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# MIXED FINITE ELEMENT METHODS FOR QUASILINEAR SECOND ORDER ELLIPTIC PROBLEMS : THE p-VERSION (*) 

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Communicated by R Scott


#### Abstract

The p-version of the finite element method is analyzed for quasilinear second order elliptic problems in mixed weak form Approximation properties of the Raviart-Thomas projection are demonstrated and $L^{2}$-error bounds for the three relevant variables in the mixed method are derived

Résumé - Nous analysons la version-p de la méthode d'éléments finus mıxtes pour des problèmes quasilınéaires ellıptıques du second ordre en forme faıble mixte Nous démontrons des proprıétés d'approxımatıon de la projectıon de Raviart-Thomas et on dérıve des bornes de l'erreur dans $L^{2}(\Omega)$ pour les troıs varıables d'intérêt dans la méthode mixte


## I. INTRODUCTION

We consider here the numerical solution of the following boundary-value problem :

$$
\left\{\begin{align*}
\mathscr{D}(u)=-\underset{\sim}{\nabla} \cdot(a(u) \underset{\sim}{\nabla} u+\underset{\sim}{b}(u))+c(u) & =0 & & \text { in } \quad \Omega,  \tag{1.1}\\
u & =-g & & \text { on } \quad \partial \Omega,
\end{align*}\right.
$$

where $\Omega$ is a convex polygon with boundary $\partial \Omega, \nabla w$ denotes the gradient of the scalar function $w$ and $\underset{\sim}{\nabla} \cdot \underset{\sim}{v}$ and $\operatorname{div} \underset{\sim}{v}$ denote the divergence of the vector
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[^0]function $\underset{\sim}{v}$. We shall assume that for $r \geqslant 2$ and for each $g \in H^{r-1 / 2}(\partial \Omega)$ there exists a unique isolated solution $u \in H^{r}(\Omega)$ of (1.1) (that is, a solution not situated at a bifurcation point). Note that Sobolev's embedding theorem implies then that $u \in W^{r-1-\varepsilon, \infty}(\Omega), \varepsilon>0, \varepsilon \ll 1$, which will be needed throughout the paper.

We shall also assume that the coefficients $a: \bar{\Omega} \times \mathbb{R} \rightarrow \mathbb{R}$, $\underset{\sim}{b}: \bar{\Omega} \times \mathbb{R} \rightarrow \mathbb{R}^{2}$ and $c: \bar{\Omega} \times \mathbb{R} \rightarrow \mathbb{R}$ are twice continuously differentiable with bounded derivatives through second order, and that $a(\underset{\sim}{x}, q) \geqslant a_{1}>0$. The variable $\underset{\sim}{x}$ will be omitted as explicit argument of all functions, except when necessary to avoid ambiguity.

For $1 \leqslant s \leqslant \infty$ and $k$ any nonnegative integer, we let

$$
W^{k, s}(\Omega)=\left\{f \in L^{s}(\Omega): D^{\alpha} f \in L^{s}(\Omega),|\alpha| \leqslant k\right\}
$$

denote the Sobolev space endowed with its standard norm

$$
\begin{aligned}
\|f\|_{k, s, \Omega} & =\left(\sum_{|\alpha| \leqslant k}\left\|D^{\alpha} f\right\|_{L^{s}(\Omega)}^{s}\right)^{1 / s}, s \leqslant \infty \\
\|f\|_{k, \infty, \Omega} & =\max _{|\alpha| \leqslant k}\left\|D^{\alpha} f\right\|_{L^{\infty}(\Omega)}
\end{aligned}
$$

The subscript $\Omega$ in the norms will be omitted. Let $H^{k}(\Omega)=W^{k, 2}(\Omega)$ with norm $\|\cdot\|_{k}=\| \|_{k, 2}$. In particular, the notation $\|\cdot\|_{0}$ will mean $\|\cdot\|_{L^{2}(\Omega)}$ or $\|\cdot\|_{L^{2}(\Omega)^{2}}$ For $0 \leqslant r<\infty$ let $W^{r, s}(\Omega) . W^{r, s}(\partial \Omega), H^{r}(\Omega)$, and $H^{r}(\partial \Omega)$ denote the fractional order Sobolev spaces with norms $\|\cdot\|_{r, s, \Omega}$, $\|\cdot\|_{r, s, \partial \Omega},\|\cdot\|_{r, \Omega}$ and $\|\cdot\|_{r, \partial \Omega}$, respectively, defined by interpolation [7].

We shall denote by (., .) the Hilbert inner product in either $L^{2}(\Omega)$ or $L^{2}(\Omega)^{2}$ and by $\langle.,$.$\rangle the L^{2}$-inner product on the boundary of $\Omega$. The same notation will be used to indicate the dualities between $W^{r, s}(\Omega)$ and $W^{r, s}(\Omega)^{\prime}$ and $H^{s}(\partial \Omega)$ and $H^{-s}(\partial \Omega)$, respectively. Throughout the paper, $C$, $Q$, and $K$ will denote generic positive constants which need not have the same value in all their occurrences.

Let

$$
\underset{\sim}{V}=\underset{\sim}{H}(\operatorname{div} ; \Omega)=\left\{\underset{\sim}{v} \in L^{2}(\Omega)^{2}: \operatorname{div} \underset{\sim}{v} \in L^{2}(\Omega)\right\},
$$

normed by

$$
\|\underset{\sim}{v}\|_{\underline{H}(\operatorname{div}, \Omega)}^{2}=\|\underset{\sim}{v}\|_{0}^{2}+\|\operatorname{div} \underset{\sim}{v}\|_{0}^{2},
$$

and

$$
W=L^{2}(\Omega)
$$

The mixed finite element method we shall consider seeks simultaneous approximations of the solution of (1.1), $u$, and of the flux

$$
\begin{equation*}
\underset{\sim}{z}=-a(u) \underset{\sim}{\nabla} u-\underset{\sim}{b}(u) . \tag{1.2}
\end{equation*}
$$

The mixed weak formulation of (1.1) consists of finding $(z, u) \in \underset{\sim}{V} \times W$ such that

$$
\left\{\begin{array}{rlrl}
(\alpha(u) \underset{\sim}{z}, \underset{\sim}{v})- & (u, \operatorname{div} \underset{\sim}{v})+(\underset{\sim}{\beta}(u), \underset{\sim}{v})=\langle g, \underset{\sim}{v} \cdot \underset{\sim}{v}\rangle, & \underset{\sim}{v} \in \underset{\sim}{V},  \tag{1.3}\\
& (\operatorname{div} \underset{\sim}{z}, \underset{w}{ })+(c(u), w)=0, & & \underset{W}{w},
\end{array}\right.
$$

where we have set

$$
\begin{equation*}
\alpha(u)=1 / a(u), \quad \underset{\sim}{\beta}(u)=\alpha(u) \underset{\sim}{b}(u), \tag{1.4}
\end{equation*}
$$

and $\underset{\sim}{\boldsymbol{\nu}}$ is the outward unit normal vector on $\partial \Omega$. Our mixed finite element method is a discrete form of (1.3).

Let $\mathscr{G}$ be a decomposition of $\Omega$ by parallelograms which will be the «elements» $E$ and let $\mathscr{P}_{p, q}(E)=$ \{polynomials $f(x, y)$ on $E$, of degree $\leqslant p$ in $x$ and degree $\leqslant q$ in $y\}, \mathscr{Q}_{p}(E)=\{$ polynomials of degree $\leqslant p$ on $E\}$; next define, for each element $E$,

$$
\underline{V}^{p}(E)=\mathscr{P}_{p+1, p}(E) \times \mathscr{P}_{p, p+1}(E),
$$

and let

$$
\underline{V}^{p} \times W^{p} \subset \underset{\sim}{V} \times W
$$

be the Raviart-Thomas-Nedelec space of index $p \geqslant 0$ associated with this decomposition [3, 5], given by

$$
\left\{\begin{aligned}
{\underset{\sim}{V}}^{p} & =\left(\prod_{E \in \mathcal{C}} \underset{\sim}{V^{p}}(E)\right) \cap\left\{\underset{\sim}{f}: \Omega^{2} \rightarrow \mathbb{R} \mid \underset{\sim}{f} \cdot \nu_{E}\right. \\
& \left.=\underset{\sim}{f} \cdot \nu_{E^{\prime}} \text { on } E \cap E^{\prime}, E, E^{\prime} \in \mathfrak{C}\right\} \\
W^{p} & =\prod_{E \in \mathcal{C}} \mathscr{Q}_{p}(E)
\end{aligned}\right.
$$

where ${\underset{\sim}{\nu}}_{E}$ denotes the outward unit normal vector along $\partial E, E \in \mathcal{C}$. It is known [3,5] that $\operatorname{div}{\underset{V}{V}}^{p} \subset W^{p}$, a property we shall exploit later.

We seek $\left(\underset{\sim}{z}, u^{p}\right) \in{\underset{\sim}{V}}^{p} \times W^{p}$ so that

$$
\left\{\begin{array}{rlrl}
\left(\alpha\left(u^{p}\right) \underset{\sim}{p}, \underset{\sim}{v}\right)- & \left(u^{p}, \operatorname{div} \underset{\sim}{v}\right)+\left(\underset{\sim}{\beta}\left(u^{p}\right), \underset{\sim}{v}\right) & =\langle g, \underset{\sim}{v} \cdot \underset{\sim}{\nu}\rangle, \underset{\sim}{v} \in{\underset{\sim}{V}}^{p},  \tag{1.5}\\
& (\operatorname{div} \underset{\sim}{p}, w)+\left(c\left(u^{p}\right), w\right)=0, & & W^{p} .
\end{array}\right.
$$

Equations (1.5) define the $p$-version of the mixed finite element approximation for (1.3). This version is based on using a fixed mesh and increasing vol. $26, n^{\circ} 7,1992$
the degree of the finite elements (as opposed to the $h$-version that keeps the degree fixed and refines the mesh). The $p$-version has been analyzed for the linearized version of (1.1) in terms of the standard variational form in [1] and in terms of the mixed variational form in [6]. In this paper, we extend the results obtained in [6] for the linear problem to the quasilinear case. We also obtain an improved version of lemma 3.1 of [6] by reducing the regularity assumed there. We restrict our attention to the mixed method, the corresponding generalization for the standard finite element method is more straightforward.

Milner [3] described the $h$-version of this method for the same problem, demonstrated the unique solvability (for small $h$ ) of the nonlinear algebraic system (1.5) and derived error estimates in $L^{s}(\Omega), 2 \leqslant s \leqslant+\infty$, for the error in $u$, and in $\underset{\sim}{H}(\operatorname{div} ; \Omega)$ for the error in $z$. The assumption there was that the solution of (1.1) was in the space $H^{2+\varepsilon}(\Omega)$. In contrast, for the present paper we shall need an extra half derivative, that is, $u \in H^{5 / 2+\varepsilon}(\Omega)$.

We shall follow very closely the analysis of [3]. In order to do so we shall use the $L^{2}$-projection onto $W^{p}, P^{p}: L^{2} \rightarrow W^{p}$, given by

$$
\begin{equation*}
\left(P^{p} w-w, \chi\right)=0, \quad \chi \in W^{p}, \quad w \in W, \tag{1.6}
\end{equation*}
$$

for which the following approximation properties follow by repeating the arguments of [4] in two dimensions and using interpolation from the cases $s=2$ and $s=\infty$ :

$$
\begin{equation*}
\left\|\underline{P}^{p} w-w\right\|_{0, s} \leqslant Q p^{-m+3 / 2-3 / s}\|w\|_{m}, \quad s \geqslant 2, \quad 3 / 2-3 / s \leqslant m, \tag{1.7}
\end{equation*}
$$

if $w \in H^{m}(\Omega)$. We shall also use the Raviart-Thomas projection of $\underset{\sim}{V}$ onto ${\underset{\sim}{V}}^{p}, \pi^{p}: \underset{\sim}{V} \rightarrow{\underset{\sim}{V}}^{p}$, [5] for which we shall demonstrate in Section 2 the following approximation property :

$$
\begin{equation*}
\left\|\pi^{p} \underset{\sim}{v}-\underset{\sim}{v}\right\|_{0} \leqslant Q p^{1 / 2-r}\|\underset{\sim}{v}\|_{r}, \quad r>1 / 2, \quad \underset{\sim}{v} \in H^{r}(\Omega)^{2} \cap \underset{\sim}{V} . \tag{1.8}
\end{equation*}
$$

Our proof of (1.8) improves upon the one presented in [6], which imposed extra regularity on $\underset{\sim}{v}$ by requiring that $r>1$. In contrast, the condition $r>1 / 2$ is optimal (see remark 2.1). We also obtain estimates for the approximation properties of $\pi^{p}$ in the $W^{0, s}(\Omega)$-norm.

We shall find very useful the following inverse-type inequalities, the two dimensional form of the ones found in [4]:

$$
\begin{align*}
&\|\chi\|_{0, s} \leqslant Q p^{4 / r-4 / s}\|\chi\|_{0, r}, \quad 1 \leqslant r \leqslant s \leqslant \infty \\
& \chi \in L^{s}(\Omega) \cap W^{p}\left(\text { or } \chi \in L^{s}(\Omega)^{2} \cap{\underset{\sim}{V}}^{p}\right) . \tag{1.9}
\end{align*}
$$

The plan of the paper is as follows : in Section 2 we demonstrate (1.8), in Section 3 we prove that, for $p$ sufficiently large, (1.5) is uniquely solvable

[^1]and its solution $\left(\not \underset{\sim}{p}, u^{p}\right)$ converges to $(\underset{z}{z}, u)$ in $\underset{\sim}{V} \cap L^{2+\varepsilon}(\Omega)^{2} \times$ $L^{(2+4 \varepsilon) / \varepsilon}(\Omega)$ for any fixed $\varepsilon, 0 \ll \varepsilon \ll 1$, and in Section 4 we establish the rate of convergence of the approximation to the exact solution.

## II. THE APPROXIMATION PROPERTIES OF $\boldsymbol{\pi}^{\boldsymbol{p}}$

We recall that $\pi^{p} \underset{\sim}{v}$ is given locally (on every element $E$ ) by the following relations (2.1) and (2.2) (see [5]) :

$$
\begin{equation*}
\left\langle\left[\pi^{p} \underset{\sim}{v}-\underset{\sim}{v}\right] \cdot{\underset{\sim}{\nu}}_{E}, \varphi\right\rangle_{S_{t}}=0, \quad \varphi \in \mathscr{P}_{p}, \tag{2.1}
\end{equation*}
$$

where $\langle., .\rangle_{S}, 1 \leqslant i \leqslant 4$, denotes the line integral along each side $S_{\imath}$ of the element $E$ and $\mathscr{P}_{p}$ is the set of all polynomials in one variable of degree less than or equal to $p$,

$$
\begin{equation*}
\left(\pi^{p} \underset{\sim}{v}-\underset{\sim}{v}, \underset{\sim}{\psi}\right)_{E}=0, \quad \underset{\sim}{\psi} \in{\underset{\sim}{V}}^{p}(E), \tag{2.2}
\end{equation*}
$$

where (.,. $)_{E}$ denotes the standard $L^{2}(E)$-inner product.
Now, let $R=[-1,1] \times[-1,1]$ and let $\left\{P_{l}\right\}_{l \geqslant 0}$ denote the $L^{2}([-1,1])$-complete orthogonal Legendre polynomials. For any $\underset{\sim}{v} \in \underset{\sim}{H}(\operatorname{div} ; R)$, let

$$
\begin{equation*}
\underset{\sim}{v}(x, y)=\left[\sum_{i=0}^{\infty} \sum_{j=0}^{\infty} a_{i, j} P_{i}(x) P_{j}(y), \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} b_{i, j} P_{i}(x) P_{J}(y)\right], \tag{2.3}
\end{equation*}
$$

and let

$$
\begin{equation*}
\pi^{p} \underset{\sim}{\boldsymbol{v}}(x, y)=\left[\sum_{i=0}^{p+1} \sum_{j=0}^{p} \tilde{a}_{i, J} P_{l}(x) P_{J}(y), \sum_{i=0}^{p} \sum_{j=0}^{p+1} \tilde{b}_{l, j} P_{l}(x) P_{J}(y)\right] . \tag{2.4}
\end{equation*}
$$

It follows from (2.2)-(2.4) that

$$
\left\{\begin{array}{lll}
a_{i, j}=\tilde{a}_{l, j}, & 0 \leqslant i \leqslant p-1, & 0 \leqslant j \leqslant p  \tag{2.5}\\
b_{k l}=\tilde{b}_{k l}, & 0 \leqslant k \leqslant p, & 0 \leqslant l \leqslant p-1
\end{array}\right.
$$

Next, we see that (2.1), (2.3)-(2.5) imply that

$$
\begin{cases}\sum_{l=p}^{p+1} \tilde{a}_{l, J} P_{l}( \pm 1)=\sum_{i=p}^{\infty} a_{i, J} P_{l}( \pm 1), & 0 \leqslant j \leqslant p  \tag{2.6}\\ \sum_{J=p}^{p+1} \tilde{b}_{l, J} P_{J}( \pm 1)=\sum_{J=p}^{\infty} b_{l, J} P_{l}( \pm 1), & 0 \leqslant i \leqslant p\end{cases}
$$

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Since $P_{t}(-1)=(-1)^{t}$ and $P_{t}(1)=1$, (2.6) implies that

Proposition 2.1 : Let $\underset{\sim}{v} \in \underset{\sim}{V}$ and let $\pi^{p} \underset{\sim}{v}$ be its Raviart-Thomas projection in ${\underset{\sim}{V}}^{p}$ given by (2.1)-(2.2). Then, if $\underset{\sim}{v} \in H^{r}(\Omega)^{2}$, we have

$$
\left\|\underset{\sim}{v}-\pi^{p} \underset{\sim}{v}\right\|_{0} \leqslant Q p^{1 / 2-r}\|\underset{\sim}{v}\|_{r}, \quad r>1 / 2,
$$

where $Q>0$ is a constant independent of $p$ and $\underset{\sim}{v}$ but depending on $r$.
Proof: We first assume that $\Omega=R$ and that the decomposition consists of just one element. Then, $\underset{\sim}{v} \in \underset{\sim}{V}$ and $\pi^{p} \underset{\sim}{v} \in{\underset{\sim}{V}}^{p}$ can be given, respectively, by (2.3) and (2.4).

The following relation is a trivial consequence of well known properties of the Legendre polynomials,

$$
\begin{aligned}
& \left\|\underset{\sim}{v}-\pi^{p} \underset{\sim}{v}\right\|_{0}^{2}=\sum_{i=p}^{p+1} \sum_{j=0}^{p} \frac{4\left(a_{i, j}-\tilde{a}_{i, j}\right)^{2}}{(2 i+1)(2 j+1)}+\sum_{i=0}^{p} \sum_{j=p}^{p+1} \frac{4\left(b_{i, j}-\tilde{b}_{i, j}\right)^{2}}{(2 i+1)(2 j+1)}+ \\
& +\sum_{i=0}^{p+1} \sum_{j=p+i}^{\infty} \frac{4 a_{i, j}^{2}}{(2 i+1)(2 j+1)}+\sum_{i-p+1}^{\infty} \sum_{j-0}^{p+1} \frac{4 b_{i, j}^{2}}{(2 i+1)(2 j+1)} \\
& +\sum_{i=p+2}^{\infty} \sum_{j=0}^{p} \frac{4 a_{i, j}^{2}}{(2 i+1)(2 j+1)}+\sum_{i=0}^{p} \sum_{j=p+2}^{\infty} \frac{4 b_{i, j}^{2}}{(2 i+1)(2 j+1)} \\
& +\sum_{i=p+2}^{\infty} \sum_{j=p+1}^{\infty} \frac{4 a_{i, j}^{2}}{(2 i+1)(2 j+1)}+\sum_{i=p+1}^{\infty} \sum_{j=p+2}^{\infty} \frac{4 b_{i, j}^{2}}{(2 i+1)(2 j+1)} \\
& =\mathrm{I}+\mathrm{II}+\cdots+\text { VIII } \text {. }
\end{aligned}
$$

Note that III-VIII can be bounded as follows :

$$
\mathrm{III}+\mathrm{V}+\mathrm{VII} \leqslant Q p^{-2 r} \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} \frac{a_{i, j}^{2}\left(1+i^{2}+j^{2}\right)^{r}}{(2 i+1)(2 j+1)}, \quad r \geqslant 0
$$

while

$$
\mathrm{IV}+\mathrm{VI}+\mathrm{VIII} \leqslant Q p^{-2 r} \sum_{i=0}^{\infty} \sum_{j=0}^{\infty} \frac{b_{i, j}^{2}\left(1+i^{2}+j^{2}\right)^{r}}{(2 i+1)(2 j+1)}, \quad r \geqslant 0
$$

which implies that (see [6])

$$
\begin{equation*}
\mathrm{III}+\cdots+\mathrm{VIII} \leqslant Q p^{-2 r}\|\underset{\sim}{v}\|_{r}^{2}, \quad r \geqslant 0 \tag{2.9}
\end{equation*}
$$

On the other hand, it follows from (2.7) that

$$
\begin{align*}
\mathrm{I}=\frac{4}{2 p+1} \sum_{J=0}^{p} & \frac{\left(a_{p, J}-\sum_{j=0}^{\infty} a_{2 \imath+p, J}\right)^{2}}{2 j+1}+\frac{4}{2 p+3} \sum_{J=0}^{p} \frac{\left(a_{p+1, J}-\sum_{i=0}^{\infty} a_{2 \imath+p+1, J}\right)^{2}}{2 j+1} \\
& =\frac{4}{2 p+1} \sum_{J=0}^{p} \frac{\left(\sum_{i=1}^{\infty} a_{2 \imath+p, J}\right)^{2}}{2 j+1}+\frac{4}{2 p+3} \sum_{J=0}^{p} \frac{\left(\sum_{\imath=1}^{\infty} a_{2 \imath+p+1, J}\right)^{2}}{2 j+1} \tag{2.10}
\end{align*}
$$

and

$$
\begin{align*}
\mathrm{II}= & \frac{4}{2 p+1} \sum_{i=0}^{p} \frac{\left(b_{l p}-\sum_{J=0}^{\infty} b_{l, 2 J+p}\right)^{2}}{2 i+1} \\
& +\frac{4}{2 p+3} \sum_{i=0}^{p} \frac{\left(b_{i, p+1}-\sum_{J=0}^{\infty} b_{l, 2 J+p+1}\right)^{2}}{2 i+1} \\
= & \frac{4}{2 p+1} \sum_{i=0}^{p} \frac{\left(\sum_{J=1}^{\infty} b_{l, 2 J+p}\right)^{2}}{2 i+1}+\frac{4}{2 p+3} \sum_{i=0}^{\infty} \frac{\left(\sum_{J=1}^{\infty} b_{l, 2}+p+1\right)^{2}}{2 i+1} . \tag{2.11}
\end{align*}
$$

Next observe that bounding the series $\sum_{k=p+1}^{\infty}\left(c+k^{2}\right)^{-s}(1+2 k)$ using the integral method for $\int_{p}^{\infty}\left(c+t^{2}\right)^{-s}(1+2 t) d t \leqslant \frac{K}{s-1} p^{2-2 s} \quad(s>1)$, and using the Cauchy-Schwarz inequality, we see that, for $s$ bounded away from 1 ,

$$
\begin{align*}
\left(\sum_{i=1}^{\infty}\right. & \left.a_{2 i+p, J}\right)^{2} \leqslant \\
\leqslant & \sum_{i=1}^{\infty} \frac{a_{2 i+p, J}^{2}}{4 i+2 p+1}\left[1+(2 i+p)^{2}+j^{2}\right]^{s} \sum_{i=1}^{\infty}\left[1+(2 i+p)^{2}+j^{2}\right]^{-s} \times \\
& \times(1+4 i+2 p) \\
& \leqslant \sum_{i=0}^{\infty} \frac{a_{i, j}^{2}\left(1+i^{2}+j^{2}\right)^{s}}{2 i+1} \sum_{t=p+2}^{\infty}\left(1+i^{2}+j^{2}\right)^{-s}(1+2 i) \\
& \leqslant Q p^{2-2 s} \sum_{i=0}^{\infty} \frac{a_{i, j}^{2}\left(1+i^{2}+j^{2}\right)^{s}}{2 i+1} \tag{2.12}
\end{align*}
$$

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with exactly the same final bound holding for $\left(\sum_{t=1}^{\infty} a_{2 t+p+1, J}\right)^{2}$. It follows from (2.10) and (2.12) that, for $s$ bounded away from 1 ,

$$
\begin{equation*}
\mathrm{I} \leqslant Q p^{1-2 s} \sum_{J=0}^{\infty} \sum_{t=0}^{\infty} \frac{a_{i, J}^{2}\left(1+i^{2}+j^{2}\right)^{s}}{(2 i+1)(2 j+1)} \tag{2.13}
\end{equation*}
$$

In an entirely analogous way (replacing $a_{i, j}$ by $b_{i, j}$ and reversing the roles of $i$ and $J$ ) we deduce from (2.11) that, for $s$ bounded away from 1,

$$
\begin{equation*}
\mathrm{II} \leqslant Q p^{1-2 s} \sum_{i=0}^{\infty} \sum_{j=1}^{\infty} \frac{b_{t, j}^{2}\left(1+i^{2}+j^{2}\right)^{s}}{(2 i+1)(2 j+1)} \tag{2.14}
\end{equation*}
$$

Combining (2.13) and (2.14) results, for $s$ bounded away from 1 , in

$$
\begin{equation*}
\mathrm{I}+\mathrm{II} \leqslant Q p^{1-2 s}\|\underset{\sim}{v}\|_{s}^{2} . \tag{2.15}
\end{equation*}
$$

Next note that

$$
\begin{align*}
& \left.\underset{\sim}{v} \cdot \underset{\sim}{\nu}\right|_{\partial R}= \\
&  \tag{2.16}\\
& v_{1}( \pm 1, y)=\sum_{i=0}^{\infty} \sum_{j=0}^{\infty} a_{i, j} P_{t}( \pm 1) P_{J}(y),-1 \leqslant y \leqslant 1, \\
& v_{2}(x, \pm 1)=\sum_{i=0}^{\infty} \sum_{j=0}^{\infty} \dot{b}_{i, j} P_{i}(y) P_{J}( \pm 1), \quad-1 \leqslant x \leqslant 1 .
\end{align*}
$$

The trace theorem (Sobolev's embedding theorem) implies that $v_{1}, v_{2} \in$ $L^{2}(\partial R)$ for $s>1 / 2$. Consequently, since $P_{l}(1)=1$ and $P_{l}(-1)=(-1)^{l}$, we see from (2.16) that

$$
\begin{align*}
& \left\|v_{1}( \pm 1, .)\right\|_{0, \partial \Omega}^{2}=2 \sum_{j=0}^{\infty} \frac{\left[\sum_{t=0}^{\infty}( \pm 1)^{t} a_{\imath, J}\right]^{2}}{2 j+1}<\infty, \\
& \left\|v_{2}(., \pm 1)\right\|_{0, \partial \Omega}^{2}=2 \sum_{j=0}^{\infty} \frac{\left[\sum_{j=0}^{\infty}( \pm 1)^{\prime} b_{l, J}\right]^{2}}{2 i+1}<\infty . \tag{2.17}
\end{align*}
$$

Let now $\underset{\sim}{v} \in H^{1 / 2+\varepsilon}(\Omega)^{2}$. We shall prove that

$$
\begin{equation*}
\left\|\pi^{p} \underset{\sim}{v}-\underset{\sim}{v}\right\|_{0} \leqslant Q p^{-\varepsilon}\|\underset{\sim}{v}\|_{1 / 2+\varepsilon} \tag{2.18}
\end{equation*}
$$

In view of (2.8) and (2.9) it is sufficient to prove that I , $\mathrm{II}<Q p^{-2 \varepsilon}\|\underset{\sim}{v}\|_{1 / 2+\varepsilon}^{2}$. It follows from (2.10) that

$$
\begin{align*}
& \mathrm{I} \leqslant 4 p^{-1} \sum_{j=0}^{p}(2 j+1)^{-1}\left[\left(\sum_{t=1}^{\infty} a_{2 i+p, j}\right)^{2}+\left(\sum_{\imath=1}^{\infty} a_{2 l+p+1, \jmath}\right)^{2}\right] \\
& =2 p^{-1} \sum_{j=0}^{p}(2 j+1)^{-1}\left[\left(\sum_{i=p+2}^{\infty} a_{\imath, J}\right)^{2}+\left(\sum_{t=p+2}^{\infty}(-1)^{i} a_{i, j}\right)^{2}\right] \\
& =2 p^{-1} \sum_{j=0}^{p}(2 j+1)^{-1}\left[\left(\sum_{i-0}^{\infty} a_{i, j}-\sum_{i=0}^{p+1} a_{i, j}\right)^{2}+\right. \\
& \left.+\left(\sum_{i=0}^{\infty}(-1)^{t} a_{t, \jmath}-\sum_{t=0}^{p+1}(-1)^{t} a_{t, j}\right)^{2}\right] \\
& \leqslant 4 p^{-1} \sum_{j=0}^{p}(2 j+1)^{-1}\left[\left(\sum_{t=0}^{\infty} a_{l, j}\right)^{2}+\left(\sum_{i=0}^{\infty}(-1)^{l} a_{i, j}\right)^{2}+\left(\sum_{t=0}^{p+1} a_{l, j}\right)^{2}+\right. \\
& \left.+\left(\sum_{i=0}^{p+1}(-1)^{l} a_{i, j}\right)^{2}\right] \\
& \leqslant 4 p^{-1}\left(\left\|v_{1}(1, .)\right\|_{0, \partial \Omega}^{2}+\left\|v_{1}(-1, .)\right\|_{0, \partial \Omega}^{2}\right)+ \\
& +4 p^{-1} \sum_{j=0}^{p}(2 j+1)^{-1}\left[\left(\sum_{l=0}^{p+1} a_{l, J}\right)^{2}+\left(\sum_{t=0}^{p+1}(-1)^{l} a_{l, J}\right)^{2}\right] . \tag{2.19}
\end{align*}
$$

Note that the next to last term on the right hand side of (2.19) can be bounded, using the integral method for

$$
\int_{0}^{p+1}(2 i+1)\left(1+i^{2}+j^{2}\right)^{-1 / 2-\varepsilon}=O\left(p^{1-2 \varepsilon}\right)
$$

as $p \rightarrow \infty$, as follows :

$$
\begin{align*}
\sum_{J=0}^{p}(2 j+1)^{-1}\left(\sum_{t=0}^{p+1} a_{t, j}\right)^{2} \leqslant & \sum_{j=0}^{p}(2 j+1)^{-1} \sum_{t=0}^{p+1} \frac{a_{t, j}^{2}\left(1+i^{2}+j^{2}\right)^{1 / 2+\varepsilon}}{2 i+1} \\
& \times \sum_{t=0}^{p+1}(2 i+1)\left(1+i^{2}+j^{2}\right)^{-1 / 2-\varepsilon} \\
\leqslant & Q\|\underset{\sim}{v}\|_{1 / 2+\varepsilon}^{2} p^{1-2 \varepsilon} \tag{2.20}
\end{align*}
$$

with an identical bound holding for the last term of (2.19). Combining (2.19) and (2.20) and using Sobolev's embedding theorem yields

$$
\begin{equation*}
\mathrm{I} \leqslant Q p^{-2 \varepsilon}\|\underset{\sim}{v}\|_{1 / 2+\varepsilon}^{2} \tag{2.21}
\end{equation*}
$$

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In an entirely analogous fashion we can see that

$$
\mathrm{II} \leqslant Q p^{-2 \varepsilon}\|\underset{\sim}{v}\|_{1 / 2+\varepsilon}^{2}
$$

which together with (2.21) yields (2.18). Using interpolation [7], it follows from (2.15) and (2.18) that, for $s$ bounded away from $1 / 2$,

$$
\mathrm{I}+\mathrm{II} \leqslant Q p^{1-2 s}\|\underset{\sim}{v}\|_{s}^{2}
$$

which together with (2.8) and (2.9) concludes the proof for the case $\Omega=R$. For the case when $\Omega$ is a disjoint union of parallelograms the result follows on each element by using affine mappings onto $R$. The proposition then follows by summing over all the elements (see [6] for details).

Remark 2.1: This result differs from the one known for the $h$-version of the finite element method [3, (1.5)],

$$
\begin{equation*}
\left\|\underset{\sim}{\boldsymbol{v}}-\pi^{h} \underset{\sim}{v}\right\|_{0} \leqslant Q h^{r}\|\underset{\sim}{\boldsymbol{v}}\|_{r}, \quad r>1 / 2 . \tag{2.22}
\end{equation*}
$$

The constraint $r>1 / 2$ (or $r \geqslant 1 / 2+\varepsilon$ ) stems from the fact that, according to the trace theorem, this is the minimal requirement to ensure that $\underset{\sim}{v}$ has a trace on the boundary which is an $L^{2}$-function (not just a distribution). In [6] the corresponding result required an additional half derivative on $\underset{\sim}{v}(r>1)$. In contrast, proposition 2.1 assumes the minimum regularity necessary. It is possible, however, that the bound still holds with the exponent of $p$ replaced by $-r$, as suggested by (2.22).

COROLLARY $2.1:$ For $s \geqslant 2, r>\max \{1 / 2 ; 3 / 2-3 / s\}$,

$$
\left\|\underset{\sim}{v}-\pi^{p} \underset{\sim}{v}\right\|_{0, s} \leqslant Q p^{5 / 2-r-4 / s}\|\underset{\sim}{v}\|_{r} .
$$

Proof : Let ${\underset{\sim}{\sim}}^{p} \underset{\sim}{v}$ be the $L^{2}$-projection $P^{p} \times P^{p}: \underset{\sim}{V} \rightarrow{\underset{\sim}{V}}^{p}$. Then the following analogue of (1.7) holds :

$$
\begin{equation*}
\left\|{\underset{\sim}{P}}^{p} \underset{\sim}{v}-\underset{\sim}{v}\right\|_{0, s} \leqslant Q p^{-r+3 / 2-3 / s}\|\underset{\sim}{v}\|_{r}, \quad s \geqslant 2, \quad 3 / 2-3 / s \leqslant r . \tag{2.23}
\end{equation*}
$$

Also,

$$
\begin{equation*}
\left\|\pi^{p} \underset{\sim}{v}-\underset{\sim}{v}\right\|_{0, s} \leqslant\left\|{\underset{\sim}{P}}^{p} \underset{\sim}{v}-\underset{\sim}{v}\right\|_{0, s}+\left\|\pi^{p} \underset{\sim}{v}-{\underset{\sim}{P}}^{p} \underset{\sim}{v}\right\|_{0, s} . \tag{2.24}
\end{equation*}
$$

The second term in this expression may be bounded using the inverse inequality (1.9) as follows :

$$
\begin{align*}
\| \pi^{p} \underset{\sim}{v}-{\underset{\sim}{P}}^{p} & \underset{\sim}{v} \|_{0, s}
\end{align*} \leqslant Q p^{2-4 / s}\left\|\pi^{p} \underset{\sim}{v}-{\underset{\sim}{P}}^{p} \underset{\sim}{v}\right\|_{0} .
$$

Combining (2.23)-(2.25) and using proposition 2.1, we obtain the corollary.

## III. SOLVABILITY OF THE DISCRETE PROBLEM

Following [3] we introduce, for $\rho \in W^{p}$, the notation
$\alpha(\rho)-\alpha(u)=-\tilde{\alpha}_{u}(\rho)(u-\rho)=-\alpha_{u}(u)(u-\rho)+\tilde{\alpha}_{u u}(\rho)(u-\rho)^{2}$,
where

$$
\tilde{\alpha}_{u}(\rho)=\int_{0}^{1} \alpha_{u}(\rho+t[u-\rho]) d t
$$

and

$$
\tilde{\alpha}_{u u}(\rho)=\int_{0}^{1}(1-t) \alpha_{u u}(u+t[\rho-u]) d t
$$

are bounded functions in $\bar{\Omega}$. Similarly, we write
$\underset{\sim}{\beta}(\rho)-\underset{\sim}{\beta}(u)=-{\underset{\sim}{\underset{\sim}{\beta}}}_{u}(\rho)(u-\rho)=-{\underset{\sim}{\beta}}_{u}(u)(u-\rho)+{\underset{\sim}{\underset{\beta}{\beta}}}_{u u}(\rho)(u-\rho)^{2}$,
and
$c(\rho)-c(u)=-\tilde{c}_{u}(\rho)(u-\rho)=-c_{u}(u)(u-\rho)+\tilde{c}_{u u}(\rho)(u-\rho)^{2}$,
where $\tilde{\sim}_{\sim}^{\beta}(\rho), \tilde{\sim}_{\sim u}(\rho), \tilde{c}_{u}(\rho)$, and $\tilde{c}_{u u}(\rho)$ are bounded functions in $\bar{\Omega}$. Also, let

$$
\begin{equation*}
\underset{\sim}{\Gamma}=\alpha_{u}(u) \underset{\sim}{z}+{\underset{\sim}{\beta}}_{u}(u), \quad \gamma=c_{u}(u) \tag{3.4}
\end{equation*}
$$

With the notation of (3.1)-(3.4), the following error equations follow from (1.3) and (1.5), [3] :

$$
\begin{align*}
\left(\alpha(u)\left[\pi^{p} \underset{\sim}{z}-\underset{\sim}{z}\right], \underset{\sim}{v}\right)-\left(\operatorname{div} \underset{\sim}{v}, P^{p} u-u^{p}\right) & +\left(\left[P^{p} u-u^{p}\right] \underset{\sim}{\Gamma}, \underset{\sim}{v}\right)= \\
& =\left(\underset{\sim}{q}\left(u^{p}, \underset{\sim}{p}\right), \underset{\sim}{v}\right), \quad \underset{\sim}{v} \in{\underset{\sim}{V}}^{p}, \\
\left(\operatorname{div}\left[\pi^{p} \underset{\sim}{z}-\underset{\sim}{p}\right], w\right)+\left(\gamma\left[P^{p} u-u^{p}\right], w\right) & =\left(\eta\left(u^{p}\right), w\right), \quad w \in W^{p}, \tag{3.5}
\end{align*}
$$

where

$$
\begin{align*}
\underset{\sim}{q}\left(u^{p}, \underset{\sim}{p}\right) & =\alpha(u)\left[\pi^{p} z-\underset{z}{z}\right]+\left[P^{p} u-u\right] \underset{\sim}{\Gamma}+ \\
& +\left(u-u^{p}\right)^{2}\left[\tilde{\alpha}_{u u}\left(u^{p}\right) z+{\underset{\sim}{\tilde{\beta}}}_{u u}\left(u^{p}\right)\right]+\tilde{\alpha}_{u}\left(u^{p}\right)\left(u-u^{p}\right)\left(\underset{\sim}{z}-\not z^{p}\right), \tag{3.6}
\end{align*}
$$

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and

$$
\begin{equation*}
\eta\left(u^{p}\right)=\gamma\left[P^{p} u-u\right]+\tilde{c}_{u u}\left(u^{p}\right)\left(u-u^{p}\right)^{2} . \tag{3.7}
\end{equation*}
$$

Just as in [3], we let

$$
\Phi:{\underset{\sim}{V}}^{p} \times W^{p} \rightarrow \underline{V}^{p} \times W^{p}
$$

be given by $\Phi((\underset{\sim}{\mu}, \rho))=(\underset{\sim}{\lambda}, \kappa),(\underset{\sim}{\lambda}, \kappa)$ being the (unique) solution of the system

$$
\begin{aligned}
\left(\alpha(u)\left[\pi^{p} z-\underset{\sim}{\lambda}\right], \underset{\sim}{v}\right)-\left(\operatorname{div} \underset{\sim}{v}, P^{p} u-\kappa\right) & +\left(\left[P^{p} u-\kappa\right] \underset{\sim}{\Gamma}, \underset{\sim}{v}\right)= \\
& =(\underset{\sim}{q}(\rho, \underset{\sim}{\mu}), \underset{\sim}{v}), \underset{\sim}{v} \in{\underset{V}{V}}^{p},
\end{aligned}
$$

$$
\begin{equation*}
\left(\operatorname{div}\left[\pi^{p} \underset{\sim}{z}-\underset{\sim}{\lambda}\right], w\right)+\left(\gamma\left[P^{p} u-\kappa\right], w\right)=(\eta(\rho), w), \quad w \in W^{p} \tag{3.8}
\end{equation*}
$$

where $\underset{\sim}{q}(\rho, \mu)$ and $\eta(\rho)$ are given by (3.6) and (3.7), respectively, replacing $u^{p}$ by $\rho$ and $\underset{\sim}{z}$ by $\underset{\sim}{\mu}$. The unique solvability of this (linear) system follows, for $p$ sufficiently large, from [2], since the left hand side of (3.8) corresponds to the mixed method for the operator $M: H^{2}(\Omega) \cap H_{0}^{1}(\Omega) \rightarrow$ $L^{2}(\Omega)$ given by

$$
M w=-\underset{\sim}{\nabla} \cdot(a(u) \underset{\sim}{\nabla} w+a(u) w \underset{\sim}{\Gamma})+\gamma \omega,
$$

which has a bounded inverse. In fact, note that (1.2), (1.4) and (3.4) give

$$
\begin{aligned}
M w= & -\underset{\sim}{\nabla} \cdot\left[a(u) \underset{\sim}{\nabla} w+a(u) w\left(\alpha_{u}(u) \underset{\sim}{z}+{\underset{\sim}{\beta}}_{u}(u)\right)\right]+c_{u}(u) w \\
= & -\underset{\sim}{\nabla} \cdot\left[a(u) \underset{\sim}{\nabla} w+a(u) w\left[-\frac{a_{u}(u)}{a^{2}(u)}(-a(u) \nabla u)+\alpha(u){\underset{\sim}{u}}_{u}(u)\right]\right]+ \\
& +c_{u}(u) w \\
= & -\underset{\sim}{\nabla} \cdot\left[a(u) \underset{\sim}{\nabla} w+\left(a_{u}(u) \nabla u+{\underset{\sim}{b}}_{u}(u)\right) w\right]+c_{u}(u) w,
\end{aligned}
$$

which shows that $M$ is the linearization of the operator $\mathscr{D}$ in (1.1) about the function $u$, and, thus, it has a bounded inverse since we have assumed that (1.1) admits unique isolated solutions.

The solvability of (1.5) is now equivalent to showing that $\Phi$ has a fixed point. This will follow from the Brouwer fixed point theorem if we show that $\Phi$ maps a ball of $V^{p} \times W^{p}$ into itself. We shall need the following technical result, a $p$-version of lemma 2.1 of [3]. Let $\varepsilon>0$ be fixed for the rest of the paper, $\varepsilon \ll 1$.
 If $\tau \in W^{p}$ satisfies

$$
\left\{\begin{aligned}
(\alpha(u) \underset{\sim}{\omega}, \underset{\sim}{v})- & (\operatorname{div} \underset{\sim}{v}, \tau)+(\tau \underset{\sim}{\Gamma}, \underset{\sim}{v})= \\
& (\operatorname{div} \underset{\sim}{q}, \underset{\sim}{v}), \underset{\sim}{v} \in{\underset{\sim}{V}}^{p}, \\
(\gamma \tau, \omega)= & (\eta, w), w \in W^{p},
\end{aligned}\right.
$$

then, there exists a constant $C=C(\theta, u, \alpha, \underset{\sim}{\Gamma}, \gamma, \Omega, \varepsilon)$ such that, for $p$ sufficiently large, depending upon $\varepsilon$,

$$
\|\tau\|_{0, \theta} \leqslant C\left[p^{1 / 2-2 / \theta}\|\underset{\sim}{\omega}\|_{0}+p^{-1-2 / \theta}\|\operatorname{div} \underset{\sim}{\omega}\|_{0}+\|\underset{\sim}{q}\|_{0}+\|\eta\|_{0}\right]
$$

Proof: We follow the proof of lemma 2.1 of [3]. Let $\theta^{\prime}=\theta /(\theta-1)$ be the conjugate exponent of $\theta$. For $\psi \in L^{\theta^{\prime}}(\Omega)$ let $\phi \in W^{2, \theta^{\prime}}(\Omega)$ be the (unique) solution of $M^{*} \phi=\psi$ in $\Omega, \psi=0$ on $\partial \Omega$, where $M^{*}$ is the formal adjoint of $M$. It follows that $\|\phi\|_{2, \theta^{\prime}} \leqslant Q\|\psi\|_{0, \theta^{\prime}}$. We then have [3],

$$
\begin{align*}
& (\tau, \psi)=(\underset{\sim}{q}, a(u) \underset{\sim}{\nabla} \phi)+\left(\underset{\sim}{q}, \pi^{P} a(u) \underset{\sim}{\nabla} \phi-a(u) \underset{\sim}{\nabla} \phi\right)+ \\
& \quad+\left(\operatorname{div} \underset{\sim}{\underset{\omega}{*}}+\gamma \tau, \phi-P^{p} \phi\right) \\
& \quad+\left(\alpha(u) \underset{\sim}{\omega}+\tau \underset{\sim}{\underset{\sim}{r}}, a(u) \underset{\sim}{\nabla} \phi-\pi^{p} a(u) \underset{\sim}{\nabla} \phi\right)+(\eta, \phi)+\left(\eta, P^{p} \phi-\phi\right) . \tag{3.9}
\end{align*}
$$

Note that Sobolev's embedding theorem implies that

$$
\begin{equation*}
(\underset{\sim}{q}, a(u) \underset{\sim}{\nabla} \phi) \leqslant C\|\underset{\sim}{q}\|_{0}\|\phi\|_{1} \leqslant C\|\underset{\sim}{q}\|_{0}\|\phi\|_{2, \theta^{\prime}} \tag{3.10}
\end{equation*}
$$

Next, (1.8) and Sobolev's embedding theorem imply that

$$
\begin{align*}
\left(\underset{\sim}{q}-\alpha(u) \underset{\sim}{\omega}, \pi^{p} a(u) \underset{\sim}{\nabla} \phi\right. & -a(u) \underset{\sim}{\nabla} \phi) \leqslant \\
& \leqslant C\left(\|\underset{\sim}{q}\|_{0}+\|\underset{\sim}{\omega}\|_{0}\right) p^{1 / 2-2 / \theta}\|\underset{\sim}{\nabla} \phi\|_{2 / \theta} \\
& \leqslant C\left(\|\underset{\sim}{q}\|_{0}+\|\underset{\sim}{\omega}\|_{0}\right) p^{1 / 2-2 / \theta}\|\phi\|_{2, \theta^{\prime}}, \tag{3.11}
\end{align*}
$$

and that

$$
\begin{align*}
\left(\tau \underset{\sim}{\Gamma}, a(u) \underset{\sim}{\nabla} \phi-\pi^{p} a(u) \underset{\sim}{\nabla} \phi\right) & \leqslant C\|\tau\|_{0, \theta}\left\|a(u) \underset{\sim}{\nabla} \phi-\pi^{p} a(u) \underset{\sim}{\nabla} \phi\right\|_{0} \\
& \leqslant C\|\tau\|_{0, \theta} p^{-\varepsilon / 8}\|\phi\|_{2, \theta^{\prime}} \tag{3.12}
\end{align*}
$$

On the other hand, (1.7) and Sobolev's embedding theorem lead to

$$
\begin{align*}
\left(\operatorname{div} \underset{\sim}{\omega}, \phi-P^{p} \phi\right) & \leqslant K\|\operatorname{div} \underset{\sim}{\omega}\|_{0} p^{-1-2 / \theta}\|\phi\|_{2, \theta^{\prime}},  \tag{3.13}\\
\left(\gamma \tau, \phi-P^{p} \phi\right) & \leqslant K\|\tau\|_{0, \theta} p^{-1-2 / \theta}\|\phi\|_{2, \theta^{\prime}}, \tag{3.14}
\end{align*}
$$

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and

$$
\begin{equation*}
(\eta, \phi)+\left(\eta, P^{p} \phi-\phi\right) \leqslant K\|\eta\|_{0}\|\phi\|_{0} \leqslant K\|\eta\|_{0}\|\phi\|_{2, \theta^{\prime}} \tag{3.15}
\end{equation*}
$$

Collecting (3.9)-(3.15) we see that

$$
\begin{aligned}
&(\tau, \psi) \leqslant K\|\psi\|_{0, \theta^{\prime}}\left[p^{1 / 2-2 / \theta}\|\stackrel{\omega}{\sim}\|_{0}+p^{-1-2 / \theta}\|\operatorname{div} \underset{\sim}{\omega}\|_{0}+\right. \\
&\left.+p^{-\varepsilon / 8}\|\tau\|_{0, \theta}+\|\underset{\sim}{q}\|_{0}+\|\eta\|_{0}\right]
\end{aligned}
$$

which, for $p$ sufficiently large, yields the desired estimate.
Now let ${\underset{\sim}{\mathscr{V}}}^{p}={\underset{\sim}{V}}^{p}$ with the stronger norm $\|\underset{\sim}{v}\|_{{\underset{V}{V}}^{p}}=\|\underset{\sim}{v}\|_{0,2+\varepsilon}+\|\operatorname{div} \underset{\sim}{v}\|_{0}$ and let $\mathscr{W}^{p}=W^{p}$ with the stronger norm $\|w\|_{\mathscr{W}^{p}}=\|w\|_{0, t}$, where $t=\frac{4+2 \varepsilon}{\varepsilon}$. We can prove now the existence of a solution of (1.5).

THEOREM 3.1: For $\delta>0$ sufficiently small (dependent on $p$ ) and for $p$ sufficiently large, $\Phi$ maps a ball of radius $\delta$ centered at $\left(\pi^{p} z, P^{p} u\right)$ of ${\underset{\sim}{V}}^{p} \times \mathscr{W}^{p}$ into itself.

Proof: Note that $1 / t+1 /(2+\varepsilon)=1 / 2$. Let

$$
\left\|\pi^{p} z-\underset{\sim}{\mu}\right\|_{{\underset{\sim}{x}}^{p}} \leqslant \delta \quad \text { and } \quad\left\|P^{p} u-\rho\right\|_{\mathscr{W}^{p}} \leqslant \delta<1
$$

Ler us use lemma 3.1 on (3.8) with $\tau=P^{p} u-\kappa, \underset{\sim}{\underset{w}{w}}=\pi^{p} z-\lambda$, $\underset{\sim}{q}=\underset{\sim}{q}(\rho, \underset{\sim}{\mu}), \quad \eta=\eta(\rho)$ and $\theta=4-\varepsilon$. Observe that (1.7)-(1.9) and corollary 2.1 imply that, for $r>1 / 2, m=r+1$,

$$
\begin{align*}
\| \underset{\sim}{q}(\rho, \underset{\sim}{\mu}) & \left\|_{0}+\right\| \eta(\rho) \|_{0} \leqslant \mathscr{Q}\left[p^{1 / 2-r}\|z\|_{r}+p^{-m}\|u\|_{m}+\|u-\rho\|_{0,4}^{2}+\right. \\
& \left.+\|u-\rho\|_{0, t}\|z-\underset{\sim}{\mu}\|_{0,2+\varepsilon}\right] \\
\leqslant & \mathscr{Q}\left[p^{1 / 2-r}\|u\|_{r+1}+\left(\left\|u-P^{p} u\right\|_{0,4}+\left\|P^{p} u-\rho\right\|_{0,4}\right)^{2}+\right. \\
& +\left(\left\|u-P^{p} u\right\|_{0, t}+\left\|P^{p} u-\rho\right\|_{0, t}\right) \times \\
& \left.\times\left(\left\|z-\pi^{p} z\right\|_{0,2+\varepsilon}+\left\|\pi^{p} \underset{\sim}{z}-\underset{\sim}{\mu}\right\|_{0,2+\varepsilon}\right)\right] \\
\leqslant & \mathscr{Q}\left[p^{1 / 2-r}\|u\|_{r+1}+\left(p^{-m+3 / 4}\|u\|_{m}+\delta\right)^{2}+\right. \\
& \left.+\left(p^{5 / 2-r-4 /(2+\varepsilon)}\|u\|_{r+1}+\delta\right)\left(p^{-m+3 / 2-3 / t}\|u\|_{m}+\delta\right)\right] \\
\leqslant & \mathscr{Q}\left(\delta^{2}+p^{1 / 2-r}\|u\|_{r+1}\right), \tag{3.16}
\end{align*}
$$

where $\mathscr{2}$ depends on $\|u\|_{m}$. Therefore,

$$
\begin{align*}
\left\|P^{p} u-\kappa\right\|_{0,4-\varepsilon} \leqslant \mathscr{Q}\left[p^{-\varepsilon / 8} \|\right. & \pi^{p} z-\underset{\sim}{\lambda} \|_{0}+p^{-1-2 /(4-\varepsilon)} \times \\
& \left.\times\left\|\operatorname{div}\left(\pi^{p} z-\underset{\sim}{\lambda}\right)\right\|_{0}+\delta^{2}+p^{1 / 2-r}\right] \tag{3.17}
\end{align*}
$$

On the other hand, taking $\underset{\sim}{v}=\pi^{p} \underset{\sim}{z}-\underset{\sim}{\lambda}$ and $w=P^{p} u-\kappa$ in (3.8), we see that

$$
\begin{equation*}
\left\|\pi^{p} \underset{z}{z}-\underset{\sim}{\lambda}\right\|_{0} \leqslant \mathscr{2}\left[\left\|P^{p} u-\kappa\right\|_{0}+\|\underset{\sim}{q}\|_{0}+\|\eta\|_{0}\right] \tag{3.18}
\end{equation*}
$$

and, taking $w=\operatorname{div}\left(\pi^{p} \underset{\sim}{z}-\underset{\sim}{\lambda}\right)$ in the second equation of (3.8) results in

$$
\begin{equation*}
\left\|\operatorname{div}\left(\pi^{p} z-\underset{\sim}{\lambda}\right)\right\|_{0} \leqslant \mathscr{Q}\left[\left\|P^{p} u-\kappa\right\|_{0}+\|\underset{\sim}{q}\|_{0}+\left\|\eta_{0}\right\|\right] . \tag{3.19}
\end{equation*}
$$

Combining (3.17)-(3.19) yields the relation

$$
\left\|P^{p} u-\kappa\right\|_{0,4-\varepsilon} \leqslant \mathscr{Q}\left[p^{-\varepsilon / 8}\left\|P^{p} u-\kappa\right\|_{0}+\delta^{2}+p^{1 / 2-r}\right]
$$

which, for $p$ sufficiently large and $r=5 / 2$, implies that

$$
\begin{equation*}
\left\|P^{p} u-\kappa\right\|_{0,4-\varepsilon} \leqslant \mathscr{Q}\left[\delta^{2}+p^{-2}\right] \tag{3.20}
\end{equation*}
$$

where the constant $\mathscr{2}$ depends on $\|u\|_{7 / 2}$. Combining (3.20) with (1.9) we see that

$$
\begin{align*}
\left\|P^{p} u-\kappa\right\|_{0, t} & \leqslant \mathscr{2} p^{\frac{4}{4-\varepsilon}-\frac{2 t}{\varepsilon+2}}\left\|P^{p} u-\kappa\right\|_{0,4-\varepsilon}  \tag{3.21}\\
& \leqslant \mathscr{Q}\left(p^{1-\epsilon / 4} \delta^{2}+p^{-1-\varepsilon / 4}\right)
\end{align*}
$$

while (1.9), (3.18), (3.16), and (3.20) imply that

$$
\begin{align*}
\left\|\pi^{p} z-\underset{\sim}{\lambda}\right\|_{0,2+\varepsilon} & \leqslant \mathscr{Q} p^{2 \varepsilon /(2+\varepsilon)}\left\|\pi^{p} z-\underset{\sim}{\lambda}\right\|_{0} \\
& \leqslant \mathscr{Q}\left(p^{\varepsilon} \delta^{2}+p^{-2+\varepsilon}\right) . \tag{3.22}
\end{align*}
$$

Combining (3.19) and (3.22) yields

$$
\begin{equation*}
\left\|\pi^{p} z-\underset{\sim}{\lambda}\right\|_{\mathcal{V}^{p}} \leqslant \mathscr{Q}\left(p^{\varepsilon} \delta^{2}+p^{-2+\varepsilon}\right) \tag{3.23}
\end{equation*}
$$

We can now combine (3.21) and (3.23) in the bound

$$
\begin{equation*}
\left\|P^{p} u-\kappa\right\|_{\mathscr{W}^{p}}+\left\|\pi^{p} z-\underset{\sim}{\lambda}\right\|_{\boldsymbol{v}^{p}} \leqslant \mathscr{Q}_{1}\left(p^{1-\varepsilon / 4} \delta^{2}+p^{-1-\varepsilon / 4}\right) \tag{3.24}
\end{equation*}
$$

We want to choose $p$ and $\delta$ so that $\mathscr{Q}_{1} p^{1-\varepsilon / 4} \delta^{2} \leqslant \frac{\delta}{2}$ and $\mathscr{Q}_{1} p^{-1-\varepsilon / 4} \leqslant \frac{\delta}{2}$. vol. $26, n^{\circ} 7,1992$

Let $p \geqslant\left(2 \mathscr{Q}_{1}\right)^{4 / \varepsilon}$, so that $I=\left[2 \mathscr{Q}_{1} p^{-1-\varepsilon / 4}, \frac{p^{\varepsilon / 4-1}}{2 \mathscr{Q}_{1}}\right]$ is not empty. Then, for $\delta \in I$, (3.24) implies that

$$
\left\|P^{p} u-\kappa\right\|_{\mathscr{W}^{p}} \leqslant \delta \quad \text { and } \quad\left\|\pi^{p} z-\underset{\sim}{\lambda}\right\|_{\boldsymbol{\chi}^{p}} \leqslant \delta
$$

as we needed.
Remark 3.1 : Note that the choice $\delta=2 \mathscr{Q}_{1} p^{-1-\varepsilon / 4}$ in theorem 3.1 shows (using (1.7) and (1.8)) not only that (1.5) is solvable but also that, for $p \rightarrow \infty$, the solution of (1.5), $\left({\underset{z}{z}}^{p}, u^{p}\right)$, differs from $(\underset{\sim}{z}, u)$ in the ${\underset{\sim}{\mathcal{V}}}^{p} \times$ $\mathscr{W}^{p}$ norm by $0\left(p^{-1-\varepsilon / 4}\right)$ at most. We shall need this observation in order to arrive at the correct error estimates.

## 4. THE $L^{2}$-ERROR BOUNDS

Just as in [3], using (3.1)-(3.3) we now rewrite (3.5) in the form

$$
\left\{\begin{align*}
&(\alpha(u) \underset{\sim}{\zeta}, \underset{\sim}{v})-(\operatorname{div} \underset{\sim}{v}, \tau)+(\tau \underset{\sim}{\tilde{\sim}}, \underset{\sim}{v})=(\underset{\sim}{q}, \underset{\sim}{v}), \quad \underset{\sim}{v} \in V^{p}  \tag{4.1}\\
&(\operatorname{div} \underset{\sim}{\zeta}, w)+(\tilde{\gamma} \tau, w)=(\eta, w), \quad w \in W^{p},
\end{align*}\right.
$$

where $\underset{\sim}{\underset{\sim}{r}}=\underset{\sim}{z}-\underset{\sim}{z}, \quad \tau=P^{p} u-u^{p}, \quad \underset{\sim}{\tilde{\Gamma}}=\tilde{\boldsymbol{\alpha}}_{u}\left(u^{p}\right) \underset{\sim}{p}+{\underset{\sim}{\underset{\sim}{\underset{\sim}{x}}}}_{u}\left(u^{p}\right), \quad \tilde{\gamma}=\tilde{c}_{u}\left(u^{p}\right)$, $\underset{\sim}{q}=\left(P^{p} u-u\right) \underset{\sim}{\tilde{\tilde{q}}}$, and $\eta=\left(P^{p} u-u\right) \tilde{\gamma}$. Note that the left hand side of (4.1) corresponds to the mixed method for the operator $N: H^{2}(\Omega) \rightarrow$ $L^{2}(\Omega)$ given by

$$
N w=-\underset{\sim}{\nabla} \cdot(a(u) \underset{\sim}{\nabla} w+a(u) w \tilde{\sim} \tilde{\sim})+\tilde{\gamma} w .
$$

Therefore, if we show that its formal adjoint, $N^{*}$, has a bounded inverse $L^{2} \rightarrow H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$, then lemma 3.1 would apply to (4.1) without any change in the proof. Since we know that $M^{*}$ has a bounded inverse, all we need to do is to check that the operator norm of $M^{*}-N^{*}$ can be made arbitrarily small by taking $p$ large enough.

LEMMA 4.1: There exists a positive integer $p_{0}$ such that, for all $p \geqslant p_{0}, N^{*}$ has a bounded inverse $L^{2}(\Omega) \rightarrow H^{2}(\Omega) \cap H_{0}^{1}(\Omega) .\left(N^{*}\right.$ depends on $p$ through $\tilde{\gamma}$ and $\tilde{\sim}$ ).

Proof: Just as in [3], we have

$$
\begin{aligned}
&\left(M^{*}-N^{*}\right) \chi=a(u)\left\{\left[\bar{\alpha}_{u u} z+{\underset{\sim}{\beta}}_{u u}\right]\left(u-u^{p}\right)+\tilde{\alpha}_{u}\left(u^{p}\right)(z-\underset{\sim}{p})\right\} \times \\
& \times \underset{\sim}{\nabla} \chi+\bar{c}_{u u}\left(u-u^{p}\right) \chi, \quad \chi \in L^{2}(\Omega),
\end{aligned}
$$

where $\bar{\alpha}_{u u}=\frac{\alpha_{u}(u)-\tilde{\alpha}_{u}\left(u^{p}\right)}{u-u^{p}}$ and ${\underset{\sim}{\underset{\sim}{*}}}_{u u}$, and $\bar{c}_{u u}$, defined by analogous relations, are bounded functions in $\bar{\Omega}$. It follows from remark 31 and Sobolev's embedding theorem that

$$
\begin{aligned}
\left\|\left(M^{*}-N^{*}\right) \chi\right\|_{0} \leqslant & K\left[\|z\|_{0 \infty}\left\|u-u^{p}\right\|_{0 t}\|\nabla \chi\|_{02+\varepsilon}+\right. \\
& \left.+\|z-\underset{\sim}{p}\|_{02+\varepsilon}\|\nabla \chi\|_{0 t}+\left\|u-u^{p}\right\|_{0}\|\chi\|_{0 \infty}\right] \\
\leqslant & K\left(\|\nabla \chi \chi\|_{1}+\|\chi\|_{1+\varepsilon}\right) p^{-1-\varepsilon / 4} \\
\leqslant & K p^{-1}\|\chi\|_{2}
\end{aligned}
$$

as needed
To conclude, we establish the rate of convergence of $\left(\underset{\sim}{p}, u^{p}\right)$ to ( $z, u$ ).

Theorem 4.1 : Assume that the solution $u$ of (1.1) is in $H^{7 / 2}(\Omega)$ There is a positive constant $Q$, independent of $p$ but dependent on $\|u\|_{7 / 2+2 \varepsilon}$, such that, for $p$ suffictently large and $m \geqslant 7 / 2$,
ı)

$$
\left\|u-u^{p}\right\|_{0} \leqslant Q p^{1-m}\|u\|_{m}
$$

ir)

$$
\left\|z-\not z_{0} \leqslant Q p^{3 / 2-m}\right\| u \|_{m}
$$

llı)

$$
\|\operatorname{div}(z-\not z)\|_{0} \leqslant Q p^{2-m}\|u\|_{m}
$$

Proof In view of remark 31 and lemma 4.1, we can use lemma 3.1 on (41) with $\theta=2$. Thus,

Note that remark 3.1 together with (1.7) lead to the following estimate for $r \geqslant 0, m>3 / 2$,

$$
\begin{align*}
\left\|{\underset{\sim}{q}}_{0}+\right\| \eta \|_{0} & =\left\|\left(P^{p} u-u\right) \tilde{\sim}\right\|+\left\|\left(P^{p} u-u\right) \tilde{\gamma}\right\| \leqslant \\
& \leqslant K\left(\left\|P^{p} u-u\right\|_{0}+\left\|\left(P^{p} u-u\right) \underset{z}{l}\right\|_{0}\right) \\
& \leqslant K\left[p^{-r}\|u\|_{r}\left(1+\|z\|_{0 \infty}\right)+\left\|\not{ }^{p}-z\right\|_{0,2+\varepsilon}\left\|P^{p} u-u\right\|_{0, t}\right] \\
& \leqslant K\left(p^{-r}\|u\|_{r}+p^{-1-\varepsilon / 4} p^{3 / 2-3 \varepsilon /[2(2+\varepsilon)]-m}\|u\|_{m}\right) \\
& \leqslant K p^{1-m-\varepsilon}\|u\|_{m} \tag{array}
\end{align*}
$$

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Combining (4.2), (4.3), (3.18), (3.19), (1.7) and (1.8) yields,

$$
\begin{aligned}
& \|\tau\|_{0} \leqslant C\left[p^{-1 / 2}\left(\left\|z-\pi^{p} z\right\|_{0}+\left\|\pi^{p} z-\not z^{p}\right\|_{0}\right)+p^{-2}\left(\left\|\operatorname{div} \underset{z}{ }-P^{p} \operatorname{div} z\right\|_{0}\right.\right. \\
& \left.\left.\quad \quad+\left\|\operatorname{div}\left(\pi^{p} z-\not z^{p}\right)\right\|_{0}\right)+p^{1-m-\varepsilon}\|u\|_{m}\right] \\
& \leqslant \\
& \quad C\left[p^{-1 / 2}\|\tau\|_{0}+p^{-1 / 2} p^{1 / 2-r}\|u\|_{r+1}+p^{-2} p^{-s}\|u\|_{s+2}+\right. \\
& \left.\quad+p^{-1 / 2} p^{-1-\varepsilon / 4}+p^{1-m-\varepsilon}\|u\|_{m}\right], r>1 / 2, \quad s \geqslant 0, m>3 / 2,
\end{aligned}
$$

which, for $p$ sufficiently large, leads to

$$
\begin{equation*}
\|\tau\|_{0} \leqslant C p^{1-m}\|u\|_{m}, \quad m \geqslant 2 \tag{4.4}
\end{equation*}
$$

where the constant $C$ depends on $\|u\|_{7 / 2}$. The first part of the theorem is an immediate consequence of (1.7) and (4.4). On the other hand, it follows from (1.8), (3.18), (4.3) and (4.4), that

$$
\begin{aligned}
\|z-\not z\|_{0} & \leqslant\left\|z-\pi^{p} z\right\|_{0}+\left\|\pi^{p} z-\not z^{p}\right\|_{0} \\
& \leqslant C\left[p^{3 / 2-m}\|u\|_{m}+p^{1-m}\|u\|_{m}\right]
\end{aligned}
$$

which proves the second part of the theorem.
Finally, we deduce from (3.19), (1.7), (4.3) and (4.4) that

$$
\begin{aligned}
\left\|\operatorname{div}\left(z-\not z^{p}\right)\right\|_{0} & \leqslant\left\|\operatorname{div} z-P^{p} \operatorname{div} z\right\|_{0}+\left\|\operatorname{div}\left(\pi^{p} z-\not z^{p}\right)\right\|_{0} \\
& \leqslant C\left[p^{2-m}\|u\|_{m}+p^{1-m}\|u\|_{m}\right]
\end{aligned}
$$

which gives iii).
Remark 4.1: The estimate for the error in $\underset{\sim}{ }$ is the best we could hope for in view of (1.8). The estimate for the error in $\operatorname{div} \underset{\sim}{z}$ is optimal in rate and regularity, while the one for $u$ is probably not sharp in view of (1.7).

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