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# Houde Han <br> The error estimates for the infinite element method for eigenvalue problems 

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# THE ERROR ESTIMATES FOR THE INFINITE ELEMENT METHOD FOR EIGENVALUE PROBLEMS (*) 

by Houde Han ( ${ }^{1}$ )<br>Communicated by P G Ciarlet

Resume - Cet article analyse les estımations d erreur pour la methode des elements «infinis» Le domaine est partage en un nombre infint de triangles On utlise des elements lineaires pour les deux espaces en cause
Abstract - This paper analyses the error estimates for infinite element method appled to eigenvalue problems The domain is divided into infintely many triangles and linear elements are used for the trial space and test space

## 1. INTRODUCTION

In the numerical solution of elliptic boundary value problems, it is wellknown that the presence of corners in the domain can cause a loss of accuracy in the solution Many methods have been developed to overcome the loss of accuracy, such as the use of singular functions [9], mesh refinements [10], and the infinite element method [11,2] The infinite element method may be considered as a kind of mesh refinement, but has the advantages that the refinement is easy to construct, the stiffness matrix is calculated more efficiently, and an approximate solution is obtamed which itself has a singularity at the corner Recently, we showed how the infinite element method may be applied to eigenvalue problems on domans with corners [1] In this paper, we obtain the error estimates for the infinite element method, when applied to an elgenvalue problem

[^0]We consider the following eigenvalue problem.

$$
\begin{align*}
\Delta u+\lambda u & =0, & & \text { in } \Omega  \tag{1.1}\\
\frac{\partial u}{\partial v} & =0, & & \text { on } \Gamma_{0} \text { and } \Gamma_{M},  \tag{1.2}\\
u & =0, & & \text { on } \Gamma^{0} . \tag{1.3}
\end{align*}
$$

Here $\Omega$ is an open polygonal domain in the $x_{1}, x_{2}$ plane and $\Gamma=\bigcup_{J=0}^{M} \Gamma_{J}$ is the boundary of $\Omega$, with the $\Gamma$, s denoting the side of $\Omega$.

$$
\Gamma^{0}=\bigcup_{J=1}^{M-1} \Gamma_{J} \text { and } \frac{\partial u}{\partial v} \text { denotes the outward normal derivative of } u \text { on } \Gamma_{0}
$$ and $\Gamma_{M} . A_{i}(i=0,1, \ldots, M)$ denotes the vertice of $\Omega$ with $\varphi_{\imath}$ being the interior angle of $\Omega$ as shown in figure 1 .



Figure 1.

For the sake of simplicity, we suppose

$$
\begin{equation*}
\pi<\varphi_{0} \leqslant 2 \pi ; \quad 0<\varphi_{1}, \quad \varphi_{M} \leqslant \frac{\pi}{2} ; \quad 0<\varphi_{J} \leqslant \pi \quad(j=2, \ldots, M-1) \tag{1.4}
\end{equation*}
$$

Without losing generality, we assume $A_{0}$ is the origin of the rectangular coordinate system.

Let $W^{m, p}(\Omega)$ denote the Sobolev space on $\Omega$ with norm

$$
\|u\|_{W^{m, p}(\Omega)}=\left\{\sum_{i=0}^{m} \iint_{\Omega} \sum_{\alpha_{1}+\alpha_{2}=\imath}\left|\frac{\partial^{l} u}{\partial x_{1}^{\alpha_{1}} \partial x_{2}^{\alpha_{2}}}\right|^{p} d x_{1} d x_{2}\right\}^{1 / p}
$$

where $m$ is a non-negative integer and $p$ is a positive real number. As usual, when $p=2, w^{m, 2}(\Omega)$ is denoted by $H^{m}(\Omega)$; when $m=0, w^{0, p}(\Omega)$ is denoted by $L_{p}(\Omega)$. Moreover, we shall introduce the Sobolev space with weight, $H^{m, t}(\Omega)$, with the norm

$$
\left.\begin{array}{rl}
\|u\|_{H^{m, t}(\Omega)}^{2}= & \|u\|_{H^{m-1}(\Omega)}^{2}+ \\
& \quad+\iint_{\Omega} \sum_{\alpha_{1}+\alpha_{2}=m} r^{2 t}\left|\frac{\partial^{m} u}{\partial x_{1}^{\alpha_{1}} \partial x_{2}^{\alpha_{2}}}\right|^{2} d x_{1} d x_{2}, m \geqslant 1 \tag{1.5}
\end{array}\right\}
$$

where $r^{2}=x_{1}^{2}+x_{2}^{2}$ and $t$ is a real number.
Let

$$
\stackrel{*}{H}^{1}(\Omega)=\left\{u \mid u \in H^{1}(\Omega) \text { and } u=0 \text { on } \Gamma^{0} \quad \text { (in the trace sense) }\right\} .
$$

$\stackrel{*}{H}^{1}(\Omega)$ is a subspace of $H^{1}(\Omega)$.
We know that the eigenvalue problem (1.1)-(1.3) has the following variational form : find a complex numbers $\lambda$ and a nonzero $u \in \stackrel{*}{H}^{1}(\Omega)$ such that

$$
\begin{equation*}
B(u, v)=\lambda J(u, v), \quad \forall v \in \stackrel{*}{H}^{1}(\Omega) \tag{1.6}
\end{equation*}
$$

where

$$
\begin{aligned}
B(u, v) & =\iint_{\Omega}\left(\frac{\partial u}{\partial x_{1}} \frac{\partial \bar{v}}{\partial x_{1}}+\frac{\partial u}{\partial x_{2}} \frac{\partial \bar{v}}{\partial x_{2}}\right) d x_{1} d x_{2} \\
J(u, v) & =\iint_{\Omega} u \bar{v} d x_{1} d x_{2}
\end{aligned}
$$

Let us now recall the procedure of obtaining the approximate solution of (1.6) using the infinite element method [1]. In the first step the domain $\Omega$ is divided into infinitely many similar element layers $D_{1}, D_{2}, \ldots, D_{k}, \ldots$, where $D_{k}$ denotes the $k$-th layer. Every layer is divided into several triangles in the same manner as in [1]. $0<\xi<1$ is the constant of proportionality. Point $A_{0}$ is the center of similarity. Therefore

$$
\Omega=\bigcup_{k=1}^{\infty} D_{k}
$$

Let

$$
\Omega_{N}=\bigcup_{k=N+1}^{\infty} D_{k} \quad(N=1,2, \ldots)
$$

Let $h$ denote the length of the side which is the longest among all triangles in $\Omega$. Moreover, we suppose the angles of all triangles are greater than $\gamma_{0}$, where $0<\gamma_{0}<\pi / 3$ is a constant. This criterion is called the smallest interior angle condition.

For this partition we introduce a closed subspace of $\stackrel{*}{H}^{1}(\Omega)$ denoted by $\stackrel{\circ}{S}(\Omega)$ :
$\stackrel{\circ}{S}(\Omega)=\left\{u \mid u \in \stackrel{*}{H}^{1}(\Omega) \cap C(\Omega)\right.$ and $u$ is a linear function on each triangle $\}$.
Using the space $\stackrel{\circ}{S}(\Omega)$ instead of $\stackrel{*}{H}^{1}(\Omega)$ in the problem (1.6) we obtain the following eigenvalue problem : find a complex numbers $\lambda$ and a nonzero $u \in \stackrel{\circ}{S}(\Omega)$, such that

$$
\begin{equation*}
B(u, v)=\lambda J(u, v), \quad \forall v \in \stackrel{\circ}{S}(\Omega) \tag{1.7}
\end{equation*}
$$

The eigenvalue problem (1.7) is the discretized model of (1.6). (1.7) is equivalent to the following eigenvalue problem of the pencil of the infinite dimensional matrices

$$
\left[Q_{1}-\lambda Q_{2}\right]\left[\begin{array}{c}
y_{1}  \tag{1.8}\\
y_{2} \\
\vdots \\
y_{N} \\
\vdots
\end{array}\right]=0
$$

where

$$
\begin{aligned}
& Q_{1}=\left(\begin{array}{ccc}
K & -A^{T} & \\
-A & K & -A^{T} \\
& \ldots \ldots \ldots \ldots \ldots
\end{array}\right)_{\infty \times \infty},
\end{aligned}
$$

$K=K_{0}+K_{0}^{\prime}$ and $L=\xi^{2} L_{0}+L_{0}^{\prime}$. The matrices $K_{0}, K_{0}^{\prime}, A$ and $L_{0}, L_{0}^{\prime}$, $D$ constitute the stiffness matrix of the $k$-th layer

$$
\left(\begin{array}{rr}
K_{0} & -A^{T} \\
-A & K_{0}^{\prime}
\end{array}\right)-\lambda \xi^{2(k-1)}\left(\begin{array}{rr}
L_{0} & -D^{T} \\
-D & L_{0}^{\prime}
\end{array}\right)
$$

corresponding to equation (1.1). We do not have a method to solve the eigenvalue problem (1.8) in the present form. By means of a technique in [1], the eigenvalue problem (1.8) was changed to a eigenvalue problem of the pencil of finite dimensional matrices :

$$
\left[Q_{1}^{N}-\lambda Q_{2}^{N}\right]\left[\begin{array}{c}
y_{1}  \tag{1.9}\\
\vdots \\
y_{N}
\end{array}\right]=0,
$$

where $y_{1}, \ldots, y_{N}$ are $M$-dimensional column vectors,

$$
\begin{aligned}
& Q_{1}^{N}=\left(\begin{array}{rrrrl}
K & -A^{T} & & & \\
-A & K & -A^{T} & & \\
\\
& \ldots \ldots \ldots \ldots \cdots & & \\
\\
& & -A & K & -A^{T} \\
& & & -A & K-A^{T} X(0)
\end{array}\right)_{(M \times N) \times(M \times N)},
\end{aligned}
$$

Here $X(0)$ is a solution of the matrix equation

$$
-A+K X(0)-A^{T}(X(0))^{2}=0
$$

of which the solution can be obtained using the direct method or the iterative method [2], [3].

The solutions of (1.9) are the approximate solutions of (1.6). In this paper we shall study error estimates.

## 2. THE VARIATIONAL FORM OF (1.9)

Prior to discussing the error estimate for the solution of (1.6) and (1.9) we introduce a variational form of the matrix eigenvalue problem (1.9). Consider the following eigenvalue problem. Find a complex numbers $\lambda$ and a nonzero function $u \in \stackrel{\circ}{S}(\Omega)$ such that

$$
\begin{equation*}
B(u, v)=\lambda J_{N}(u, v), \quad \forall v \in \stackrel{\circ}{S}(\Omega) \tag{2.1}
\end{equation*}
$$

where

$$
\begin{equation*}
J_{N}(u, v)=\iint_{\Omega / \Omega_{N}} u \bar{v} d x_{1} d x_{2} \tag{2.2}
\end{equation*}
$$

For the problem (2.1), if $k \leqq N$ the stiffness matrix of the $k$-th layer is

$$
\left(\begin{array}{cr}
K_{0} & -A^{T} \\
-A & K_{0}^{\prime}
\end{array}\right)-\lambda \xi^{2(k-1)}\left(\begin{array}{rr}
L_{0} & -D^{T} \\
-D & L_{0}^{\prime}
\end{array}\right)
$$

and, when $k>N$, the stiffness matrix of $k$-th layer is

$$
\left(\begin{array}{cr}
K_{0} & -A^{T} \\
-A & K_{0}^{\prime}
\end{array}\right)
$$

It is straightforward to show that the problem (2.1) is equivalent to the following eigenvalue problem :

$$
\begin{aligned}
& \quad(K-\lambda L) y_{1}-\left(A^{T}-\lambda \xi^{2} D^{T}\right) y_{2}=0, \\
& -\left(A-\lambda \xi^{2(k-1)} D\right) y_{k-1}+\left(K-\lambda \xi^{2(k-1)} L\right) y_{k}- \\
& -\left(A^{T}-\lambda \xi^{2 k} D^{T}\right) y_{k+1}=0, \quad k=2, \ldots, N-1, \quad(2 \\
& -\left(A-\lambda \xi^{2(N-1)} D\right) y_{N-1}+\left(K-\lambda \xi^{2(N-1)} L_{0}^{\prime}\right) y_{N}-A^{T} y_{N+1}=0 \\
& -A y_{k-1}+K y_{k}-A^{T} y_{k+1}=0, \quad k=N+1, N+2, \ldots
\end{aligned}
$$

From the above equations it is seen that $\lambda$ does not appear in the last part of (2.3). Therefore, we consider the system of infinitely many equations

$$
\left.\begin{array}{r}
-A y_{k-1}+K y_{k}-A^{T} y_{k+1}=0  \tag{2.4}\\
k=N+1, N+2, \ldots
\end{array}\right\}
$$

From Lemma 1.4 in [1] we know that, for any given $y_{N}$, problem (2.4) has a unique solution $\left\{y_{N}, y_{N+1}, \ldots\right\}$ which corresponds to the function $u_{N} \in H^{1}\left(\Omega_{N}\right)$ and

$$
\begin{equation*}
y_{k+1}=X(0) y_{k}, \quad k=N, N+1, \ldots \tag{2.5}
\end{equation*}
$$

We have:
Lemma 2.1 : Suppose $\lambda^{h, N}$ is an eigenvalue of (2.3) and $u^{h, N}$ is an eigenfunction corresponding to $\lambda^{h, N}$. Then $y_{0}^{h, N}, y_{1}^{h, N}, \ldots, y_{N}^{h, N}$ are not all zero-vectors ( $u^{h, N}$ corresponds to the sequence $y_{0}^{h, N}, y_{1}^{h, N}, \ldots, y_{k}^{h, N}, \ldots$ ).
Proof: Suppose the conclusion is false, then $y_{0}^{h, N}=y_{i}^{h, N}=\cdots=y_{N}^{h, N}=0$. Since $u^{h, N}$ is an eigenfunction corresponding to $\lambda^{h, N}$, we know that

$$
y_{1}^{h_{1}, N}, y_{2}^{h_{2}^{N}}, \ldots, y_{k}^{h_{k}, N}, \ldots
$$

satisfy the system (2.3) in which $\lambda^{h, N}$ is used instead of $\lambda$. Moreover, from (3.5) we obtain

$$
\mathrm{y}_{N+l}^{h, N}=(X(0))^{l} y_{N}^{h, N}=0, \quad l=1,2, \ldots
$$

Consequently, we have $y_{0}^{h, N}=y_{1}^{h, N}=\cdots=y_{N}^{h, N}=y_{N+1}^{h, N}=\cdots=0$. Namely $u^{h, N} \equiv 0$. This contradicts the fact that $u^{h, N}$ is an eigenfunction. This contradiction shows that our conclusion is correct.

Lemma 2.2 : The eigenvalue problem (2.3) is equivalent to (1.9).
Proof: Suppose $\lambda^{h, N}$ is an eigenvalue of (2.3) and $\left\{y_{1}^{h, N}, \ldots, y_{k}^{h, N}, \ldots\right\}$ is an eigenvector corresponding to $\lambda^{h, N}$. Since $y_{N+1}^{h, N}=X(0) y_{N}^{h, N}$, we know that $\lambda^{h, N} ; y_{1}^{h, N}, y_{2}^{h, N}, \ldots, y_{N}^{h, N}$ satisfies (1.9) and that, from Lemma 2.1,

$$
y_{1}^{h, N}, \ldots, y_{N}^{h, N}
$$

are not all zero vectors. Therefore $\lambda^{h, N} ; y_{1}^{h, N}, \ldots, y_{N}^{h, N}$ is a solution of the eigenvalue problem (1.9).

On the other hand, if $\lambda^{h, N}, y_{1}^{h, N}, \ldots, y_{N}^{h, N}$ is a solution of (1.9) let

$$
y_{N+l}^{h_{N}, N}=(X(0))^{l} y_{N}^{k, N}, \quad l=1,2, \ldots
$$

Obviously, $\lambda^{h, N}, y_{1}^{h, N}, y_{2}^{h, N}, \ldots, y_{N}^{h, N}, y_{N+1}^{h, N}, \ldots$ is a solution of the eigenvalue problem (2.3).

From the above lemma we have
Lemma 2.3 : The eigenvalue problem (2.1) is equivalent to (1.9). The variational form (2.1), instead of (1.9), will be used for the following discussion.

## 3. ERROR ESTIMATE

Before the discussion of the error estimate we recall some results which are used in this paper.

Lemma 3.1 : There exists a constant $\alpha>0$ such that

$$
\begin{equation*}
B(u, u) \geqslant \alpha\|u\|_{H^{1}(\Omega)}^{2}, \quad \forall u \in \stackrel{*}{H}^{1}(\Omega) \tag{3.1}
\end{equation*}
$$

Lemma 3.2: For any $f \in H^{0, t}(\Omega)$ (where $1-\pi / \varphi_{0}<t<1$ ), then

$$
\begin{equation*}
B(u, v)=J(f, v), \quad \forall v \in \stackrel{*}{H}^{1}(\Omega) \tag{3.2}
\end{equation*}
$$

has a unique solution $u \in H^{2, t}(\Omega) \cap \stackrel{*}{H}^{1}(\Omega)$ and there exists a constant $c>0$ independent of $f$ such that

$$
\begin{equation*}
\|u\|_{H^{2, t}(\Omega)} \leqq c\|f\|_{H^{0, t}(\Omega)} \tag{3.3}
\end{equation*}
$$

Lemma 3.3 : For any $u \in H^{2, t}(\Omega) \cap \stackrel{*}{H}^{1}(\Omega)\left(1-\pi / \varphi_{0}<t<1\right)$ there exists a function $u_{I} \in \dot{S}(\Omega)$ such that

$$
\begin{equation*}
\left\|u-u_{I}\right\|_{H^{1}(\Omega)} \leqslant c h\|u\|_{H^{2, t}(\Omega)} \tag{3.4}
\end{equation*}
$$

where $c$ is a constant which is independent of $u$ and $h$. The proof of Iemma 3.1 can be found in [8]. The Lemma 3.2 is quoted from [4], whereas Lemma 3.3 is from Theorem 1 in [5].

Lemma 3.4 : For any $f \in H^{0}(\Omega)$, the problem

$$
\begin{equation*}
B(u, v)=j(f, v), \quad \forall v \in \stackrel{\circ}{S}(\Omega) \tag{3.5}
\end{equation*}
$$

has a unique solution $u^{h} \in \stackrel{\circ}{S}(\Omega)$ and

$$
\begin{equation*}
\left\|u^{h}\right\|_{H^{1}(\Omega)} \leqslant \frac{1}{\alpha}\|f\|_{H^{0}(\Omega)} \tag{3.6}
\end{equation*}
$$

Proof : Since $f \in H^{0}(\Omega)$, we know that $f$ is a bounded linear functional on the space $\stackrel{\circ}{S}(\Omega)$. Moreover, from Lemma 3.1 we note that $B(u, v)$ is a positive definite bilinear form. Thus there exists a unique solution $u^{h} \in S(\Omega)$ such that

$$
\begin{equation*}
B\left(u^{h}, v\right)=J(f, v), \quad \forall v \in \stackrel{\circ}{S}(\mathbf{\Omega}) \tag{3.5}
\end{equation*}
$$

Taking $v=u^{h}$ in (3.5)' we obtain (3.6) from Lemma 3.1. Similarly, we can prove the following Lemma.

Lemma 3.5 : For any $f \in H^{0}(\Omega)$ the problem

$$
\begin{equation*}
B(u, v)=J_{N}(f, v), \quad \forall v \in \stackrel{\circ}{S}(\Omega) \tag{3.7}
\end{equation*}
$$

has a unique solution $u^{h, N} \in \stackrel{\circ}{S}(\Omega)$ and

$$
\begin{equation*}
\left\|u^{h, N}\right\|_{H^{1}(\Omega)} \leqq \frac{1}{\alpha}\|f\|_{H^{0}(\Omega)} \tag{3.8}
\end{equation*}
$$

From Lemma 3.2 we know that, for any $f \in H^{0}(\Omega)$, problem (3.2) always has a unique solution $u \in H^{2, t}(\Omega) \cap \stackrel{*}{H}^{1}(\Omega)$. Therefore, we denote by $T$ the linear operator which maps $f$ to $u$. Moreover, we know that $T$ is a compact operator from $H^{0}(\Omega)$ to $H^{0}(\Omega)$. It satisfies

$$
\begin{equation*}
B(T f, v)=J(f, v), \quad \forall v \in \stackrel{*}{H}^{1}(\Omega) . \tag{3.9}
\end{equation*}
$$

An eigenvalue of $T$ is a real number $\mu$ (because $T$ is a self-adjoint operator) such that $T u=\mu u$ for some non-zero function $u \in H^{0}(\Omega)$. Clearly, for any non-zero eigenvalue $\mu$ of $T$, we have that $\lambda=1 / \mu$ is an eigenvalue of (1.6). On the other hand, since $\lambda$ is an eigenvalue of (1.6), $\mu=1 / \lambda$ is also an eigenvalue of $T$.

Similarly, from Lemma 3.4 we know that, for any $f \in H^{0}(\Omega)$, (3.5) has a unique solution $u^{h} \in \stackrel{\circ}{S}(\Omega) \subset H^{1}(\Omega)$. We denote by $T_{h}$ the linear operator which maps $f$ to $u^{h}$. $T_{h}$ is compact from $H^{0}(\Omega)$ to $H^{0}(\Omega)$. $T_{h}$ satisfies

$$
\begin{equation*}
B\left(T_{h} f, v\right)=J(f, v), \quad \forall v \in \stackrel{\circ}{S}(\Omega) \tag{3.10}
\end{equation*}
$$

Obviously, $\lambda^{h}$ is an eigenvalue of (1.7) if and only if $\mu(h)=1 / \lambda^{h}$ is an eigenvalue of $T_{h}$.

Let $T_{h}^{N}$ denote the linear operator which takes $f \in H^{0}(\Omega)$ to $u^{h, N} ; T_{h}^{N}$ satisfies

$$
\begin{equation*}
B\left(T_{h}^{N} f, v\right)=J_{N}(f, v), \quad \forall v \in \stackrel{\circ}{S}(\Omega) \tag{3.11}
\end{equation*}
$$

Obviously, for any non-zero eigenvalue $\mu^{N}(h)$ of $T_{h}^{N}, \lambda^{h, N}=1 / \mu^{N}(h)$ is an eigenvalue of problem (2.1). Namely, $\lambda^{h, N}=1 / \mu^{N}(h)$ is an eigenvalue of (1.9). Since $u^{h, N} \in \stackrel{\circ}{S}(\Omega) \subset H^{1}(\Omega)$, we know that $T_{h}^{N}$ is a compact operator from $H^{0}(\Omega)$ to $H^{0}(\Omega)$.

Now the error estimates for the solutions of (1.6) and (1.9) are reduced to those for the eigenvalues and eigenfunctions of the compact operators $T$ and $T_{h}^{N}$.

Let :

$$
\|T\|=\sup _{\substack{f \neq 0 \\ f \in H^{0}(\Omega)}} \frac{\|T f\|_{H^{0}(\Omega)}}{\|f\|_{H^{0}(\Omega)}}
$$

denote the norm of the operator $T$. We need to estimate the error \| $T-T_{h}^{N} \|$. vol $16, n^{\circ} 2,1982$

Moreover, we have

$$
\begin{equation*}
\left\|T-T_{h}^{N}\right\| \leqslant\left\|T-T_{h}\right\|+\left\|T_{h}-T_{h}^{N}\right\| . \tag{3.12}
\end{equation*}
$$

In order to get the error $\left\|T-T_{h}^{N}\right\|$, we estimate $\left\|T-T_{h}\right\|$ and $\left\|T_{h}-T_{h}^{N}\right\|$. We have :

Lemma 3.6 : For any $f \in H^{0}(\Omega)$ there exists a constant $C$ which is independent off such that

$$
\begin{equation*}
\inf _{v \in \mathcal{S}_{(\Omega)}}\|T f-v\|_{H^{1}(\Omega)} \leqq C h\|f\|_{H^{0}(\Omega)} \tag{3.13}
\end{equation*}
$$

Proof : By Lemma 3.3, we know that $T f \in H^{2, t}(\Omega)$ and there exists a function $v_{I} \in \stackrel{\circ}{S}(\Omega)$ such that

$$
\left\|T f-v_{I}\right\|_{H^{1}(\Omega)} \leqq C h\|T f\|_{H^{2, t}(\Omega)}
$$

On the other hand, we know from Lemma 3.2

$$
\|T f\|_{H^{2, t}(\Omega)} \leqq C\|f\|_{H^{0, t}(\Omega)} \leqq C\|f\|_{H^{0}(\Omega)}
$$

here $C$ is a constant independent of $h$ and $f$. Therefore we get

$$
\left\|T f-v_{I}\right\|_{H^{1}(\Omega)} \leqq C h\|f\|_{H^{0}(\Omega)}
$$

Thus

$$
\inf _{v \in \mathcal{S}(\Omega)}\|T f-v\|_{H^{1}(\Omega)} \leqq\left\|T f-v_{I}\right\| \leqq C h\|f\|_{H^{0}(\Omega)}
$$

Lemma 3.7 : There exists a constant $C$ such that

$$
\begin{equation*}
\left\|T-T_{h}\right\| \leqq C h^{2} \tag{3.14}
\end{equation*}
$$

Proof : Since $\stackrel{\circ}{S}(\Omega)$ is a subspace of $\stackrel{*}{H}^{1}(\Omega)$, for any $f \in H^{\circ}(\Omega)$ from (3.2) and (3.5)', we obtain

$$
\begin{equation*}
B\left(\left(T-T_{h}\right) f, v\right)=0, \quad \forall v \in \stackrel{\circ}{S}(\Omega) \tag{3.15}
\end{equation*}
$$

By Lemma 3.1 we have

$$
\begin{aligned}
\left\|\left(T-T_{h}\right) f\right\|_{H^{1}(\Omega)}^{2} & \leqq \frac{1}{\alpha}\left|B\left(\left(T-T_{h}\right) f,\left(T-T_{h}\right) f\right)\right| \\
& =\frac{1}{\alpha}\left|B\left(\left(T-T_{h}\right) f,\left(T-T_{h}\right) f-v\right)\right|, \quad \forall v \in \stackrel{\circ}{S}(\Omega)
\end{aligned}
$$

Therefore, we get

$$
\left\|\left(T-T_{h}\right) f\right\|_{H^{1}(\Omega)} \leqslant \frac{1}{\alpha} \inf _{v \in S_{(\Omega)}}\|T f-v\|_{H^{1}(\Omega)}
$$

From Lemma 3.6 it follows that

$$
\begin{equation*}
\left\|\left(T-T_{h}\right) f\right\|_{H^{1}(\Omega)} \leqslant C h\|f\|_{H^{0}(\Omega)} \tag{3.16}
\end{equation*}
$$

On the other hand, for any $f, \psi \in H^{0}(\Omega)$ we have

$$
\begin{aligned}
J\left(\left(T-T_{h}\right) f, \psi\right) & =B\left(\left(T-T_{h}\right) f, T \psi\right) \\
& =B\left(\left(T-T_{h}\right) f, T \psi-v\right), \quad \forall v \in \stackrel{\circ}{S}(\Omega)
\end{aligned}
$$

Consequently

$$
\begin{aligned}
\left|J\left(\left(T-T_{h}\right) f, \psi\right)\right| & \leqq\left\|\left(T-T_{h}\right) f\right\|_{H^{1}(\Omega)} \inf _{v \in \mathcal{S}(\Omega)}\|T \psi-v\|_{H^{1}(\Omega)} \\
& \leqslant C h^{2}\|f\|_{H^{0}(\Omega)}\|\psi\|_{H^{0}(\Omega)}
\end{aligned}
$$

The last inequality is from (3.16) and Lemma 3.6. Thus we have

$$
\left\|\left(T-T_{h}\right) f\right\|_{H^{0}(\Omega)}=\sup _{\|\Psi\|_{H^{0}(\Omega)}=1}\left|J\left(\left(T-T_{h}\right) f, \psi\right)\right| \leqq C h^{2}\|f\|_{H^{0}(\Omega)}
$$

Finally, from the above inequality we get (3.14).
. Lemma 3.8 : For any fixed $0<\varepsilon<1$, there exists a constant $C(\varepsilon)$ which is independent of $N$ and $\xi$ such that

$$
\begin{align*}
&\left|J(u, v)-J_{N}(u, v)\right| \leqq C(\varepsilon)\left(\xi^{N}\right)^{1-\varepsilon}\|u\|_{H^{0}(\Omega)}\|v\|_{H^{1}(\Omega)} \\
& \forall u \in H^{0}(\Omega), \quad v \in H^{1}(\Omega) \tag{3.17}
\end{align*}
$$

Proof: Based on the imbedding theorem of the Sobolev space [7] we know that for any real number $p \geqslant 1$ there exists a constant $C_{1}(p)$ such that

$$
\begin{equation*}
\|v\|_{L_{p}} \leqslant C_{1}(p)\|v\|_{H^{1}(\Omega)}, \quad \forall v \in H^{1}(\Omega) \tag{3.18}
\end{equation*}
$$

From the definition of $J_{N}(u, v)$ and upon repeated use of the Cauchy inequality, we obtain

$$
\begin{aligned}
\left|J(u, v)-J_{N}(u, v)\right| & =\left|\iint_{\Omega_{N}} u \bar{v} d x_{1} d x_{2}\right| \\
& \leqq\left\{\iint_{\Omega_{N}}|u|^{2} d x_{1} d x_{2}\right\}^{1 / 2}\left\{\iint_{\Omega_{N}}|v|^{2} d x_{1} d x_{2}\right\}^{1 / 2}
\end{aligned}
$$

$$
\begin{aligned}
& \leqq\|u\|_{H^{0}(\Omega)}\left\{\iint_{\Omega_{N}}|v|^{2} d x_{1} d x_{2}\right\}^{1 / 2} \\
& \leqq\|u\|_{H^{0}(\Omega)}\left(\operatorname{meas} \Omega_{N}\right)^{\frac{1}{2}-\frac{1}{p}}\left(\iint_{\Omega_{N}}|v|^{p} d x_{1} d x_{2}\right)^{\frac{1}{p}} \\
& \leqq(\operatorname{meas} \Omega)^{\frac{1}{2}-\frac{1}{p}}\left(\xi^{N}\right)^{1-\frac{2}{p}}\|u\|_{H^{0}(\Omega)}\|v\|_{L_{p}(\Omega)} \\
& \quad \forall p>2, \quad \forall u \in H^{0}(\Omega), \quad v \in H^{1}(\Omega)
\end{aligned}
$$

Taking $p=2 / \varepsilon, C(\varepsilon)=(\text { meas } \Omega)^{\frac{1-\varepsilon}{2}} C_{1}(2 / \varepsilon)$ and using (3.18), we obtain (3.17).

Lemma 3.9 : For any fixed $0<\varepsilon<1$, there exists a constant $C(\varepsilon)$, such that

$$
\begin{equation*}
\left\|T_{h}-T_{h}^{N}\right\| \leqq \frac{C(\varepsilon)}{\alpha}\left(\xi^{N}\right)^{1-\varepsilon} . \tag{3.19}
\end{equation*}
$$

Proof : Equations (3.10) and (3.11) yield

$$
\begin{equation*}
B\left(\left(T_{h}-T_{h}^{N}\right) f, v\right)=J(f, v)-J_{N}(f, v), \quad \forall v \in \stackrel{\circ}{S}(\Omega) \tag{3.20}
\end{equation*}
$$

and from Lemma 3.8 we derive

$$
\left|B\left(\left(T_{h}-T_{h}^{N}\right) f, v\right)\right| \leqslant \frac{C(\varepsilon)}{\alpha}\left(\xi^{N}\right)^{1-\varepsilon}\|f\|_{H^{0}(\Omega)}\|v\|_{H^{1}(\Omega)}, \quad \forall v \in \stackrel{\circ}{S}(\Omega)
$$

Taking $v=\left(T_{h}-T_{h}^{N}\right) f \in \stackrel{\circ}{S}(\Omega)$, and using the Lemma 3.1 we obtain

$$
\left\|\left(T_{h}-T_{h}^{N}\right) f\right\|_{H^{1}(\Omega)} \leqq \frac{C(\varepsilon)}{\alpha}\left(\xi^{N}\right)^{1-\varepsilon}\|f\|_{H^{0}(\Omega)} .
$$

Moreover,

$$
\left\|\left(T_{h}-T_{h}^{N}\right) f\right\|_{H^{0}(\Omega)} \leqq\left\|\left(T_{h}-T_{h}^{N}\right) f\right\|_{H^{1}(\Omega)} .
$$

Therefore, we get (3.19).
From (3.14) and (3.19) we obtain the following result.
Theorem 3.1: For fixed $0<\varepsilon<1$, there exist two constants $C$ and $C(\varepsilon)$ such that

$$
\begin{equation*}
\left\|T-T_{h}^{N}\right\| \leqq C h^{2}+\frac{C(\varepsilon)}{\alpha}\left(\xi^{N}\right)^{1-\varepsilon} \tag{3.21}
\end{equation*}
$$

Suppose that $\mu$ is a non-zero eigenvalue of $T$ with algebraic multiplicity $m$ and $\delta>0$ is a fixed constant such that the interval $[\mu-\varepsilon, \mu+\varepsilon]$ contains no other eigenvalues of $T$. Then there is an $h_{0}$ such that for $0<h \leqq h_{0}$ there exists an integer $N_{0}(h)$ such that if $N \geqq N_{0}(h)$, then there are exactly $m$ eigenvalues (counting algebraic multiplicities) of $T_{h}^{N}$ lying on $[\mu-\delta, \mu+\delta]$. All other eigenvalues of $T_{h}$ are located beyond $[\mu-\delta, \mu+\delta]$ (cf. [6]). In this case the operators $T$ and $T_{h}^{N}$ are self-adjoint. Let $\mu_{1}^{N}(h), \mu_{2}^{N}(h), \ldots, \mu_{m}^{N}(h)$ denote the eigenvalues of $T_{h}^{N}$ lying on the interval $[\mu-\delta, \mu+\delta]$. From Theorem 3 in [6] and Theorem 3.1 we have the following theorem.

Theorem 3.2: There are two constants $C$ and $C(\varepsilon)$ such that, for $0<h<h_{0}$, $N \geqq N_{0}(h)$

$$
\begin{equation*}
\max _{1 \leqslant i \leqslant m}\left|\mu-\mu_{i}^{N}(h)\right| \leqq C h^{2}+\frac{C(\varepsilon)}{\alpha}\left(\xi^{N}\right)^{1-\varepsilon} \tag{3.22}
\end{equation*}
$$

With regard to the eigenfunctions, from Theorem 4 in [6], we have :
Theorem 3.3: Let $\mu^{N}(h)$ be an eigenvalue of $T_{h}^{N}$ such that $\mu^{N}(h) \rightarrow \mu$ as $h \rightarrow 0, N \rightarrow \infty$ and suppose that, for each pair $h, N, w$ is a unit eigenfunction satisfying $\left(\mu^{N}(h)-T_{h}^{N}\right) w=0$. Then there is a function $u$ which is an eigenfunction of $T$ corresponding to $\mu$ such that

$$
\begin{equation*}
\|u-w\|_{H^{0}(\Omega)} \leqq C h^{2}+\frac{C(\varepsilon)}{\alpha}\left(\xi^{N}\right)^{1-\varepsilon} \tag{3.23}
\end{equation*}
$$

where $C$ and $C(\varepsilon)$ are two constants independent of $h$ and $N$.

## 4. CHOOSING THE INTEGER N

The formula (3.22) and (3.23) are the error estimates of the infinite element method for eigenvalue problems. Each of them contains two terms $C h^{2}$ and $\frac{C(\varepsilon)}{\alpha}\left(\xi^{N}\right)^{1-\varepsilon}$. Therefore only when $h$ is sufficiently small and $N$ is sufficiently large we can get accurate approximate solutions. Now we shall discuss a question regarding how to choose $N$ for a fixed partition so that

$$
\left.\begin{array}{c}
\left|\mu-\mu_{i}^{N}(h)\right| \leqslant C h^{2}, \quad i=1,2, \ldots, m  \tag{4.1}\\
\left\|w^{h, N}-u\right\|_{H^{0}(\Omega)} \leqslant C h^{2}
\end{array}\right\} .
$$

From (3.22), (3.16) and (3.23) we know that either $\left(\xi^{N}\right)^{1-\varepsilon}=h^{2}$ or $\left(\xi^{N}\right)^{1-\varepsilon}=C h^{2}$ needs to be chosen (where $c$ is constant) and that the constant $\xi$
is dependent on $h$, i.e., $\xi(h) \rightarrow 1$ when $h \rightarrow 0$. Furthermore, we have the following Lemma.

Lemma 4.1 : Suppose that the partitions satisfy the smallest interior angle condition. Then there exist two positive constants $C_{2}$ and $C_{3}$ which are independent of h such that

$$
\begin{equation*}
C_{2} h \leqslant 1-\xi \leqslant C_{3} h . \tag{4.1}
\end{equation*}
$$

Proof : We consider the triangle of which the length of one side is equal to $h$. This triangle belongs to a quadrangle of the first layer. Let $\triangle A_{i} A_{i+1} A B$ denote this quadrangle.


Figure 2.
Let

$$
L_{\max }=\max \left\{\left|A_{0} A_{1}\right|,\left|A_{0} A_{2}\right|, \ldots,\left|A_{0} A_{M}\right|\right\}
$$

and

$$
L_{\min }=\min \left\{\left|A_{0} A_{1}\right|,\left|A_{0} A_{2}\right|, \ldots,\left|A_{0} A_{m}\right|\right\}
$$

i) If $\left|A A_{i+1}\right|=h$ or $\left|B A_{i}\right|=h$, then

$$
L_{\min }(1-\xi) \leqslant\left|A A_{i+1}\right|\left(\text { or }\left|B A_{i}\right|\right) \leqslant L_{\max }(1-\xi) .
$$

Taking $C_{2}=\frac{1}{L_{\max }}$ and $C_{3}=\frac{1}{L_{\min }}$, we obtain (4.1).
ii) $|B A|=h$ is impossible, because of $\left|A_{i} A_{i+1}\right|>|B A|$.
iii) If $\left|A_{i} A_{i+1}\right|=h$, we have

$$
\begin{equation*}
\frac{\left|A_{i} A_{i+1}\right|}{\sin { }_{\llcorner } A_{i} A A_{i+1}}=\frac{\left|A_{i+1} A\right|}{\sin { }_{\perp} A_{i+1} A_{i} A} . \tag{4.2}
\end{equation*}
$$

Because of $\left|A_{i} A_{i+1}\right|=h \geqslant\left|A_{i+1} A\right|$, then ${ }_{1} A_{i} A A_{i+1} \geqslant{ }_{\llcorner } A_{i+1} A_{i} A$. Namely, $\pi-\gamma_{0} \geqslant{ }_{1} A_{i} A A_{1+1} \geqslant{ }_{1} A_{i+1} A_{i} A \geqslant \gamma_{0}$. From (4.2) we get

$$
\frac{(1-\xi)}{h}=\frac{\sin _{{ }_{1} A_{i+1} A_{i} A}^{\sin ^{1} A_{i} A A_{i+1}}\left|A_{0} A_{i+1}\right| . ~ . ~ . ~}{\text {. }}
$$

Consequently,

$$
L_{\min } \sin \gamma_{0} \leqslant \frac{(1-\xi)}{h} \leqslant \frac{L_{\max }}{\sin \gamma_{0}} .
$$

In this case $C_{2}=L_{\text {min }} \sin \gamma_{0}$ and $C_{3}=\frac{L_{\text {max }}}{\sin \gamma_{0}}$.
iv) If $\left|A A_{A}\right|=h$, similarly we get some results as in case (iii). This proof is completed.
Let $1-\xi=\Delta$. From $\left(\xi^{N}\right)^{1-\varepsilon}=h^{2}$ we obtain $(1-\varepsilon) N \ln (1-\Delta)=2 \ln h$. The conditions $h \ll 1$, and $\Delta \ll 1$, lead to $\ln (1-\Delta) \approx-\Delta$. Furthermore, we have $(1-\varepsilon) N \Delta \cong-2 \ln h$. Namely,

$$
N \approx \frac{-2 \ln h}{(1-\varepsilon) \Delta}
$$

By means of Lemma 4.1, we obtain

$$
\begin{equation*}
\frac{1}{(1-\varepsilon) C_{3}}\left(-\frac{2 \ln h}{h}\right) \leqslant N \leqslant \frac{1}{(1-\varepsilon) C_{2}}\left(-\frac{2 \ln h}{h}\right) . \tag{4.3}
\end{equation*}
$$

From (4.3) we know that, if $N \approx C\left(-\frac{\ln h}{h}\right)$, then we can get the error estimates (4.1). In this case the dimension of matrices $Q_{1}^{N}$ and $Q_{2}^{N}$ is $M \times N$ which can be approximated by $C\left(\frac{-\ln h}{h^{2}}\right)$. It is shown that, in order to get the approximate solutions of Equation (1.6) satisfying the error estimates (4.1), we only need to calculate an eigenvalue problem of a symmetric matrix with the dimension which is approximately $C\left(\frac{-\ln h}{h^{2}}\right)$.

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## REFERENCES

[1] H Han, The infinite element method for eigenvalue problems Tech TR-80-70, MD80-80-HH, October 1980 Department of Mathematics, Unıversity of Maryland
[2] L A Ying, The infinte similar element method for calculating stress intensity factors Scientia Sinica, 21, 1978, 19-43
[3] H Han and L A Ying, An iterative method in the infinite element Mathematicae Numericae Sinica, 1, 1979, 91-99
[4] V A Kondrat'ev, Boundary problems for elliptic equations with conical or angular points Trans Moscow Math Soc, 16, 1967, 227-313
[5] L A Ying, On the convergence of infinite similar element method Acta Mathematicae Applagatae Sinica, 2, 1979, 149-166
[6] J H Bramble and J E Osborn, Rate of convergence estimates for non-self-adjoint elgenvalue approximations Math Comp, 27, 1973, 525-549
[7] R A Adams, Sobolev spaces New York, Academic Press, 1975
[8] H Han, The numerical solutions of interface problems by infinte element method Tech TR-80-7, MD80-7-HH, February 1980 Department of Mathematics, Unıversity of Maryland
[9] P Destuynder and M Diaoua, Estimation de l'erreur sur le coefficient de la singularite de la solution d'un probleme elliptique sur un ouvert avec com R A I R O Analyse Numerique/Numberical Analysis, Vol 14, no 3, 1980, pp 239 a 248
[10] I Babuška, R B Kellogg and J Pitkaranta, Direct and inverse error estimates for finite elements with mesh refinements Numer Math, 33, 1979, 447-471
[11] R W Thatcher, The use of infinite grid refinements at singularities in the solution of Laplace's equation Numer Math 25, 1976, 163-178


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