## Annales de l'institut Fourier

H. Hueber<br>M. Sieveking<br>Uniform bounds for quotients of Green functions on $C^{1,1}$-domains

Annales de l'institut Fourier, tome 32, no 1 (1982), p. 105-117
[http://www.numdam.org/item?id=AIF_1982__32_1_105_0](http://www.numdam.org/item?id=AIF_1982__32_1_105_0)
© Annales de l'institut Fourier, 1982, tous droits réservés.
L'accès aux archives de la revue « Annales de l'institut Fourier » (http://annalif.ujf-grenoble.fr/) implique l'accord avec les conditions générales d'utilisation (http://www.numdam.org/conditions). Toute utilisation commerciale ou impression systématique est constitutive d'une infraction pénale. Toute copie ou impression de ce fichier doit contenir la présente mention de copyright.

Numdam

# UNIFORM BOUNDS FOR QUOTIENTS OF GREEN FUNCTIONS ON $\mathrm{C}^{1,1}$-DOMAINS 

by H. HUEBER and M. SIEVEKING

Consider a partial differential operator L on $\mathbf{R}^{n}$ which has the form

$$
\mathrm{L} u=\sum_{i, j=1}^{n} a_{i j} \mathrm{D}_{i j} u+\sum_{i=1}^{n} b_{i} \mathrm{D}_{i} u+c u .
$$

We assume that L is strictly elliptic and has Hölder continuous coefficients. We will also assume $c \leqslant 0$. Hence every bounded domain $\Omega$ has a Green function which we will denote by $\mathrm{G}_{\mathrm{L}}^{\Omega}$. As usual $\Delta$ denotes the Laplace operator. The aim of our paper is to prove the following theorem:

Theorem. - For every bounded $\mathrm{C}^{1,1}$-domain $\Omega$ there exists a constant C such that we have

$$
\mathrm{C}^{-1} \mathrm{G}_{\Delta}^{\Omega} \leqslant \mathrm{G}_{\mathrm{L}}^{\Omega} \leqslant \mathrm{CG}_{\Delta}^{\Omega}
$$

on $\Omega \times \Omega$. The constant C may be chosen depending only on the ellipticity constant of L , on the Hölder norms of the coefficients of L , on the diameter of $\Omega$, and on the curvature of $\partial \Omega$.

We proceed as follows: In § 1 we introduce notations concerning the geometry of $\Omega$ and give two elementary lemmas. In § 2 we introduce the harmonic space and the adjoint harmonic space associated to L and quote some results which we will use in our proof. The proof is given in $\S 3$. We finish our paper with some remarks which include an application of our result to the Dirichlet problem for $\mathrm{L} u=f$.

## 1. The geometrical situation.

For $r>0$ and $x \in \mathrm{R}^{n}$ let $\mathrm{B}_{r}(x)$ denote the open ball with radius $r$ and center $x$ and let $\mathrm{S}_{r}(x):=\mathrm{B}_{r}(x) \backslash \mathrm{B}_{\frac{1}{3} r}(x)$. Throughout this paper $\Omega$ will denote a bounded domain in $R^{n}$ which is of class $\mathrm{C}^{1,1}$. The distance from $x \in \mathrm{R}^{n}$ to $\partial \Omega$ will be denoted by $d_{x}$.

For every $x_{0} \in \partial \Omega$ there exists a unique inner normal $n\left(x_{0}\right)$ at $\partial \Omega$ with $\left|n\left(x_{0}\right)\right|=1$. By definition the mapping $x_{0} \longrightarrow n\left(x_{0}\right)$ fulfills a local Lipschitz condition, and a simple compactness argument shows that there exists a constant $\mathrm{C}_{\Omega} \geqslant 1$ such that we have $\left|n\left(x_{0}\right)-n\left(y_{0}\right)\right| \leqslant \mathrm{C}_{\Omega}\left|x_{0}-y_{0}\right|$ for all $x_{0}, y_{0} \in \partial \Omega$. Furthermore there exists $\left.r_{\Omega} \in\right] 0,1\left[\right.$ such that we have $\mathrm{B}_{r_{\Omega}}\left(x_{0}+r_{\Omega} n\left(x_{0}\right)\right) \subset \Omega$ and $\mathrm{B}_{r_{\Omega}}\left(x_{0}-r_{\Omega} n\left(x_{0}\right)\right) \subset \mathrm{R}^{n} \backslash \Omega$ for all $x_{0} \in \partial \Omega$. We may assume $\mathrm{C}_{\Omega} r_{\Omega} \leqslant 1$.

For every $x \in \bar{\Omega}$ with $d_{x}<r_{\Omega}$ there exists exactly one point $x_{0} \in \partial \Omega$ with $x=x_{0}+d_{x} n\left(x_{0}\right)$. Hence for such a point $x \in \bar{\Omega}$ the following definitions make sense:

$$
\begin{aligned}
x_{\alpha} & :=x_{0}+\alpha n\left(x_{0}\right) \quad \alpha \in \mathbf{R} \\
n(x) & :=n\left(x_{0}\right) \\
\mathrm{Z}_{\alpha}(x) & :=\left\{y \in \Omega| | y-x_{\beta} \mid<2^{-12} \alpha \text { for some } \beta \in\right]-\alpha, \alpha[ \} \quad \alpha \in \mathbf{R}_{+} .
\end{aligned}
$$

Lemma 1.1. - Let $x, z \in \bar{\Omega}$ with $d_{x}, d_{z}<r_{\Omega}$. Let $\beta, \alpha, \lambda \in \mathbf{R}^{+}$ such that $0<\beta<\alpha<\frac{1}{2} r_{\Omega}$ and $0<\lambda \leqslant \frac{1}{10}$. Assume further $\left|z_{0}-x_{\beta}\right| \leqslant \beta+2 \lambda \alpha$. Then one has $\left|z_{0}-x_{0}\right| \leqslant 4 \alpha \sqrt{\lambda}$.

Proof. - Regarding the situation in a plane which contains $x_{0}$, $x_{r_{\Omega}}$ and $z_{0}$ the problem becomes a two-dimensional one. The statement then follows from the theorem of Pythagoras and elementary estimates which we will leave to the reader.

Lemma 1.2. - a) Let $x \in \Omega$ with $d(x, \partial \Omega)<r_{\Omega}$, let $\alpha \in] 0, \frac{1}{2} r_{\Omega}\left[\right.$ and let $z \in Z_{\alpha} \overline{(x)}$. Then we have $\left|z_{0}-x_{0}\right| \leqslant \frac{1}{16} \alpha$.
b) Let $x, y \in \Omega$ with $d_{x}<r_{\Omega}$ and $d_{y}<r_{\Omega}$. Assume

$$
\overline{\mathrm{Z}_{\alpha}(x)} \cap \overline{\mathrm{Z}_{\alpha}(y)} \neq \varnothing
$$

for some $\alpha \in] 0, r_{\Omega}\left[\right.$. Hence we have $\left|x_{\beta}-y_{\beta}\right|<\frac{1}{4} \alpha$ for all $\beta \in\left[0, r_{\Omega}\right]$.

Proof. - a) Since $z \in Z_{\alpha}(x)$ there exists $\left.\beta \in\right]-\alpha, \alpha[$ such that $\left|x_{\beta}-z\right|<2^{-12} \alpha$. Hence we have

$$
\begin{aligned}
\left|z_{0}-x_{\beta}\right| & \leqslant\left|z_{0}-z\right|+\left|z-x_{\beta}\right| \leqslant\left|x_{0}-z\right|+\left|z-x_{\beta}\right| \\
& <\left|x_{0}-x_{\beta}\right|+2\left|z-x_{\beta}\right|<|\beta|+2 \cdot 2^{-12} \alpha .
\end{aligned}
$$

Let $\lambda=2^{-12}$. Since $z_{0} \notin \mathrm{~B}_{r_{\Omega}}\left(x_{r_{\Omega}}\right)$ the assertion follows from Lemma 1.1 at least for $\beta \geqslant 0$. For $\beta<0$ the assertion is obviously true.
b) Let $z \in \overline{Z_{\alpha}(x)} \cap \overline{Z_{\alpha}(y)}$. By a) we have

$$
\left|y_{0}-x_{0}\right| \leqslant\left|y_{0}-z_{0}\right|+\left|z_{0}-x_{0}\right| \leqslant \frac{1}{8} \alpha
$$

Hence for $\beta \in\left[0, r_{\Omega}\right]$ we get

$$
\begin{aligned}
\left|x_{\beta}-y_{\beta}\right| & =\left|x_{0}+\beta n\left(x_{0}\right)-y_{0}-\beta n\left(y_{0}\right)\right| \\
& \leqslant\left|x_{0}-y_{0}\right|+\beta\left|n\left(x_{0}\right)-n\left(y_{0}\right)\right| \leqslant\left|x_{0}-y_{0}\right|\left(1+\beta \mathrm{C}_{\Omega}\right) \\
& \leqslant \frac{1}{8} \alpha\left(1+r_{\Omega} \mathrm{C}_{\Omega}\right) \leqslant \frac{1}{4} \alpha
\end{aligned}
$$

## 2. Potential theory of $L$.

Following [2] we say that the operator

$$
\mathrm{L}=\sum_{i, j=1}^{n} a_{i j} \mathrm{D}_{i j}+\sum_{i=1}^{n} b_{i} \mathrm{D}_{i}+c
$$

belongs to the class $\mathscr{E}\left(\lambda, \alpha_{0}\right)$ with $\lambda \geqslant 1$ and $\left.\alpha_{0} \in\right] 0,1[$ if it fulfills the following properties:
(i) For all $x \in \mathbf{R}^{n}$ and all $\xi \in \mathbf{R}^{n}$ we have

$$
\sum_{i, j=1}^{n} a_{i j}(x) \xi_{i} \xi_{j} \geqslant \lambda^{-1}|\xi|^{2}
$$

(ii) For all $x, y \in \mathbf{R}^{n}$ we have

$$
\begin{aligned}
\sum_{i, j=1}^{n}\left|a_{i j}(x)-a_{i j}(y)\right|+\sum_{i=1}^{n} \mid b_{i}(x)- & b_{i}(y) \mid \\
& +|c(x)-c(y)| \leqslant \lambda|x-y|^{\alpha_{0}}
\end{aligned}
$$

(iii) For all $x \in \mathbf{R}^{n}$ we have

$$
\sum_{i, j=1}^{n}\left|a_{i j}(x)\right|+\sum_{i=1}^{n}\left|b_{i}(x)\right|+|c(x)| \leqslant \lambda
$$

From now on we will always assume $L \in \mathscr{f}\left(\lambda, \alpha_{0}\right)$, and furthermore we will assume $c \leqslant 0$.

The sheaf $\mathcal{H}_{\mathrm{L}}$ of solutions of $\mathrm{L} u=0$ is a harmonic sheaf which gives rise to a Brelot space ([3], [4], [6]). Every bounded domain V has a Green function $\mathrm{G}_{\mathrm{L}}^{\mathrm{V}}$ which may be characterized by the following properties:
(i) $\mathrm{G}_{\mathrm{L}}^{\mathrm{V}}(\cdot, y)$ is a potential on V for every $y \in \mathrm{~V}$. Its support is $\{y\}$.
(ii) For all $y \in \mathrm{~V}$ we have

$$
\lim _{x \rightarrow y} \frac{\mathrm{G}_{\mathrm{L}}^{\mathrm{V}}(x, y)}{\left(\sqrt{(x-y) \mathrm{B}(y)(x-y))^{2-n}}\right.}=(n-2) n \omega_{n} \sqrt{\operatorname{det}\left(a_{i j}(y)\right)}\left(^{1}\right) .
$$

Here $\mathrm{B}(y)$ is the inverse of $\left(a_{i j}(y)\right)_{i, j}$ and $\omega_{n}$ denotes as usual the volume of the unit ball in $\mathbf{R}^{n}$.
R.M. Herve [6] has shown that there exists a unique Brelot space on $\mathbf{R}^{n}$ such that the extremal potentials on V for this adjoint harmonic space are just the functions $\mathrm{G}_{\mathrm{L}}^{\mathrm{V}}(x, \cdot)$ with $x \in \Omega$. We use the symbol $L^{*}$ to indicate this adjoint space. If the adjoint operator of $L$ exists as a differential operator and is denoted by $L^{*}$ the notation will stay consistent.

For every L-regular bounded open set $V \subset R^{n}$ and every $f \in \mathrm{e}(\partial \mathrm{V})$ we denote by $\mathrm{H}_{\mathrm{V}}^{\mathrm{L}} f$ the unique continuous function on $\overline{\mathrm{V}}$ which is equal to $f$ on $\partial \mathrm{V}$ and L-harmonic in V . For $\mathrm{L}^{*}$-regular V the meaning of $\mathrm{H}_{\mathrm{V}}^{\mathrm{L}^{*}} f$ is obvious. For our purposes it is enough to know that Lipschitz domains are L-regular and $L^{*}$ regular (cf. [1] Théorème 4.1), and we will use this without further comment.

It is clear that $\Delta \in \mathscr{F}(n, \alpha)$ for every $\alpha \in] 0,1[$. To be able to work with the same constants for $L$ and $\Delta$ we will henceforth assume $\lambda \geqslant n$.

[^0]Lemma 2.1. - There exists a constant $\mathrm{C}_{1}=\mathrm{C}_{1}\left(\lambda, \alpha_{0}, n\right)$ such that for every $\left.x \in \mathbf{R}^{n}, r \in\right] 0,1\left[\right.$ and $f \in \mathbb{C}_{+}\left(\partial \mathrm{B}_{r}(x)\right)$ the following inequalities hold on $\mathrm{B}:=\mathrm{B}_{r}(x)$ :

$$
\begin{aligned}
& \mathrm{C}_{1}^{-1} \mathrm{H}_{\mathrm{B}}^{\Delta} f \leqslant \mathrm{H}_{\mathrm{B}}^{\mathrm{L}} f \leqslant \mathrm{C}_{1} \mathrm{H}_{\mathrm{B}}^{\Delta} f \\
& \mathrm{C}_{1}^{-1} \mathrm{H}_{\mathrm{B}}^{\Delta} f \leqslant \mathrm{H}_{\mathrm{B}}^{\mathrm{L}^{*}} f \leqslant \mathrm{C}_{1} \mathrm{H}_{\mathrm{B}}^{\Delta} f .
\end{aligned}
$$

Proof. - The first inequality is due to Serrin and the second one is due to Ancona. They can be found in [2] page 13 and page 20 (for $r \leqslant 1$ the constant C in Proposition 10 on page 20 does not depend on $r$ ).

Lemma 2.2. - There exists a constant $\mathrm{C}_{2}:=\mathrm{C}_{2}\left(\lambda, \alpha_{0}, n\right)$ such that for every $\left.x \in \mathbf{R}^{n}, r \in\right] 0,1[$, and every positive L-harmonic or $\mathrm{L}^{*}$-harmonic function $h$ on $\mathrm{B}_{r}(x)$ the following inequality holds for all $y_{1}, y_{2} \in \mathrm{~B}_{\frac{1}{2} r}(x): \quad h\left(y_{1}\right) \leqslant \mathrm{C}_{2} h\left(y_{2}\right)$.

Proof. - With the use of the ordinary Harnack inequality this statement follows immediately from Lemma 2.1.

Lemma 2.3. - Let $x \in \bar{\Omega}$ with $d_{x}<r_{\Omega}$ and let $\left.\alpha \in\right] 0, \frac{1}{2} r_{\Omega}[$. Let $h$ be a positive L-harmonic or $L^{*}$-harmonic function in $\mathrm{Z}_{\alpha}(x)$ which tends to zero at each $y \in \overline{\mathrm{Z}_{\alpha}(x)} \cap \partial \Omega$. Then we have

$$
\mathrm{C}_{3}^{-1} \frac{\gamma}{\alpha} h\left(x_{\alpha}\right) \leqslant h\left(x_{\gamma}\right) \leqslant \mathrm{C}_{3} \frac{\gamma}{\alpha} h\left(x_{\alpha}\right)
$$

for all $\gamma \in[0, \alpha]$ with a constant $C_{3}$ depending only on $\lambda, \alpha_{0}, n$ and $r_{\Omega}$.

Proof. - In the following the symbol L has to be replaced by $\mathrm{L}^{*}$ if $h$ is $\mathrm{L}^{*}$-harmonic.

Let $\beta:=5^{-1} 2^{-12} \alpha$ and let

$$
\Gamma:=\left\{z| | z-x_{\beta}\left|=\beta ;\left|z-x_{2 \beta}\right| \leqslant \beta\right\} .\right.
$$

By Lemma 2.2 we have $h\left(x_{\beta}\right) \leqslant \mathrm{C}_{2} \inf h(\Gamma)$.
By [1] Théorème 2.2 and Théorème 4.1 there exists a constant $\mathrm{C}=\mathrm{C}\left(\lambda, \alpha_{0}, n, r_{\Omega}\right)$ such that

$$
\sup h\left(\mathrm{~B}_{4 \beta}\left(x_{0}\right) \cap \Omega\right) \leqslant \mathrm{C} h\left(x_{\beta}\right)
$$

Using the maximum principle we get

$$
\mathrm{C}_{2}^{-1} h\left(x_{\beta}\right) \mathrm{H}_{\mathrm{B}_{\beta}\left(x_{\beta}\right)}^{\mathrm{L}} 1_{\Gamma}\left(x_{\gamma}\right) \leqslant h\left(x_{\gamma}\right) \leqslant \mathrm{C} h\left(x_{\beta}\right) \mathrm{H}_{\mathrm{S}_{3 \beta}\left(x_{-\beta}\right)}^{\mathrm{L}} 1_{\partial \mathrm{B}_{3 \beta}\left(x_{-\beta}\right)}\left(x_{\gamma}\right)
$$

for all $\gamma \in[0, \beta]$. Using Lemma 2.1 we get

$$
\begin{aligned}
\mathrm{H}_{\mathrm{B}_{\beta}\left(x_{\beta}\right)}^{\mathrm{L}} 1_{\Gamma}\left(x_{\gamma}\right) \geqslant \mathrm{C}_{1}^{-1} \mathrm{H}_{\mathrm{B}_{\beta}\left(x_{\beta}\right)}^{\Delta} 1_{\Gamma}( & \left.x_{\gamma}\right) \\
& \geqslant \widetilde{\mathrm{C}}_{1}^{-1} \beta^{-1}\left|x_{\gamma}-x_{0}\right|=\widetilde{\mathrm{C}}_{1}^{-1} \frac{\gamma}{\beta}
\end{aligned}
$$

with $\widetilde{\mathrm{C}}_{1}=\widetilde{\mathrm{C}}_{1}\left(\lambda, \alpha_{0}, n\right)$. A standard transportation argument together with Lemma 2.1 shows

$$
\mathrm{H}_{\mathrm{S}_{3 \beta}\left(x_{-\beta}\right)}^{\mathrm{L}} 1_{\partial \mathrm{B}_{3 \beta}\left(x_{-\beta}\right)}\left(x_{\gamma}\right) \leqslant \widetilde{\mathrm{C}}_{2} \frac{\gamma}{\beta}
$$

with $\widetilde{\mathrm{C}}_{2}=\widetilde{\mathrm{C}}_{2}\left(\lambda, \alpha_{0}, n\right)$ (cf. [2] Théorème 2).
Thus we have

$$
\mathrm{C}_{2}^{-1} \widetilde{\mathrm{C}}_{1}^{-1} \frac{\gamma}{\beta} h\left(x_{\beta}\right) \leqslant h\left(x_{\gamma}\right) \leqslant \mathrm{C} \widetilde{\mathrm{C}}_{2} \frac{\gamma}{\beta} h\left(x_{\beta}\right)
$$

for all $\gamma \in[0, \beta]$.

Lemma 2.2 implies that there exists a constant $\widetilde{\mathrm{C}}_{3}$ such that we have $h\left(x_{\gamma_{1}}\right) \leqslant \widetilde{\mathrm{C}}_{3} h\left(x_{\gamma_{2}}\right)$ for all $\gamma_{1}, \gamma_{2} \in[\beta, \alpha]$. Thus for $\gamma \in[0, \beta]$ we have

$$
5.2^{12} \widetilde{\mathrm{C}}_{1}^{-1} \mathrm{C}_{2}^{-1} \widetilde{\mathrm{C}}_{3}^{-1} \frac{\gamma}{\alpha} h\left(x_{\alpha}\right) \leqslant h\left(x_{\gamma}\right) \leqslant 5.2^{12} \mathrm{C} \widetilde{\mathrm{C}}_{2} \widetilde{\mathrm{C}}_{3} \frac{\gamma}{\alpha} h\left(x_{\alpha}\right)
$$

and for $\gamma \in[\beta, \alpha]$ we have

$$
\begin{aligned}
\gamma \alpha^{-1} \widetilde{\mathrm{C}}_{3}^{-1} h\left(x_{\alpha}\right) & \leqslant \widetilde{\mathrm{C}}_{3}^{-1} h\left(x_{\alpha}\right) \leqslant h\left(x_{\gamma}\right) \leqslant \widetilde{\mathrm{C}}_{3} h\left(x_{\alpha}\right) \\
& \leqslant \gamma \alpha^{-1}\left(\alpha \beta^{-1}\right) \widetilde{\mathrm{C}}_{3} h\left(x_{\alpha}\right)=\gamma \alpha^{-1} 5 \cdot 2^{12} \widetilde{\mathrm{C}}_{3} h\left(x_{\alpha}\right)
\end{aligned}
$$

This shows that the assertion is true with $\mathrm{C}_{3}:=5.2^{12} \mathrm{C} \widetilde{\mathrm{C}}_{1} \mathrm{C}_{2} \widetilde{\mathrm{C}}_{2} \widetilde{\mathrm{C}}_{3}$.

Lemma 2.4. - There exists a constant $\mathrm{C}_{4}=\mathrm{C}_{4}\left(\lambda, \alpha_{0}, n\right)$ such that we have

$$
\mathrm{C}_{4}^{-1}|x-y|^{2-n} \leqslant \mathrm{G}_{\mathrm{L}}^{\Omega}(x, y) \leqslant \mathrm{C}_{4}|x-y|^{2-n}
$$

for all $x, y \in \Omega$ with $|x-y| \leqslant \frac{1}{2} d_{x}$ or $|x-y| \leqslant \frac{1}{2} d_{y}$.
Proof. - This Lemma is essentially due to Gilbarg/Serrin [5]. It can be found in [2] on page 14.

## 3. Proof of the theorem.

Let $K:=\left\{x \in \Omega \left\lvert\, d_{x} \geqslant \frac{1}{16} r_{\Omega}\right.\right\}$. Using Lemma 2.4 one can see that $\mathrm{C}_{4}^{-2} \mathrm{G}_{\Delta}^{\Omega} \leqslant \mathrm{G}_{\mathrm{L}}^{\Omega} \leqslant \mathrm{C}_{4}^{2} \mathrm{G}_{\Delta}^{\Omega}$ is true at least in a neighbourhood of the diagonal of $K \times K$. Since $G_{L}^{\Omega}$ and $G_{\Delta}^{\Omega}$ are strictly positive and continuous off the diagonal there exists a constant $\mathrm{C}_{5}$ such that we have

$$
\text { (*) } \mathrm{C}_{5}^{-1} \mathrm{G}_{\Delta}^{\Omega}(x, y) \leqslant \mathrm{G}_{\mathrm{L}}^{\Omega}(x, y) \leqslant \mathrm{C}_{5} \mathrm{G}_{\Delta}^{\Omega}(x, y)
$$

for all $x, y \in \mathrm{~K}$. Looking very carefully at Lemma 2.2 and Lemma 2.4 it is obvious that $C_{5}$ may be chosen depending only on $\lambda, \alpha_{0}, n, r_{\Omega}$ and the diameter of $\Omega$.

Now let $x, y \in \Omega$. Since all arguments which we will use in the following proof are also true if one interchanges the role of $\mathcal{H}_{\mathrm{L}}$ and $\mathcal{H}_{L^{*}}$ we may assume $d_{x} \leqslant d_{y}$. The proof is divided into three cases:

$$
\begin{aligned}
& \quad \text { I } d_{x} \geqslant \frac{1}{16} r_{\Omega} ; \quad d_{y} \geqslant \frac{1}{4} r_{\Omega} \\
& \text { II } d_{x} \leqslant \frac{1}{16} r_{\Omega} ; \quad d_{y} \geqslant \frac{1}{4} r_{\Omega} \\
& \text { III } d_{y} \leqslant \frac{1}{4} r_{\Omega} .
\end{aligned}
$$

Case I. - In this case we have $x, y \in K$ and our statement follows from (*).

$$
\begin{aligned}
& \text { we get } \\
& \begin{aligned}
\mathrm{G}_{\Delta}^{\Omega}(x, y) & \leqslant \mathrm{C}_{3} \frac{8 \beta}{r_{\Omega}} \mathrm{G}_{\Delta}^{\Omega}\left(x_{\frac{1}{8} r_{\Omega}}, y\right) \leqslant \mathrm{C}_{3} \mathrm{C}_{5} \frac{8 \beta}{r_{\Omega}} \mathrm{G}_{\mathrm{L}}^{\Omega}\left(x_{\frac{1}{8} r_{\Omega}}, y\right) \\
& \leqslant \mathrm{C}_{3}^{2} \mathrm{C}_{5} \mathrm{G}_{\mathrm{L}}^{\Omega}(x, y) \leqslant \mathrm{C}_{3}^{3} \mathrm{C}_{5} \frac{8 \beta}{r_{\Omega}} \mathrm{G}_{\mathrm{L}}^{\Omega}\left(x_{\frac{1}{8} r_{\Omega}}, y\right) \\
& \leqslant \mathrm{C}_{3}^{3} \mathrm{C}_{5}^{2} \frac{8 \beta}{r_{\Omega}} \mathrm{G}_{\Delta}^{\Omega}\left(x_{\frac{1}{8} r_{\Omega}}, y\right) \leqslant \mathrm{C}_{3}^{4} \mathrm{C}_{5}^{2} \mathrm{G}_{\Delta}^{\Omega}(x, y)
\end{aligned}
\end{aligned}
$$

Case II. - In this case we have $y \notin \mathrm{Z}_{\frac{1}{8} r_{\Omega}}(x)$. With $\beta:=d_{x}$

Here we have used (*), and we have applied Lemma 2.3 to the Lharmonic function $\mathrm{G}_{\mathrm{L}}^{\Omega}(\cdot, y)$ and the $\Delta$-harmonic function $\mathrm{G}_{\Delta}^{\Omega}(\cdot, y)$ on $\mathrm{Z}_{\frac{1}{8} r_{\Omega}}(x)$.

Case III. - Let $\beta:=d_{x}$ and $\gamma:=d_{y}$. By assumption we have $0<\beta \leqslant \gamma \leqslant \frac{1}{4} r_{\Omega}$. By $\alpha$ we will denote the minimum of all positive numbers $\delta$ for which we have $\overline{\mathrm{Z}_{\delta}(x)} \cap \overline{\mathrm{Z}_{\delta}(y)} \neq \varnothing$. We will consider four subcases.

Subcase 1: $\alpha \geqslant \frac{1}{4} r_{\Omega}=: \mu$.
We have $y_{\mu}, x_{\mu} \in \mathrm{K}$ and $\mathrm{Z}_{\mu}(x) \cap \mathrm{Z}_{\mu}(y)=\varnothing$. With (*) and Lemma 2.3 we get

$$
\begin{aligned}
\mathrm{G}_{\Delta}^{\Omega}(x, y) & \leqslant \mathrm{C}_{3} \frac{\beta}{\mu} \mathrm{G}_{\Delta}^{\Omega}\left(x_{\mu}, y\right) \\
& \leqslant \mathrm{C}_{3}^{2} \frac{\beta}{\mu} \frac{\gamma}{\mu} \mathrm{G}_{\Delta}^{\Omega}\left(x_{\mu}, y_{\mu}\right) \\
& \leqslant \mathrm{C}_{3}^{2} \mathrm{C}_{5} \frac{\beta}{\mu} \frac{\gamma}{\mu} \mathrm{G}_{\mathrm{L}}^{\Omega}\left(x_{\mu}, y_{\mu}\right) \\
& \leqslant \mathrm{C}_{3}^{3} \mathrm{C}_{5} \frac{\beta}{\mu} \mathrm{G}_{\mathrm{L}}^{\Omega}\left(x_{\mu}, y\right) \\
& \leqslant \mathrm{C}_{3}^{4} \mathrm{C}_{5} \mathrm{G}_{\mathrm{L}}^{\Omega}(x, y) \\
& \vdots \\
& \leqslant \mathrm{C}_{3}^{8} \mathrm{C}_{5}^{2} \mathrm{G}_{\Delta}^{\Omega}(x, y)
\end{aligned}
$$

where the dots indicate that we must repeat our arguments in the inverse order.

Subcase 2: $\frac{1}{4} r_{\Omega} \geqslant \alpha \geqslant \gamma$.
Obviously we have $\overline{\mathrm{Z}_{\alpha}(x)} \cap \overline{\mathrm{Z}_{\alpha}(y)} \neq \varnothing$, and hence Lemma 1.2 implies $\left|x_{\alpha}-y_{\alpha}\right| \leqslant \frac{1}{4} \alpha \leqslant \frac{1}{2} d_{y_{\alpha}}$. We apply Lemma 2.3 and Lemma 2.4 and get

$$
\begin{aligned}
\mathrm{G}_{\Delta}^{\Omega}(x, y) & \leqslant \mathrm{C}_{3}^{2} \frac{\beta}{\alpha} \frac{\gamma}{\alpha} \mathrm{G}_{\Delta}^{\Omega}\left(x_{\alpha}, y_{\alpha}\right) \\
& \leqslant \mathrm{C}_{3}^{2} \mathrm{C}_{4}^{2} \frac{\beta}{\alpha} \frac{\gamma}{\alpha} \mathrm{G}_{\mathrm{L}}^{\Omega}\left(x_{\alpha}, y_{\alpha}\right) \\
& \leqslant \mathrm{C}_{3}^{4} \mathrm{C}_{4}^{2} \mathrm{G}_{\mathrm{L}}^{\Omega}(x, y) \\
& \leqslant \mathrm{C}_{3}^{6} \mathrm{C}_{4}^{2} \frac{\beta}{\alpha} \frac{\gamma}{\alpha} \mathrm{G}_{\mathrm{L}}^{\Omega}\left(x_{\alpha}, y_{\alpha}\right)
\end{aligned}
$$

$$
\begin{aligned}
& \leqslant C_{3}^{6} \mathrm{C}_{4}^{4} \frac{\beta}{\alpha} \frac{\gamma}{\alpha} \mathrm{G}_{\Delta}^{\Omega}\left(x_{\alpha}, y_{\alpha}\right) \\
& \leqslant \mathrm{C}_{3}^{8} \mathrm{C}_{4}^{4} \mathrm{G}_{\Delta}^{\Omega}(x, y)
\end{aligned}
$$

Subcase 3: $\beta \geqslant \alpha$.
If $|x-y| \leqslant \frac{1}{2} \beta$ the desired inequalities follow from Lemma 2.4. Hence we may assume $|x-y|>\frac{1}{2} \beta$. Since $\beta \leqslant \gamma$ we have $y \notin \mathrm{Z}_{\beta}(x)$. From Lemma 1.2 we get $\left|x_{\frac{1}{2} \beta}-y_{\frac{1}{2} \beta}\right| \leqslant \frac{1}{4} \beta$, and $\left|y-y_{\frac{1}{2} \beta}\right| \geqslant \frac{1}{2} \beta$ is clear. Hence we get

$$
\begin{aligned}
\mathrm{G}_{\Delta}^{\Omega}(x, y) & \leqslant \mathrm{C}_{3} 2 \mathrm{G}_{\Delta}^{\Omega}\left(x_{\frac{1}{2} \beta}, y\right) \\
& \leqslant \mathrm{C}_{2} \mathrm{C}_{3} 2 \mathrm{G}_{\Delta}^{\Omega}\left(y_{\frac{1}{2} \beta}, y\right) \\
& \leqslant \mathrm{C}_{2} \mathrm{C}_{3}^{2} 2 \gamma^{-1} \beta \mathrm{G}_{\Delta}^{\Omega}\left(y_{\frac{1}{2} \gamma}, y\right) \\
& \leqslant \mathrm{C}_{2} \mathrm{C}_{3}^{2} \mathrm{C}_{4}^{2} 2 \gamma^{-1} \beta \mathrm{G}_{\mathrm{L}}^{\Omega}\left(y_{\frac{1}{2} \gamma}, y\right) \\
& \leqslant \mathrm{C}_{2} \mathrm{C}_{3}^{3} \mathrm{C}_{4}^{2} 2 \mathrm{G}_{\mathrm{L}}^{\Omega}\left(y_{\frac{1}{2} \beta}, y\right) \\
& \leqslant \mathrm{C}_{2}^{2} \mathrm{C}_{3}^{3} \mathrm{C}_{4}^{2} 2 \mathrm{G}_{\mathrm{L}}^{\Omega}\left(x_{\frac{1}{2} \beta}, y\right) \\
& \leqslant \mathrm{C}_{2}^{2} \mathrm{C}_{3}^{4} \mathrm{C}_{4}^{2} \mathrm{G}_{\mathrm{L}}^{\Omega}(x, y) \\
& \vdots \\
& \leqslant \mathrm{C}_{2}^{4} \mathrm{C}_{3}^{8} \mathrm{C}_{4}^{4} \mathrm{G}_{\Delta}^{\Omega}(x, y)
\end{aligned}
$$

To derive the first of these inequalities we applied Lemma 2.3 to the $\Delta$-harmonic function $\mathrm{G}_{\Delta}^{\Omega}(\cdot, y)$ on $\mathrm{Z}_{\beta}(x)$. In the second line we applied Lemma 2.2 to the same function in $\mathrm{B}_{\frac{1}{2} \beta}\left(y_{\frac{1}{2} \beta}\right)$. In the third inequality we used Lemma 2.3 on $\mathrm{Z}_{\frac{1}{2} \gamma}(y)$ observing that $y \notin \mathrm{Z}_{\frac{1}{2} \gamma}(y)$. The fourth inequality followed from Lemma 2.4. Then we applied again Lemma 2.3, then Lemma 2.2, then Lemma 2.3 once more, and the dots indicate that we must repeat all our arguments in the inverse order.

Subcase 4: $\gamma \geqslant \alpha \geqslant \beta$.

We apply subcase 3 to the points $x_{\alpha}, y$ and get

$$
\mathrm{C}_{2}^{-2} \mathrm{C}_{3}^{-4} \mathrm{C}_{4}^{-2} \mathrm{G}_{\Delta}^{\Omega}\left(x_{\alpha}, y\right) \leqslant \mathrm{G}_{\mathrm{L}}^{\Omega}\left(x_{\alpha}, y\right) \leqslant \mathrm{C}_{2}^{2} \mathrm{C}_{3}^{4} \mathrm{C}_{4}^{2} \mathrm{G}_{\Delta}^{\Omega}\left(x_{\alpha}, y\right)
$$

Since $y \notin \mathrm{Z}_{\alpha}(x)$ we may apply Lemma 2.3 to $\mathrm{Z}_{\alpha}(x)$, and we get

$$
\begin{aligned}
\mathrm{G}_{\Delta}^{\Omega}(x, y) & \leqslant \mathrm{C}_{3} \frac{\beta}{\alpha} \mathrm{G}_{\Delta}^{\Omega}\left(x_{\alpha}, y\right) \\
& \leqslant \mathrm{C}_{2}^{2} \mathrm{C}_{3}^{5} \mathrm{C}_{4}^{2} \frac{\beta}{\alpha} \mathrm{G}_{\mathrm{L}}^{\Omega}\left(x_{\alpha}, y\right) \\
& \leqslant \mathrm{C}_{2}^{2} \mathrm{C}_{3}^{6} \mathrm{C}_{4}^{2} \mathrm{G}_{\mathrm{L}}^{\Omega}(x, y) \\
& \leqslant \mathrm{C}_{2}^{2} \mathrm{C}_{3}^{7} \mathrm{C}_{4}^{2} \frac{\beta}{\alpha} \mathrm{G}_{\mathrm{L}}^{\Omega}\left(x_{\alpha}, y\right) \\
& \leqslant \mathrm{C}_{2}^{4} \mathrm{C}_{3}^{11} \mathrm{C}_{4}^{4} \frac{\beta}{\alpha} \mathrm{G}_{\Delta}^{\Omega}\left(x_{\alpha}, y\right) \\
& \leqslant \mathrm{C}_{2}^{4} \mathrm{C}_{3}^{12} \mathrm{C}_{4}^{4} \mathrm{G}_{\Delta}^{\Omega}(x, y)
\end{aligned}
$$

Now we see that with $C:=C_{2}^{2} C_{3}^{6} C_{4}^{2} C_{5}$ the statement of the theorem is true.

## 4. Remarks.

4.1. Obviously our theorem is true if the Laplace operator is replaced by any $L^{\prime} \in \mathscr{L}\left(\lambda, \alpha_{0}\right)$ with $c^{\prime} \leqslant 0$. Hence

$$
\mathrm{C}\left(\mathrm{~L}, \mathrm{~L}^{\prime}\right):=\inf \left\{c>0 \mid \mathrm{c}^{-1} \mathrm{G}_{\mathrm{L}^{\prime}}^{\Omega} \leqslant \mathrm{G}_{\mathrm{L}}^{\Omega} \leqslant c \mathrm{G}_{\mathrm{L}^{\prime}}^{\Omega}\right\}
$$

is a well defined real number $\geqslant 1$. We can show that for sufficiently smooth domains $C\left(L, L^{\prime}\right)$ depends continuously on the coefficients of $L, L^{\prime}$ with respect to uniform convergence. The proof is by a modification of a method presented in [7].
4.2. For Lipschitz domains our theorem is not true:

Let $\Omega:=\left\{x=\left(x_{1}, x_{2}\right) \in \mathbf{R}^{2} \mid 0<x_{1}<1 ; 0<x_{2}<x_{1}\right\}$; fix $y \in \Omega$ and let V denote the intersection of $\Omega$ and a small ball around zero whose closure does not contain $y$. Hence there exists a constant $C_{\Delta}$ such that

$$
\mathrm{C}_{\Delta}^{-1}\left(x_{1}^{3} x_{2}-x_{1} x_{2}^{3}\right) \leqslant \mathrm{G}_{\Delta}^{\Omega}(x, y) \leqslant \mathrm{C}_{\Delta}\left(x_{1}^{3} x_{2}-x_{1} x_{2}^{3}\right)
$$

for all $x=\left(x_{1}, x_{2}\right) \in \mathrm{V}$.

Let $A:=\frac{\partial^{2}}{\partial x_{1}^{2}}+2 \frac{\partial}{\partial x_{1} \partial x_{2}}+3 \frac{\partial^{2}}{\partial x_{2}^{2}}$. Then there exists $a$ constant $\mathrm{C}_{\mathrm{A}}$ such that

$$
\mathrm{C}_{\mathrm{A}}^{-1}\left(x_{1}^{2}-x_{1} x_{2}\right) \leqslant \mathrm{G}_{\mathrm{A}}^{\Omega}(x, y) \leqslant \mathrm{C}_{\mathrm{A}}\left(x_{1}^{2}-x_{1} x_{2}\right)
$$

This shows that $G_{A}^{\Omega}(\cdot, y) / G_{\Delta}^{\Omega}(\cdot, y)$ is unbounded near zero.
4.3. Let $X$ be a locally compact space with countable base, and assume that we have two harmonic sheafs $\mathcal{H}_{1}, \mathcal{H}_{2}$ on $X$ such that $\left(\mathrm{X}, \mathcal{H}_{i}\right) \quad i=1,2$ is a $\mathscr{P}$-Brelot space. Assume that $\Omega$ is a bounded domain in $X$ which is regular for both structures and happens to have Green functions $G_{x_{i}}^{\Omega}(i=1,2)$. Let $f:=\partial \Omega \longrightarrow \mathbf{R}$ be continuous and positive, and let $u_{i}$ denote the solution of the Dirichlet problem for $f$ with respect to the sheaf $\mathcal{H}_{i}$. We know that $u_{1}$ is the limite of an increasing sequence $\left(p_{n}\right)_{n \in N}$ of $\mathcal{H}_{1}$ potentials, and we may even assume that the potential theoretic carrier of $p_{n}$ is contained in $\Omega \backslash \mathrm{U}_{n}$ where $\left(\mathrm{U}_{n}\right)_{n \in \mathrm{~N}}$ is an exhaustion of $\Omega$.

Hence

$$
p_{n}(x)=\int \mathrm{G}_{x_{1}}^{\Omega}(x, y) \mu_{n}(d y)
$$

with $\mu_{n}\left(\mathrm{U}_{n}\right)=0$.
Let

$$
q_{n}(x):=\int \mathrm{G}_{\mathscr{H}_{2}}^{\Omega}(x, y) \mu_{n}(d y)
$$

Assume now that there exists a constant $C$ such that

$$
\mathrm{C}^{-1} \mathrm{G}_{\mathfrak{x}_{1}}^{\Omega} \leqslant \mathrm{G}_{\varkappa_{2}}^{\Omega} \leqslant \mathrm{CG}_{\mathscr{\varkappa}_{1}}^{\Omega}
$$

Hence

$$
\mathrm{C}^{-1} q_{n} \leqslant p_{n} \leqslant \mathrm{C} q_{n}
$$

By Azela-Ascoli a subsequence of $\left(q_{n}\right)_{n \in N}$ is locally uniformly convergent and therefore we may assume that $\left(q_{n}\right)_{n}$ itself converges locally uniformly to an $\mathcal{H e}_{2}$-harmonic function $q$.

Obviously we have $\mathrm{C}^{-1} q \leqslant u_{1} \leqslant \mathrm{C} q$ and from this we derive $\mathrm{C}^{-1} q \leqslant u_{2} \leqslant \mathrm{C}^{-1} q$. Thus we have $\mathrm{C}^{-2} u_{2} \leqslant u_{1} \leqslant \mathrm{C}^{2} u_{2}$.

This shows that our theorem implies the statement of Lemma 2.1. In a similar way we may also derive the other statements of § 2 from the theorem. This shows that our paper characterizes in a
certain way all Brelot spaces on $\mathbf{R}^{\boldsymbol{n}}$ for which the comparability statement of our theorem is true with a constant $C$ which depends only on the diameter of $\Omega$ and the curvature of $\partial \Omega$.
4.4. For $L_{1}, L_{2} \in \mathscr{L}\left(\lambda, \alpha_{0}\right)$ the following comparability statement is an easy conclusion of 4.1 and 4.3:

For any $u_{1}, u_{2} \in \mathcal{C}(\bar{\Omega}) \cap \mathfrak{C}^{2}(\Omega)$ with $u_{1}=u_{2} \geqslant 0$ on $\partial \Omega$ and $\mathrm{L}_{1} u_{1}=\mathrm{L}_{2} u_{2} \leqslant 0$ in $\Omega$ we have $u_{1} \leqslant \mathrm{C} u_{2}$ with a constant $\mathrm{C}=\mathrm{C}\left(\lambda, \alpha_{0}, n, \Omega\right)$.

## BIBLIOGRAPHY

[1] A. Ancona, Principe de Harnack à la frontière et théorème de Fatou pour un opérateur elliptique dans un domaine Lipschitzien, Ann. Inst. Fourier, 28, 4 (1978), 169-213.
[2] A. Ancona, Principe de Harnack à la frontière et problèmes de frontière de Martin, Lecture Notes in Mathematics, 787 (1980), 9-28.
[3] N. Boboc, P. Mustata, Espaces harmoniques associés aux opérateurs différentiels linéaires du second ordre de type elliptique, Lecture Notes in Mathematics, 68 (1968).
[4] C. Constantinescu, A. Cornea, Potential theory on harmonic spaces, Berlin-Heidelberg-New York, 1972.
[5] D. Gilbarg, J. Serrin, On isolated singularities of solutions of second order elliptic differential equations, J. d'Anal. Math., 4 (1954-1956), 309-340.
[6] R.M. Herve, Recherches axiomatiques sur la théorie des fonctions surhamoniques et du potentiel, Ann. Inst. Fourier, 12 (1962), 415-571.
[7] H. Hueber, M. Sifveking, On the quotients of Green functions (preliminary version), Bielefeld, September 1980 (unpublished).
[8] J. Serrin, On the Harnack inequality for linear elliptic equations, J. d'Anal. Math., 4 (1956), 292-308.
[9] J.-C. Taylor, On the Martin compactification of a bounded Lipschitz domain in a Riemannian manifold, Ann. Inst. Fourier, 28, 2 (1977), 25-52.

Manuscrit reçu le 13 avril 1981.
H. Hueber,

Fakultät für Mathematik
Universität Bielefeld
4800 Bielefeld 1 (R.F.A.)
\&
M. Sievering,

Angewandte Mathematik
Robert-Mayer-Str. 10
6000 Frankfurt/Main (R.F.A.).


[^0]:    (1) For $n \leqslant 2$ this formula has to be modified in the usual way. The same remark applies to the statement of Lemma 2.4.

