta}(x, t) \tag{11}
\end{equation*}
$$

if $\mathcal{L}_{\theta}=\Delta_{x}-2\left\langle\theta, \partial_{x}\right\rangle \partial_{t}$. We observe that $\mathcal{L}_{\theta}$ is obtained from the wave operator $\Delta_{x}-\left(\partial_{t_{0}}\right)^{2}$, after the substitution $t_{0}=t+\langle x, \theta\rangle$. Hence $t \geq 0$ is a characteristic half-space for $\mathcal{L}_{\theta}$. We let $G_{\theta}$ be the image of the fundamental solution of $\Delta_{x}-2\left\langle\theta, \partial_{x}\right\rangle \partial_{t}$ with support in the set $t_{0} \geq 0$ under the substitution above. Then $t \geq 0$ in the support of $G_{\theta}$.

We shall consider approximate solutions of the equation $\mathcal{L}_{\theta} V(x, t)=v(x) V_{N, \theta}(x, t)$, which is solved by $V_{N+1, \theta}$. The construction of such a solution will be similar to the methods of geometrical optics used in microlocal analysis, and an exact solution will then be obtained after convolving the error term $\mathcal{L}_{\theta} V(x, t)-v(x) V_{N, \theta}(x, t)$ with some fundamental solution $Q_{\theta}$ of $\mathcal{L}_{\theta}$ with support in the set $t \geq 0$. The following result shows that one has to take $Q_{\theta}=G_{\theta}$ if one hopes to obtain good bounds for the solutions.

Proposition 2. Assume that $\mathcal{L}_{\theta} u(x, t)=0$ and that $t \geq 0$ in the support of $u$. If $u(x, t)$ is temperate w.r.t. $x$ then

$$
u(x, t)=\sum_{|\alpha| \leq \mu(t)} f_{\alpha}(t) x^{\alpha}
$$

where $f_{\alpha} \in \mathcal{D}^{\prime}(\mathbf{R})$ and the integer valued function $\mu(t)$ is locally bounded.
PROOF: We may assume that $\theta=e_{n}$. The function $G(x, y)=u\left(x, y_{n}-x_{n}\right)$ then solves the equation $\left(\Delta_{x}-\Delta_{y}\right) G(x, y)=0$, and $y_{n} \geq x_{n}$ in its support. The proof of Theorem 3.5 of [M5] shows then that

$$
G(x, y)=G\left(x, y_{n}\right)=\sum_{0}^{\infty} g_{j}\left(x^{\prime}, y_{n}-x_{n}\right) x_{n}^{j}
$$

where $x^{\prime}=\left(x_{1}, \ldots, x_{n-1}\right)$ and $g_{j} \in \mathcal{D}^{\prime}\left(\mathbf{R}^{n}\right)$. Then

$$
u(x, t)=\sum_{0}^{\infty} g_{j}\left(x^{\prime}, t\right) x_{n}^{j},
$$

where only finitely many of the $g_{j}$ are $\neq 0$ when $t$ stays in any bounded open set $\omega$. The equation $\left(\Delta_{x}-2 \partial_{x_{n}} \partial_{t}\right) u=0$ implies that

$$
\Delta_{x^{\prime}} g_{j}\left(x^{\prime}, t\right)-2(j+1) \partial g_{j+1}\left(x^{\prime}, t\right) / \partial t+(j+2)(j+1) g_{j+2}\left(x^{\prime}, t\right)=0, \quad j=0,1, \ldots
$$

Hence, when $t$ is in any $\omega$ as above, then $\Delta_{x^{\prime}}^{N} g_{j}\left(x^{\prime}, t\right)=0$ for any $j$ if $N$ is large enough. Since the $g_{j}$ are temperate in $x^{\prime}$, this implies that they are polynomials in this variable and the proposition follows.

Let $\Gamma_{0}$ be the cone $|x|=t_{0}$ and $\Gamma$ be its image under the substitution $t_{0}=t+\langle x, \theta\rangle$, i.e. $\Gamma$ is defined by $t=|x|-\langle x, \theta\rangle$. The half-plane $B: t \geq 0$ corresponds to $B_{0}: t_{0} \geq\langle x, \theta\rangle$, which intersects $\Gamma_{0}$ only along the ray $\left\{\left(t_{0} \theta, t_{0}\right) ; t_{0} \geq 0\right\}$. Hence $-\Gamma_{0}$ intersects $B_{0}$ only along the opposite ray $\left\{\left(t_{0} \theta, t_{0}\right) ; t_{0} \leq 0\right\}$, and any distribution $u_{0}$ supported in $B_{0}$ and vanishing over a conic neighbourhood of $\gamma=\mathbf{R}_{-} \theta$ has to vanish identically if it satisfies the wave equation $\left(\Delta_{x}-\left(\partial_{t_{0}}\right)^{2}\right) u_{0}\left(x, t_{0}\right)=0$. In fact, if $(x, s)$ is in the wave cone $\Gamma_{0},(y, t)$ belongs to the support of $u_{0}$ and $x+y, s+t$ belong to bounded sets, then $|y|$ can not tend to infinity unless $y /|y|$ tends to $-\theta$. In $(x, t)$ space this implies that $w=G_{\theta} * u$ is defined if $t \geq 0$ in the support of $u$ and $u$ vanishes over a conic neighbourhood of the ray $\gamma=\mathbf{R}_{-} \theta$, and $w$ is the unique solution of the equation $\mathcal{L}_{\theta} w=u$ in the space of such distributions.

In order to have $G_{\theta} * u$ defined on a larger space we introduce the following definition:
Definition 3. $\mathcal{D}_{G_{\theta}}^{\prime}$ is the space of all $u$ in $\mathcal{D}^{\prime}\left(\mathbf{R}^{n} \times \mathbf{R}\right)$ such that $\lim _{j \rightarrow \infty} G_{\theta} *\left(\chi_{j} u\right)$ exists for any sequence $\chi_{j} \in C_{0}^{\infty}\left(\mathbf{R}^{n} \times \mathbf{R}\right)$ such that $\left\|\chi_{j}\right\|_{L^{\infty}}$ is bounded, $\chi_{j}$ converges pointwise and $\left\|\chi_{j}^{(\alpha)}\right\|_{L^{\infty}} \rightarrow 0$ as $j \rightarrow \infty$ if $\alpha \neq 0$.

If $u \in \mathcal{D}_{G_{\theta}}^{\prime}$, and the sequence $\chi_{j}$ above tends to 1 then we define $G_{\theta} * u$ as the limit of $G_{\theta} *\left(\chi_{j} u\right)$. This limit is independent of the choices made. Moreover, $\mathcal{D}_{G_{\theta}}^{\prime}$ is invariant
under differentiation, and $u \rightarrow G_{\theta} * u$ is a left-inverse for $\mathcal{L}_{\theta}$ on this space in the sense that $u=G_{\theta} * \mathcal{L}_{\theta} u$ when $u \in \mathcal{D}_{G_{\theta}}^{\prime}$.

One can show that $u \in \mathcal{D}_{G_{\theta}}^{\prime}$ if $t \geq 0$ in its support and if it decays as $|x|^{-1-\varepsilon}$ over some conic neighbourhood of $\gamma$ for some positive $\varepsilon$. One can even allow less restrictive conditions on the decay of $u$, however, since we are dealing with potentials in the class $\mathcal{V}$, we shall only consider the following conditions:
Definition 4. Let $f \in C^{\infty}\left(\mathbf{R}^{n}\right)$. Then we say that $f \in \mathcal{V}_{\theta}$ if af $\in \mathcal{V}$ for any $a \in S^{0}\left(\mathbf{R}^{n}\right)$ such that the support of $a$ is contained in some cone $\varepsilon|x| \leq-\langle x, \theta\rangle$, where $\varepsilon>0$.

Here the condition that $a \in S^{0}\left(\mathbf{R}^{n}\right)$ means that $\langle x\rangle^{|\alpha|} a^{(\alpha)}(x)$ is bounded for any $\alpha$. It is easy to see that $\mathcal{V}$ and $\mathcal{V}_{\theta}$ are Fréchet spaces and they are also $S^{0}$ - modules.

We let $C^{\infty}\left(\overline{\mathbf{R}_{+}}\right) \otimes \mathcal{V}_{\theta}$ be the space of smooth maps from $\overline{\mathbf{R}_{+}}$to $\mathcal{V}_{\theta}$. Set $Y_{+, j}(t)=$ $t^{j} Y_{+}(j) / j$ !. If $j$ is a non-negativ integer, $p=0$ or 1 , then $\mathcal{W}_{\theta, j, p}$ is the space of all functions on the form $Y_{+, j}(t)\langle x\rangle^{p} U(x, t)$, where $U \in C^{\infty}\left(\overline{\mathbf{R}_{+}}\right) \otimes \mathcal{V}_{\theta}$.

Theorem 5. $\mathcal{L}_{\theta}$ is bijective from $\mathcal{W}_{\theta, j+1,1}$ to $\mathcal{W}_{\theta, j, 0}$ if $j \geq 0$.
By combining this result with some uniqueness statements obtained from Proposition 2 one can easily prove Theorem $1^{\prime}$ now by induction over $N$. We leave out these details and discuss instead the proof of the theorem above.

It is clear that $\mathcal{L}_{\theta}$ maps $\mathcal{W}_{\theta, j+1,1}$ into $\mathcal{W}_{\theta, j, 0}$. The injectivity of the map follows since one can show that $\mathcal{W}_{\theta, 0,1}$ is contained in $\mathcal{D}_{G_{\theta}}^{\prime}$. Hence convolution with $G_{\theta}$ gives us a left inverse.

In order to give the main ideas of the proof of the surjectivity of the map $\mathcal{L}_{\theta}$ in the theorem we consider the corresponding situation when $j=-1$ so that $Y_{+, j}(t)=\delta(t)$. This leads us to discuss the equation

$$
\begin{equation*}
\mathcal{L}_{\theta} V(x, t)=v(x) \delta(t) \tag{12}
\end{equation*}
$$

when $v \in \mathcal{V}$. We first construct an approximate solution. We set

$$
v_{j}(x)=\left(2^{j+1} j!\right)^{-1} \int_{0}^{\infty} \Delta_{x}^{j} v(x-s \theta) d s
$$

Then

$$
\begin{align*}
& 2\left\langle\theta, \partial_{x}\right\rangle v_{j}(x)=\Delta_{x} v_{j-1}(x), \quad j>0,  \tag{13}\\
& 2\left\langle\theta, \partial_{x}\right\rangle v_{0}(x)=v(x),
\end{align*}
$$

and $\langle x\rangle^{-1} v_{j}(x) \in \mathcal{V}_{\theta}$.
Choose $\zeta(t) \in C_{0}^{\infty}\left(\mathbf{R}^{n}\right)$ such that $\zeta(t)=1$ in a neighbourhood of the origin. If the sequence $1 \leq L_{j}$ grows sufficiently fast, then it is true that the series

$$
w(x, t)=\sum_{0}^{\infty} Y_{+, j}(t) \zeta\left(L_{j} t\right) v_{j}(x)
$$

converges in $C^{\infty}$ and defines an element in $\mathcal{W}_{\theta, 0,1}$. Moreover, it follows from (13) that

$$
r(x, t)=\mathcal{L}_{\theta} w(x, t)-v(x) \delta(t)
$$

is a smooth function of $t$ with values in $\mathcal{V}_{\theta}$. Moreover, it vanishes when $t \leq 0$. Since the dimension is odd, one has a simple explicit formula for $G_{\theta}$ which allows one to conclude that $w_{r}(x, t)=\left(G_{\theta} * r\right)(x, t)$ is a smooth function of $t$ with values in $\mathcal{V}_{\theta}$ after multiplication by $\langle x\rangle^{-1}$. Hence by subtraction $w_{r}$ from $w$ we have obtained a solution $V(x, t)$ of (12) such that $V \in \mathcal{W}_{\theta, 0,1}$.
Remark. The proof shows that $V_{\theta}(x, \varepsilon) \rightarrow v_{0}(x)$ as $\varepsilon \rightarrow 0$. Hence it follows from (13) that

$$
2\left\langle\theta, \partial_{x}\right\rangle V_{\theta}(x, \varepsilon) \rightarrow v(x) \quad \text { as } \varepsilon \rightarrow 0
$$

This phenomenon was discovered by R.G. Newton and called the miracle by him ([N1, N2]).

## Remarks about the case of exponentially decaying potentials.

We shall finally discuss a situation when the potential is exponentially decreasing. Let $a$ be a positive number and assume that one has the estimates

$$
\begin{equation*}
\left|v^{(\alpha)}(x)\right| \leq C_{\alpha} e^{-(2 a+\varepsilon)|x|} \tag{14}
\end{equation*}
$$

for every $\alpha$ and some positive $\varepsilon$. In this case it turns out that $V_{\theta}(x, t)$ will be exponentially decaying in the variable $t$ except for some contributions to $V_{\theta}$ that are due to bound states and resonances:

Theorem 6. Assume that $v$ satisfies (14). Then there is a finite set $Z \subset\{k \in \mathbf{C} ; \Im k \geq$ $-a\}$ so that

$$
\begin{equation*}
V_{\theta}(x, t)=\delta(t)+Y_{+}(t) a(x, \theta, t)+\sum_{z \in Z} \sum_{\mu \leq \mu(z)} t^{\mu} e^{-i t z} a_{z, \mu}(x, \theta) \tag{15}
\end{equation*}
$$

where for some constants $C_{\alpha, \beta}$ and $C_{\alpha}$

$$
\int_{0}^{\infty}\left|\partial_{x}^{\alpha} \partial_{t}^{\beta} a(x, \theta, t)\right| e^{a t} d t \leq C_{\alpha, \beta}\langle x\rangle^{\beta} e^{a(|x|-\langle x, \theta\rangle)}
$$

and

$$
\left|a_{z, \mu}^{(\alpha)}(x, \theta)\right| \leq C_{\alpha} e^{a(|x|-\langle x, \theta\rangle)}
$$

All estimates are uniform in $\theta$.
Remark. It is also possible to prove smoothness w.r.t. $\boldsymbol{\theta}$.

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